

# Combining frequency-difference EIT with ultrasound modulated EIT (rough sketch of a new idea)

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Forward operator of EIT:

"conductivity"  $\mapsto$  "measurements"

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Main principal difficulties in inverting the forward operator:

- ▶ inverse problem is highly non-linear
- ▶ inverse problem is highly ill-posed
- ▶ typical applications involve large modelling errors (electrode positions, boundary geometry, . . .)

Approaches to alleviate principal difficulties:

- ▶ Large modelling errors
  - ↪ Use difference measurements
- ▶ High non-linearity
  - ↪ Only detect outer shape of inclusions / conductivity changes
- ▶ Ill-posedness
  - ↪ Measure additional interior data (hybrid tomography).

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In this talk: combine all of the above.

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Setting (continuous measurements):

- ▶ Apply current  $g$  to boundary of a body  $\Omega$

↪ Electric potential  $u$  that solves

$$\nabla \cdot \gamma_\omega \nabla u = 0, \quad \gamma_\omega \partial_\nu u|_{\partial\Omega} = g$$

( $\gamma_\omega$ : complex conductivity at frequency  $\omega$ )

- ▶ Measure boundary voltage  $u|_{\partial\Omega}$  for all possible currents  $g$

↪ Current-to-voltage map  $\Lambda(\gamma_\omega) : g \mapsto u|_{\partial\Omega}$ .

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$$\gamma_\omega \in L_+^\infty(\Omega) + iL^\infty(\Omega) \implies \Lambda(\gamma_\omega) \in \mathcal{L}(L_\diamond^2(\partial\Omega), L_\diamond^2(\partial\Omega))$$


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- ▶  $\overline{D} \subset \Omega$ ,  $D$  (the "inclusion")
- ▶ Conductivity at zero and non-zero frequency:

$$\gamma_0 = \gamma_0^\Omega + \gamma_0^D \chi_D, \quad \gamma_\omega = \gamma_\omega^\Omega + \gamma_\omega^D \chi_D$$

- ▶ Roughly:

$$\left\langle g, \left( \Lambda(\gamma_0) - \frac{\gamma_\omega^\Omega}{\gamma_0^\Omega} \Lambda(\gamma_\omega) \right) g \right\rangle \approx \int_D \left( \frac{\gamma_0^\Omega}{\gamma_\omega^\Omega} \gamma_\omega^D - \gamma_0^D \right) |\nabla u_0|^2.$$

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fdEIT: Reconstruct  $D$  from  $\Lambda(\gamma_0) - \frac{\gamma_\omega^\Omega}{\gamma_0^\Omega} \Lambda(\gamma_\omega)$ .

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- ▶ Conductivity at zero frequency:  $\gamma_0$
- ▶ Change conductivity in small area  $B$ :  $\gamma_0(1 + \alpha\chi_B)$ ,  $\alpha > 0$  small
- ▶ Roughly:

$$\frac{1}{\alpha} \langle g, (\Lambda(\gamma_0) - \Lambda(\gamma_0(1 + \alpha\chi_B))) g \rangle \approx \int_B \gamma_0(x) |\nabla u_0(x)|^2 dx,$$

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US-modulated EIT:

$\Lambda(\gamma_0) - \Lambda(\gamma_0(1 + \alpha\chi_B))$  yields interior energy  $\gamma_0(x) |\nabla u_0(x)|^2$ .

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Combining above rough estimates suggests:

Conjecture (rough)

If  $B \subset D$ ,  $\alpha > 0$  small enough

$$B \subseteq D \implies \Lambda(\gamma_0(1 + \alpha\chi_B)) \geq \frac{\gamma_\omega^\Omega}{\gamma_0^\Omega} \Lambda(\gamma_\omega).$$

Idea

Compare measurements at a different frequency with ultrasound-modulated measurements to find inclusions.

## Theorem

Let  $\Re\left(\frac{\gamma_0^\Omega}{\gamma_\omega^\Omega}\gamma_\omega^D\right) > 0$ ,

$D \subseteq \Omega$  open,  $\overline{D} \subseteq \Omega$  has connected complement.  $B \subseteq \Omega$  ball.  
For suff. small  $\alpha$ :

$$B \subseteq D \implies \Lambda(\gamma_0(1 + \alpha\chi_B)) \geq \frac{\gamma_\omega^\Omega}{\gamma_0^\Omega}\Lambda(\gamma_\omega).$$

For all  $\alpha > 0$ :

$$B \not\subseteq D \implies \Lambda(\gamma_0(1 + \alpha\chi_B)) \not\geq \frac{\gamma_\omega^\Omega}{\gamma_0^\Omega}\Lambda(\gamma_\omega)$$



## Why?

- ▶ ratio of background conductivities  $\frac{\sigma_\omega}{\sigma_0}$  usually known  
(or estimated by minimizing  $\Lambda(\gamma_0(1 + \alpha\chi_B)) - r\Lambda(\gamma_\omega)$ )
- ▶  $\Lambda(\gamma_0(1 + \alpha\chi_B))$  and  $\Lambda(\gamma_\omega)$  are measured.

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Approach only requires measurements.

No info on geometry or electrode position needed!

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Monotony tests are stable. Results extend to point electrode models.