

The Monotonicity Method for the Helmholtz equation

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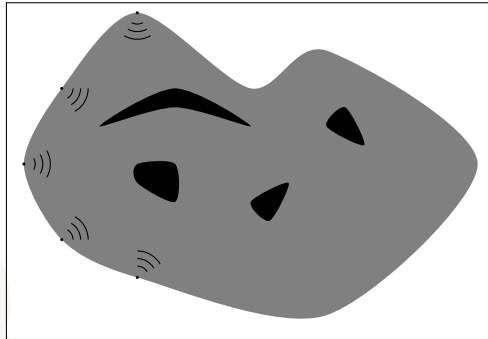
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Conference on Mathematics of Wave Phenomena
Karlsruhe Institute of Technology, July 23–27, 2018.

Scattering in a bounded domain



- ▶ Excite time-harmonic pressure wave in a bounded domain
- ▶ **Aim:** Detect defects/anomalies from scattering response
- ▶ **Applications:** Acoustic/EM tomography, non-destructive testing

Helmholtz equation

- ▶ Time-harmonic wave in bounded Lipschitz domain $\Omega \subset \mathbb{R}^n$ ($n \geq 2$)

$$(\Delta + k^2 q)u = 0 \quad \text{in } \Omega \quad (1)$$

($k > 0$: non-resonant wavenumber, $q \in L^\infty(\Omega)$: sound speed,
 $u \in H^1(\Omega)$: acoustic pressure)

- ▶ Idealized boundary measurements: **Neumann-to-Dirichlet map**

$$\Lambda(q) : L^2(\Sigma) \rightarrow L^2(\Sigma), \quad g \mapsto u|_\Sigma,$$

where u solves (1) with $\partial_\nu u|_{\partial\Omega} = \begin{cases} g & \text{on } \Sigma, \\ 0 & \text{else.} \end{cases}$

($\Sigma \subseteq \partial\Omega$: open boundary part)

Can we recover q from $\Lambda(q)$?

Inversion methods

- ▶ Linearization (Born / single scattering) & Iteration
 - ▶ generic, popular, but no convergence theory
- ▶ Linear Sampling Methods / Factorization Methods
(*Scattering: Colton, Kirsch, Cakoni, Haddar, Arens, Lechleiter, Griesmaier, ...*
EIT: Hanke, Brühl, Hyvönen, H., Seo, ...)
 - ▶ rigorous for exact data, yields uniqueness results
 - ▶ non-intuitive criterion (range/infinity tests)
 - ▶ no convergence theory for noisy data, needs definiteness
- ▶ Monotonicity Method (for EIT)
(*Tamburrino, Rubinacci, H., Ullrich, Mach, Garde, ...*)
 - ▶ rigorous theory (based on FM), yields uniqueness results
 - ▶ simple, convergent for noisy data, can treat indefinite case
 - ▶ can be combined with linearization approach

This talk: Extend monotonicity method to Helmholtz equation

Monotonicity Method (for simple test case in EIT)

- ▶ EIT: Detect $\sigma \in L_+^\infty(\Omega)$ in $\nabla \cdot (\sigma \nabla u) = 0$ from NtD $\Lambda(\sigma)$
- ▶ Inclusion detection: $\sigma = 1 + \chi_D$, D open, $\Omega \setminus \bar{D}$ connected
- ▶ Monotonicity:

$$\sigma_1 \leq \sigma_2 \implies \Lambda(\sigma_1) \geq \Lambda(\sigma_2)$$

(i.e., $\Lambda(\sigma_1) - \Lambda(\sigma_2)$ has **no** negative eigenvalues)

- ▶ Monotonicity for inclusion detection:
(„ \implies “: Tamburrino/Rubinacci 2002, „ \longleftarrow “ & Linearization: H./Ullrich 2013)

$$\begin{aligned} B \subseteq D &\iff \Lambda(1 + \chi_B) \geq \Lambda(1 + \chi_D) \\ &\iff \Lambda(1) + \frac{1}{2}\Lambda'(1)\chi_B \geq \Lambda(1 + \chi_D) \end{aligned}$$

Inclusion can be found by testing several small balls B

Monotonicity Method for Helmholtz (simple version)

- ▶ **Helmholtz:** Detect $q \in L^\infty(\Omega)$ in $(\Delta + k^2 q)u = 0$ from NtD $\Lambda(q)$
- ▶ Scatterer detection: $q = 1 + \chi_D$, D open, $\Omega \setminus \bar{D}$ connected
- ▶ Monotonicity:

$$q_1 \leq q_2 \implies \Lambda(q_1) \leq_{\text{fin}} \Lambda(q_2)$$

(i.e., $\Lambda(q_2) - \Lambda(q_1)$ has **only finitely many** negative eigenvalues)

- ▶ Monotonicity for inverse scattering:

$$\begin{aligned} B \subseteq D &\iff \Lambda(1 + \chi_B) \leq_{\text{fin}} \Lambda(1 + \chi_D) \\ &\iff \Lambda(1) + \Lambda'(1)\chi_B \leq_{\text{fin}} \Lambda(1 + \chi_D) \end{aligned}$$

Scatterer can be found by testing several small balls B

Next slides: Full results under general assumptions

Monotonicity for Helmholtz

Theorem. (H./Pohjola/Salo, submitted)

Let $q_1, q_2 \in L^\infty(\Omega)$, $k > 0$ no resonance. Then

$$q_1 \leq q_2 \quad \text{implies} \quad \Lambda(q_1) \leq_{d(q_2)} \Lambda(q_2),$$

(i.e., $\Lambda(\sigma_2) - \Lambda(\sigma_1)$ has **less than** $d(q_2)$ negative eigenvalues)

- ▶ $d(q_2)$ = no. of positive Neumann EVs of $\Delta + k^2 q$ (always finite)

Larger sound speed leads to larger NtD-measurements

(in the sense of a modified Loewner order)

Local uniqueness for Helmholtz

Theorem. (H./Pohjola/Salo, submitted) **Let**

- ▶ $q_1, q_2 \in L^\infty(\Omega)$, $k > 0$ no resonance,
- ▶ $O \subseteq \overline{\Omega}$ rel. open set connected to Σ with $q_1|_O \leq q_2|_O$.

Then

$$q_1|_O \neq q_2|_O \quad \text{implies} \quad \Lambda(q_1) \not\leq_{\text{fin}} \Lambda(q_2).$$

Deviation in sound speed can be detected

(from eigenvalues in NtD difference)

Scatterer detection (definite case)

$\Lambda(1)$: NtD for homogeneous sound speed

$\Lambda(q)$: NtD for unknown sound speed ($q \in L^\infty(\Omega)$, $k > 0$ no resonance)

$D \subseteq \Omega$: unknown scatterer (open, $\Omega \setminus \bar{D}$ connected)

T_B : test operator for open $B \subseteq \Omega$ ($\int_\Sigma g T_B h := \int_B k^2 u_1^g u_1^h dx$)

Theorem. (H./Pohjola/Salo, submitted)

Let $1 \leq q_{\min} \leq q(x) \leq q_{\max}$ for all $x \in D$ (a.e.), then

$B \subseteq D$ implies $\alpha T_B \leq_{d(q_{\max})} \Lambda(q) - \Lambda(1)$ for all $\alpha \leq q_{\min} - 1$,

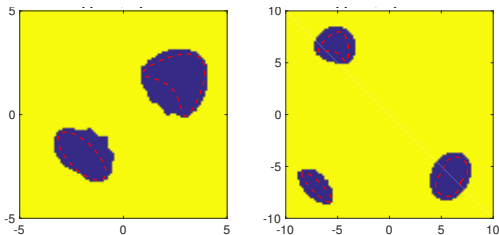
$B \not\subseteq D$ implies $\alpha T_B \not\leq_{\text{fin}} \Lambda(q) - \Lambda(1)$ for all $\alpha > 0$.

(Similar result holds for $q_{\min} \leq q(x) \leq q_{\max} < 1$)

Scatterer can be localized by monotonicity tests

Remarks and Extensions

- ▶ Monotonicity tests require no forward solutions (only for $q_0 \equiv 1$).
- ▶ Tests can be easily regularized (\leadsto convergence for noisy data)
- ▶ Extensions possible for background sound speed $q_0 \neq 1$
- ▶ Extensions possible for $\Omega \setminus \overline{D}$ not connected
(using concept of inner and outer support)
- ▶ Extension possible for indefinite case (by shrinking large test domain)
- ▶ Extension to far-field scattering: (Griesmaier/H., submitted)



Proofs (main ideas): Well-posedness

- ▶ Standard variational formulation: $u \in H^1(\Omega)$ solves

$$(\Delta + k^2 q)u = 0 \quad \text{in } \Omega, \quad \partial_\nu u|_{\partial\Omega} = \begin{cases} g & \text{on } \Sigma, \\ 0 & \text{else,} \end{cases}$$

if and only if

$$b(u, v) := \int_{\Omega} (\nabla u \cdot \nabla v - k^2 q uv) \, dx = \int_{\Sigma} g v|_{\Sigma} \, ds$$

- ▶ $b(\cdot, \cdot)$ is coercive plus compact depending analytically on $k \in \mathbb{C}$
- ▶ Analytic Fredholm theory \leadsto Unique solvability
(except for discrete set of resonance frequencies)

Proofs (main ideas): Monotonicity

- ▶ From the variational formulation one obtains

$$\begin{aligned} & \int_{\Sigma} g (\Lambda(q_2) - \Lambda(q_1)) g \, ds + \int_{\Omega} k^2 (q_1 - q_2) |u_1^{(g)}|^2 \, dx \\ &= \int_{\Omega} \left(|\nabla(u_2^{(g)} - u_1^{(g)})|^2 - k^2 q_2 |u_2^{(g)} - u_1^{(g)}|^2 \right) \, dx. \end{aligned}$$

- ▶ **Right hand side** is coercive plus compact
- ↪ **Right hand side** is non-negative is space of finite codimension
- ↪ Monotonicity inequality

$$\int_{\Sigma} g (\Lambda(q_2) - \Lambda(q_1)) g \, ds \geq_{\text{fin}} \int_{\Omega} k^2 (q_2 - q_1) |u_1^{(g)}|^2 \, dx$$

- ▶ Converse monotonicity by controlling $u_1^{(g)}|_D$ on subset $D \subset \Omega$

Proofs (main ideas): Localized potentials

Localized potentials: Control $u^{(g)}|_D$ on subset $D \subseteq \Omega$

- ▶ Neumann-to-Solution-operator:

$$L_D: L^2(\Sigma) \rightarrow L^2(D), \quad g \mapsto u^{(g)}|_D$$

- ▶ $L_D^*: L^2(D) \rightarrow L^2(\Sigma)$: Source-to-Dirichlet-operator
- ▶ Unique continuation: For "different" subsets $B, D \subseteq \Omega$

$$\mathcal{R}(L_D^*) \cap \mathcal{R}(L_B^*) = 0$$

- ▶ Duality argument: $\exists g_n \in L^2(\Sigma)$:

$$\|u^{(g_n)}|_D\| = \|L_D g_n\| \rightarrow \infty \quad \text{and} \quad \|u^{(g_n)}|_B\| = \|L_B g_n\| \rightarrow 0.$$

- ▶ $\dim \mathcal{R}(L_D^*), \dim \mathcal{R}(L_B^*) = \infty$
 $\leadsto g_n$ can be chosen from space with finite codimension

Modified Loewner order for compact selfadjoint operators:

$$A \leq_d B \quad :\iff \quad B - A \text{ has less than } d \text{ negative EVs}$$

Monotonicity and converse monotonicity for Helmholtz equation:

- ▶ Larger sound speed implies larger NtD measurements.
- ▶ Larger NtD implies that there is no boundary neighbourhood where sound speed is smaller.

Monotonicity approach yields

- ▶ Local uniqueness result for Helmholtz equation
- ▶ Simple but rigorously convergent scatterer detection algorithm