

# Detecting Interfaces in a Parabolic-Elliptic Problem from Surface Measurements

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# A parabolic-elliptic problem

Heat equation:  $\partial_t(c(x)u(x, t)) - \nabla \cdot (\kappa(x)\nabla u(x, t)) = 0$  in  $B = Q \cup \bar{\Omega}$

$u(x, t)$ : temperature

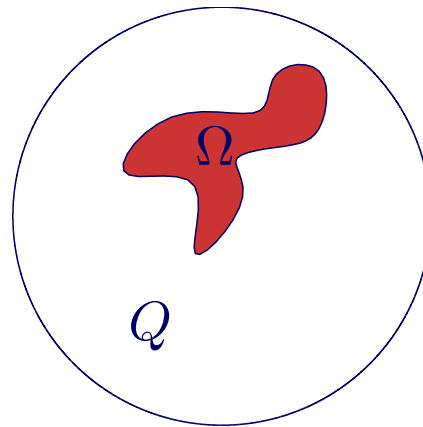
$c(x)$ : heat capacity

$\kappa(x)$ : heat conductivity

Special case:

$$c(x) = \begin{cases} 1 & \text{in } \Omega \\ 0 & \text{in } Q \end{cases}$$

$$\kappa(x) = \begin{cases} \kappa_1 > 1 & \text{in } \Omega \\ 1 & \text{in } Q \end{cases}$$



parabolic equation in  $\Omega$   
("  $\partial_t u - \Delta u = 0$  ")

elliptic equation in  $Q$   
("  $\Delta u = 0$  ")

Motivation:

- Domain with inclusions of much higher heat capacity
- Electrically conducting objects in a non-conducting background illuminated by low-frequency electromagnetic waves

# Direct problem / Inverse Problem

**Direct Problem:** For every heat flux  $g$  there is a unique solution  $u_1$  of

$$\partial_t(\chi_\Omega u_1) - \nabla \cdot (\kappa \nabla u_1) = 0, \quad (1)$$

$$\partial_\nu u_1|_{\partial B} = g, \quad (2)$$

$$u_1(x, 0)|_\Omega = 0. \quad (3)$$

(can be proven in appropriate Sobolev spaces using Lions Projection Lemma.)

**Inverse Problem:** Given a complete set of measurements

$$\Lambda_1 : g \mapsto u_1|_{\partial B}, \quad u_1 \text{ solves (1)–(3),}$$

reconstruct the interface  $\partial\Omega$  resp. the inclusion  $\Omega$ .

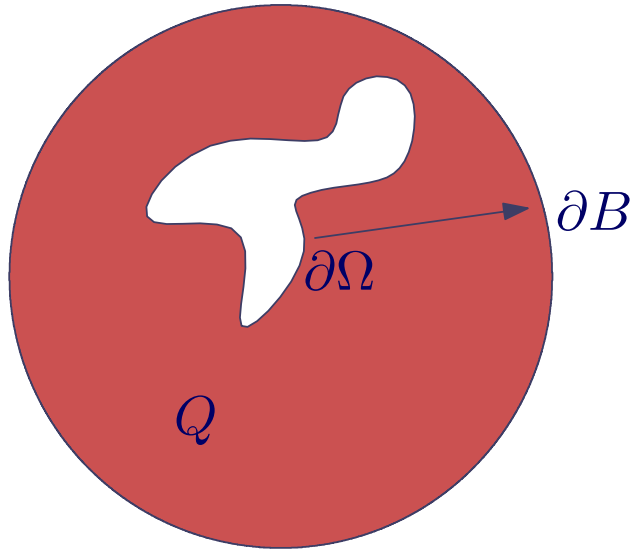
To solve the inverse problem we compare  $\Lambda_1$  to **reference measurements**

$$\Lambda_0 : g \mapsto u_0|_{\partial B}, \quad u_0 \text{ solves } \Delta u_0 = 0, \quad \partial_\nu u_0|_{\partial B} = g,$$

i. e. measurements without an inclusion  $\Omega$ .

**Goal:** Reconstruct  $\Omega$  from given  $\Lambda_0$  and  $\Lambda_1$ .

# Virtual Measurements



$\psi$ : given boundary flux on  $\partial\Omega$

$L : \psi \mapsto v|_{\partial B}$ , where

$$\Delta v(x, t) = 0 \quad \text{in } Q \times ]0, T[, \quad (4)$$

$$\partial_\nu v|_{\partial B} = 0 \quad \text{on } \partial B, \quad (5)$$

$$\partial_\nu v|_{\partial\Omega} = \psi \quad \text{on } \partial\Omega. \quad (6)$$

$\mathcal{R}(L)$  determines  $\Omega$ :

$$v_z|_{\partial B} \in \mathcal{R}(L) \quad \text{if and only if} \quad z \in \Omega$$

where  $v_z$  solves (4) in  $B \setminus \{z\}$ ,  $v_z$  solves (5),  $v_z$  suff. singular in  $z \in B$ ,  
(e. g. a partial derivative of the Green's function for the Laplacian)

# Factorization Method

Key identity of the so-called Factorization Method (for other problems!):

$$\mathcal{R}(L) = \mathcal{R}((\Lambda_0 - \Lambda_1)^{1/2}).$$

$\rightsquigarrow \mathcal{R}(L)$  (and thus  $\Omega$ ) can be computed from the measurements.

Such a range identity

- was originally developed by Kirsch for Inverse Scattering
- is known (under suitable conditions on the inclusion) for
  - Electrostatics (Hähner)
  - EIT (Brühl, Hanke), also with different electrode models (Brühl, Hanke, Hyvönen) and in the half space (Schappel)
  - Diffusion tomography (Kirsch), also with Robin B.C. (Hyvönen)
  - general real elliptic problems (G.)

*Does a similar identity hold in this parabolic-elliptic case?*

# Main Result

● Range inclusions:

$$\mathcal{R}(\tilde{\Lambda}^{1/2}) \subseteq \mathcal{R}(L),$$

$$\mathcal{R}(\tilde{\Lambda}^{1/2}) \supseteq \mathcal{R}(L|_V),$$

$\tilde{\Lambda}$ : symmetric part of  $\Lambda_1 - \Lambda_0$ ,

$V$ : space of boundary fluxes with certain temporal smoothness

● Existence of singular functions  $v_z$  with

$$v_z|_{\partial B} \in \mathcal{R}(L) \quad \text{if and only if} \quad z \in \Omega,$$

and  $\partial_\nu v_z|_{\partial\Omega} \in V$ .

↪

$$z \in \Omega \quad \text{if and only if} \quad v_z|_{\partial B} \in \mathcal{R}(\tilde{\Lambda}^{1/2}).$$

# Sketch of the proof

$$\mathcal{R}(\tilde{\Lambda}^{1/2}) \subseteq \mathcal{R}(L),$$

$$\mathcal{R}(\tilde{\Lambda}^{1/2}) \supseteq \mathcal{R}(L|_V),$$

- Factorization:

$$\tilde{\Lambda} = LFL^*$$

- If  $\|Ax\| \leq \|Bx\|$  for all  $x$  then  $\mathcal{R}(A^*) \subseteq \mathcal{R}(B^*)$ .

$$\rightsquigarrow \mathcal{R}(\tilde{\Lambda}^{1/2}) = \mathcal{R}(LF^{1/2}) \subseteq \mathcal{R}(L).$$

- Coercivity condition for  $F$

$$\rightsquigarrow \mathcal{R}(F^{1/2}) \supseteq H^{\frac{1}{4}}(0, T, H_{\diamond}^{-\frac{1}{2}}(\partial\Omega)) =: V.$$

# Consequences / Remarks

- Theoretical result:

$\partial\Omega$  is uniquely determined by  $\Lambda_1$ ,

i. e. the interface is uniquely determined by measuring all pairs of heat flux and temperature on  $\partial B$ .

- Range test  $v_z|_{\partial B} \in \mathcal{R}(\tilde{\Lambda}^{1/2})$  can be implemented numerically

↪ practical reconstruction algorithm.

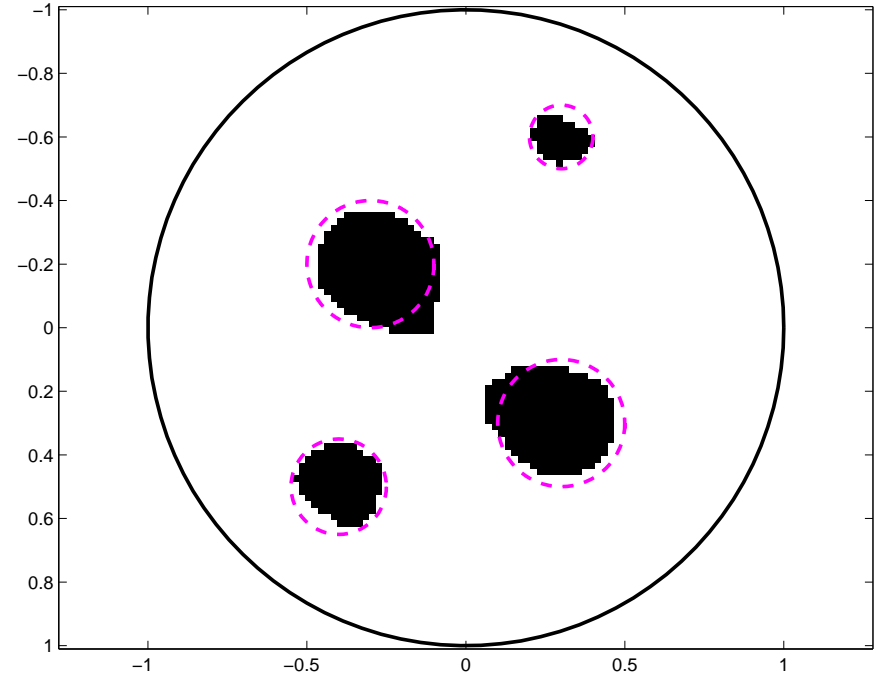
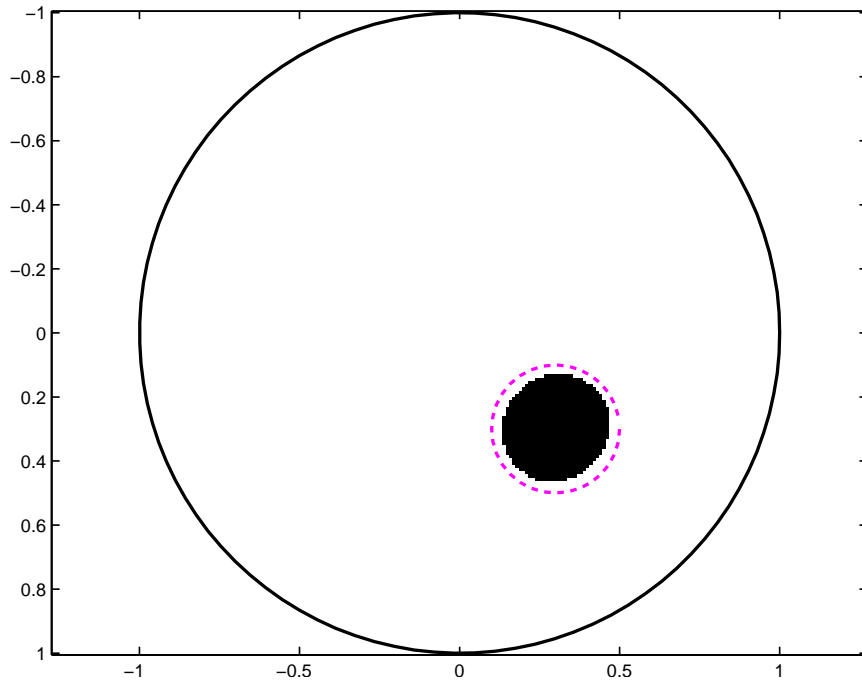
- $\Omega$  does not have to be connected. No a priori information about the number of connected components is needed.

- Usual numerical implementation of range test leads to the problem of determining the convergence of a series from finitely many summands.

↪ threshold problem.

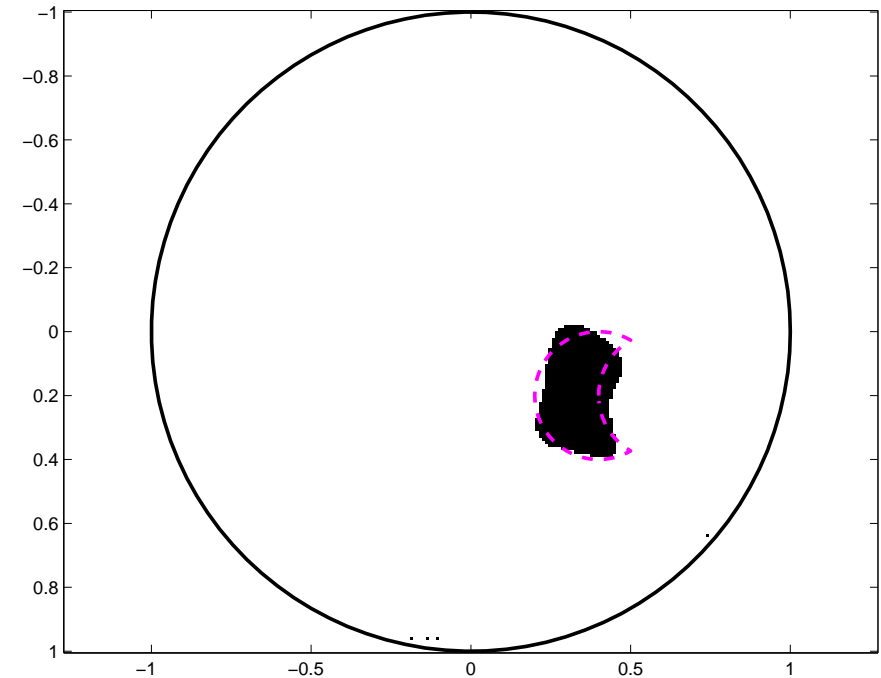
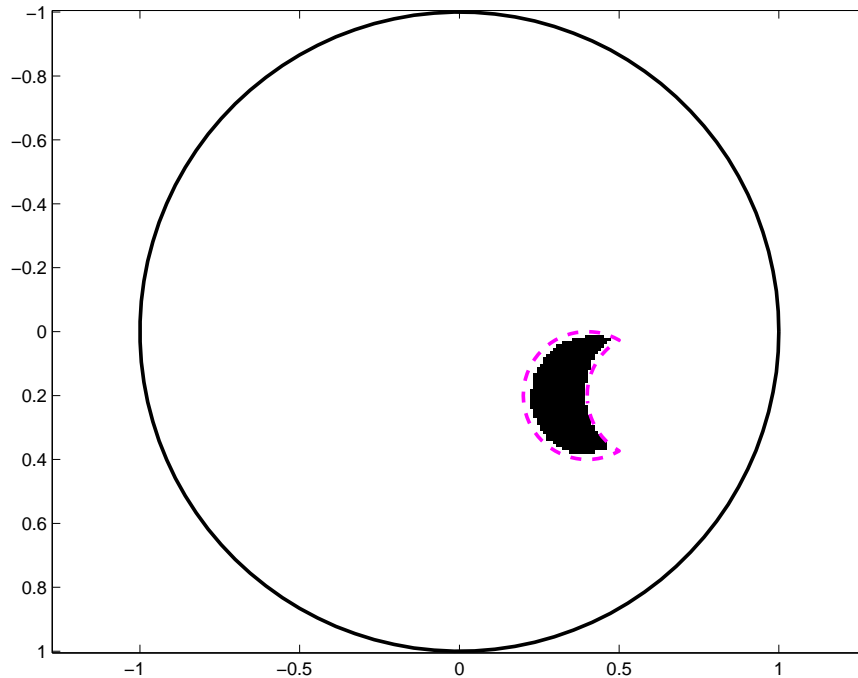


# Numerical results



Reconstruction of single and multiple inclusions.

# Numerical results



Reconstruction of a nonconvex inclusion  
(left: no noise, right: 0.1% noise)