

MONOTONICITY-BASED INVERSION OF THE FRACTIONAL SCHRÖDINGER EQUATION II. GENERAL POTENTIALS AND STABILITY

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Abstract. In this work, we use monotonicity-based methods for the fractional Schrödinger equation with general potentials $q \in L^\infty(\Omega)$ in a Lipschitz bounded open set $\Omega \subset \mathbb{R}^n$ in any dimension $n \in \mathbb{N}$. We demonstrate that if-and-only-if monotonicity relations between potentials and the Dirichlet-to-Neumann map hold up to a finite dimensional subspace. Based on these if-and-only-if monotonicity relations, we derive a constructive global uniqueness result for the fractional Calderón problem and its linearized version. We also derive a reconstruction method for unknown obstacles in a given domain that only requires the background solution of the fractional Schrödinger equation, and we prove uniqueness and Lipschitz stability from finitely many measurements for potentials lying in an a-priori known bounded set in a finite dimensional subset of $L^\infty(\Omega)$.

Key words. Fractional inverse problem, fractional Schrödinger equation, monotonicity, localized potentials, Lipschitz stability, Loewner order

AMS subject classifications. 35R30

1. Introduction . Let Ω be a Lipschitz bounded open set in \mathbb{R}^n , $n \in \mathbb{N}$, and $q \in L^\infty(\Omega)$ be a potential. For $0 < s < 1$, we consider the Dirichlet problem for the nonlocal fractional Schrödinger equation

$$(-\Delta)^s u + qu = 0 \quad \text{in } \Omega, \quad u|_{\Omega_e} = F \quad \text{in } \Omega_e := \mathbb{R}^n \setminus \Omega, \quad (1.1)$$

where the fractional Laplacian $(-\Delta)^s$ is defined by Fourier transform. We will consider the Calderón problem of reconstructing an unknown potential q from the Dirichlet-to-Neumann (DtN) operator

$$\Lambda(q) : H(\Omega_e) \rightarrow H(\Omega_e)^*, \quad F \mapsto (-\Delta)^s u|_{\Omega_e}, \quad \text{where } u \in H^s(\mathbb{R}^n) \text{ solves (1.1),}$$

cf. Section 2 for a precise definition of the DtN-operator and the function spaces, and [32, Section 3] for further properties of the nonlocal DtN map Λ_q .

In the first part of this work [39], we proved an if-and-only-if monotonicity relation between potentials $q \in L_+^\infty(\Omega)$ with positive essential infima and the associated DtN operators $\Lambda(q)$, where the DtN operators are ordered in the sense of definiteness of quadratic forms (also known as Loewner order). From this relation, we obtained a constructive uniqueness result for the Calderón problem and a shape reconstruction method to determine unknown obstacles in a given domain.

The aim of this work is to drop the positivity assumption on the potential q and extend the results from [39] to general potentials $q \in L^\infty(\Omega)$. Note that this may include resonant cases where 0 is a Dirichlet eigenvalue of $(-\Delta)^s + q$ in Ω . In such cases the Dirichlet problem (1.1) is only solvable in a subspace of the natural Dirichlet trace

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space $H(\Omega_e)$ with finite codimension, and the DtN operator $\Lambda(q)$ is defined accordingly, cf. Section 2. For general potentials $q_1, q_2 \in L^\infty(\Omega)$, we will use a combination of monotonicity arguments and localized potentials to show that

$$q_1 \leq q_2 \quad \text{if and only if} \quad \Lambda(q_1) \leq_{\text{fin}} \Lambda(q_2),$$

cf. Theorem 4.1, where $q_1 \leq q_2$ denotes that $q_1(x) \leq q_2(x)$ for almost every (a.e.) $x \in \Omega$, and $\Lambda(q_1) \leq_{\text{fin}} \Lambda(q_2)$ denotes that the quadratic form associated with $\Lambda(q_2) - \Lambda(q_1)$ is non-negative on a subspace of $H(\Omega_e)$ with finite codimension (resp. on a subspace with finite codimension of the intersection of their domains of definition in the case of resonances).

This if-and-only-if monotonicity relation yields a constructive uniqueness proof for the fractional Calderón problem, cf. Theorem 4.3. For non-resonant potentials, we show a similar if-and-only-if monotonicity relation also for the linearized DtN-operators, and deduce uniqueness for the linearized Calderón problem, cf. Theorem 4.8, and Corollary 4.9.

We then turn to the shape reconstruction (or inclusion detection) problem of locating regions where a unknown (non-resonant) coefficient function $q \in L^\infty(\Omega)$ differs from a known (non-resonant) reference function $q_0 \in L^\infty(\Omega)$. We will show that this can be done without solving the fractional Schrödinger equation for potentials other than the reference potentials q_0 . In the indefinite case, with no further assumption on q_0 and q , we characterize the support of $q - q_0$ as the intersection of all closed sets fulfilling a linearized monotonicity condition, cf. Theorem 4.10. In the definite case, that either $q \geq q_0$ or $q_0 \geq q$ in all of Ω , we also obtain an easier characterization of the (inner) support of $q - q_0$ as the union of all open balls fulfilling a linearized monotonicity condition, cf. Theorem 4.11.

Our final result uses monotonicity and localized potentials arguments to show uniqueness and Lipschitz stability for the fractional Calderón problem with finitely many measurements for the case that the potential belongs to an a-priori known bounded set in a finite dimensional subset of $L^\infty(\Omega)$.

Let us give some references of the fast growing body of literature on inverse problems involving the non-local fractional Laplacian operator, and relate our work to previous results. Fractional inverse problems appear when an imaging domain is investigated by an anomalous diffusion process and this process is more complicated than in the standard Brownian motion modeled by the Laplacian $-\Delta$. Global uniqueness for the Calderón problem for the fractional Schrödinger equation was first proven by Ghosh, Salo, and Uhlmann [32], and the recent work of Ghosh, Rüländ, Salo, and Uhlmann [31] shows uniqueness with a single measurement. Note that both results rely on a very strong unique continuation property, and we will utilize this property from [32] as a key ingredient for our results. Furthermore, for uniqueness results, [30] and [61] solved the Calderón problem for general nonlocal variable elliptic operators and the semilinear case, respectively. In addition, [18] studied the fractional Calderón problem with drift, which shows the global uniqueness result holds for drift and potential simultaneously, which is the first example to demonstrate different results between local and nonlocal inverse problems. Recently, [62] investigated the Calderón problem for a space-time fractional parabolic equation. We also refer readers to [16, 17] for further studies on the simultaneous determination of parameters in fractional inverse problems.

Arguments combining PDE-based estimates with blow-up techniques have a long history in the study of inverse coefficients problems, see, e.g., [1, 51, 54, 59, 60]. The

technique of combining monotonicity estimates with localized potentials [29] as used herein is a flexible recent approach that has already lead to a number of results, cf. [6, 8, 15, 33, 34, 35, 39, 40, 44, 45, 46, 47, 49, 72]. Also, several recent works build practical reconstruction methods on monotonicity properties [24, 25, 26, 27, 28, 38, 42, 43, 48, 64, 74, 75, 76, 77, 80]. Notably, the present work shows that monotonicity-based reconstruction methods that have been developed for standard diffusion processes can also be applied to the fractional diffusion case and that the methods even become simpler and more powerful due to the very strong unique continuation property of Ghosh, Salo, and Uhlmann [32]. Moreover, we derive in this work a new result on the existence of simultaneously localized potentials for two coefficient functions, that may be of importance also in the study of other inverse problems.

Logarithmic stability results for the fractional Schrödinger equation and their optimality were proven by Rüländ and Salo in [69, 70]. Lipschitz stability for the finite dimensional fractional Calderón problem with a specific set of finitely many measurements (that depend on the unknown potentials) was shown by Rüländ and Sincich in [71]. Note that our Lipschitz stability result in Section 5 complements the result in [71] as we show that any sufficiently high number of measurements (depending only on the a-priori data but not on the unknown potentials) uniquely determines the potential and that Lipschitz stability holds. Moreover, let us stress that the idea of using monotonicity and localized potentials arguments for proving Lipschitz stability (that was already utilized in [21, 37, 41, 72]), differs from traditional approaches that are mostly based on quantitative unique continuation or quantitative Runge approximation, cf., [2, 3, 4, 5, 7, 9, 10, 11, 12, 13, 14, 19, 52, 53, 56, 57, 58, 65, 71, 73, 78, 79]. Our new approach of showing Lipschitz stability seems conceptually simpler as it does not require quantitative analytic estimates. On the downside, our new approach does not give any analytic bounds on the Lipschitz stability constants that may characterize the asymptotic instability when the dimension of the ansatz space tends to infinity. It may however, lead to a numerical algorithm to calculate the Lipschitz constant for a given setting, cf. [36, 41], which might be important to quantify the achievable resolution and noise robustness in practical applications.

The main technical difficulty in extending the results from the positive potentials case [39] to general coefficients $q \in L^\infty(\Omega)$ is to prove two new extensions of the localized potentials approach [29]. For general potentials, the variational formulation of the fractional Schrödinger equation is no longer coercive but a compact perturbation of a coercive formulation and resonances may arise. To overcome this difficulty, we use an approach that originated in [45] and work in spaces of finite codimension where the formulation is still coercive and resonances are excluded. This makes it necessary to prove that any subspace of finite codimension contains localized potentials. The second major difficulty comes from the fact that only the simpler monotonicity inequality in [39, Lemma 3.1] can be extended to general potentials, cf. Theorem 3.3 in this work. This makes it necessary to prove that localized potentials exist for two different coefficients simultaneously (and in any subspace of finite codimension). It can be expected that the idea of simultaneously localized potentials introduced in this work will also be helpful to extend monotonicity-based methods to other applications.

The paper is structured as follows. In Section 2, we summarize the variational theory for the fractional Schrödinger equation, introduce the DtN operator and the unique continuation property from [32]. In Section 3, we define a generalized Loewner order for linear operators, which holds up to a finite dimensional subspace of a Hilbert

space. We also show that increasing potentials q monotonically increases the corresponding DtN map Λ_q in the sense of this generalized Loewner order, and prove the existence of localized potentials to control the energy terms appearing in the monotonicity relations. The last two sections contain our main results. In Section 4, we investigate a converse result for the monotonicity relations using localized potentials, to deduce if-and-only-if monotonicity relations between the DtN map and the potentials. Based on these results, we prove uniqueness for the fractional Calderón problem in a constructive way. We also prove uniqueness for the linearized fractional Calderón problem and develop an inclusion detection algorithm based on monotonicity tests. Finally, in Section 5, we use the monotonicity relations and the localized potentials, to prove uniqueness and Lipschitz stability in finite dimensional subspaces by finitely many measurements.

2. The fractional Schrödinger equation for general potentials . Throughout this work let $s \in (0, 1)$, $n \in \mathbb{N}$, $\Omega \subseteq \mathbb{R}^n$ be a Lipschitz bounded open set, and $q \in L^\infty(\Omega)$. All function spaces in this work are real-valued unless indicated otherwise. In this section, we briefly summarize some notations and results on the fractional Schrödinger equation and the associated Dirichlet problem.

2.1. Variational formulation of the fractional Schrödinger equation. As in [39] we consider the fractional Laplacian (defined by Fourier transform) as an operator

$$(-\Delta)^s : L^2(\mathbb{R}^n) \rightarrow \mathcal{S}'(\mathbb{R}^n),$$

The fractional Sobolev space is defined by

$$H^s(\mathbb{R}^n) := \{u \in L^2(\mathbb{R}^n) : (-\Delta)^{s/2}u \in L^2(\mathbb{R}^n)\}$$

and equipped with the scalar product

$$(u, v)_{H^s(\mathbb{R}^n)} := \int_{\mathbb{R}^n} \left((-\Delta)^{s/2}u \cdot (-\Delta)^{s/2}v + uv \right) dx \quad \text{for all } u, v \in H^s(\mathbb{R}^n).$$

It can be shown that $H^s(\mathbb{R}^n)$ is a Hilbert space, cf., e.g., [20]. Let

$$H_0^s(\Omega) := \text{closure of } C_c^\infty(\Omega) \text{ in } H^s(\mathbb{R}^n),$$

and note that this space is sometimes denoted as $\tilde{H}^s(\Omega)$ in the literature, e.g., [32, 30].

We also define the bilinear form

$$\mathcal{B}_q(u, w) := \int_{\mathbb{R}^n} (-\Delta)^{s/2}u \cdot (-\Delta)^{s/2}w dx + \int_{\Omega} quw dx \quad \text{for } u, w \in H^s(\mathbb{R}^n).$$

Then, for any $f \in L^2(\Omega)$, $u \in H^s(\mathbb{R}^n)$ solves (in the sense of distributions)

$$(-\Delta)^s u + qu = f \quad \text{in } \Omega$$

if and only if $u \in H^s(\mathbb{R}^n)$ fulfills the variational formulation

$$\mathcal{B}_q(u, w) = \int_{\Omega} fw dx \quad \text{for all } w \in H_0^s(\Omega), \quad (2.1)$$

cf., e.g., [39, Lemma 2.1].

2.2. The Dirichlet boundary value problem. The Dirichlet trace operator on $\Omega_e := \mathbb{R}^n \setminus \overline{\Omega}$ can be defined using abstract quotient spaces by setting

$$\gamma_{\Omega_e}^{(D)} : H^s(\mathbb{R}^n) \rightarrow H(\Omega_e) := H^s(\mathbb{R}^n)/H_0^s(\Omega), \quad u \mapsto u + H_0^s(\Omega).$$

Then, by definition, $\gamma_{\Omega_e}^{(D)}$ is surjective, $H_0^s(\Omega) = \{u \in H^s(\mathbb{R}^n) : \gamma_{\Omega_e}^{(D)}u = 0\}$. Moreover, for all $u, v \in H^s(\mathbb{R}^n)$,

$$\gamma_{\Omega_e}^{(D)}u = \gamma_{\Omega_e}^{(D)}v \quad \text{implies that} \quad u(x) = v(x) \quad \text{for } x \in \Omega_e \text{ a.e.}, \quad (2.2)$$

cf., e.g., [39, Lemma 2.2]. This implies that $\gamma_{\Omega_e}^{(D)}$ is an injective mapping from $C_c^\infty(\Omega_e)$ into $H(\Omega_e)$. For the sake of readability we will write $u|_{\Omega_e}$ instead of $\gamma_{\Omega_e}^{(D)}u$ throughout this work, and identify $C_c^\infty(\Omega_e)$ with its image in $H(\Omega_e)$.

Throughout this work, we will use that for all $u, w \in H_0^s(\Omega)$

$$\mathcal{B}_q(u, w) = ((I - \iota^* \iota + \iota^* M_q \iota)u, w)_{H_0^s(\Omega)}$$

with the bounded linear operators

$$\begin{aligned} I &: H_0^s(\Omega) \rightarrow H_0^s(\Omega), \\ \iota &: H_0^s(\Omega) \rightarrow L^2(\Omega), \\ M_q &: L^2(\Omega) \rightarrow L^2(\Omega), \end{aligned}$$

denoting the identity operator, the compact restriction and embedding, cf. [66, Lemma 10], and the multiplication operator by q .

We then have the following result on the solvability of the Dirichlet boundary value problem.

LEMMA 2.1. *Let $F \in H(\Omega_e)$, $f \in L^2(\Omega)$, and*

$$N_q := \{u \in H_0^s(\Omega) : (-\Delta)^s u + qu = 0 \quad \text{in } \Omega\}.$$

(a) *$u \in H^s(\mathbb{R}^n)$ solves the Dirichlet problem*

$$(-\Delta)^s u + qu = f \quad \text{in } \Omega, \quad u|_{\Omega_e} = F, \quad (2.3)$$

if and only if $u = u^{(0)} + u^{(F)}$, where $u^{(F)} \in H^s(\mathbb{R}^n)$ fulfills $u^{(F)}|_{\Omega_e} = F$, and $u^{(0)} \in H_0^s(\Omega)$ solves

$$\mathcal{B}_q(u^{(0)}, w) = -\mathcal{B}_q(u^{(F)}, w) + \int_{\Omega} f w \, dx \quad \text{for all } w \in H_0^s(\Omega).$$

Note that for $F \in C_c^\infty(\Omega_e)$ one can simply choose $u^{(F)} := F$.

(b) *N_q is finite-dimensional. The Dirichlet problem (2.3) is solvable if and only if*

$$\mathcal{B}_q(u^{(F)}, w) = \int_{\Omega} f w \, dx \quad \text{for all } w \in N_q. \quad (2.4)$$

The solution $u \in H^s(\mathbb{R}^n)$ of (2.3) is unique up to addition of a function in N_q , and $u + N_q \in H^s(\mathbb{R}^n)/N_q$ depends linearly and continuously on $F \in H(\Omega_e)$ and $f \in L^2(\Omega)$.

Proof. (a) immediately follows from the variational formulation (2.1).

To prove (b), we use the Riesz representation theorem to obtain $v_f^F \in H_0^s(\Omega)$ fulfilling

$$(v_f^F, w)_{H_0^s(\Omega)} = -\mathcal{B}_q(u^{(F)}, w) + \int_{\Omega} f w \, dx \quad \text{for all } w \in H_0^s(\Omega).$$

Using (a), and that $w \in H_0^s(\Omega)$ implies $w(x) = 0$ for $x \in \Omega_e$ a.e., we obtain that $u \in H^s(\mathbb{R}^n)$ solves (2.3) if and only if $u = u^{(0)} + u^{(F)}$ with $u^{(0)} \in H_0^s(\Omega)$ solving

$$\begin{aligned} & \left((I - \iota^* \iota + \iota^* M_q \iota) u^{(0)}, w \right)_{H_0^s(\Omega)} \\ &= \mathcal{B}_q(u^{(0)}, w) = -\mathcal{B}_q(u^{(F)}, w) + \int_{\Omega} f w \, dx = (v_f^F, w)_{H_0^s(\Omega)} \quad \text{for all } w \in H_0^s(\Omega), \end{aligned}$$

i.e.

$$(I - \iota^* \iota + \iota^* M_q \iota) u^{(0)} = v_f^F,$$

and that

$$N_q = \mathcal{N}(I - \iota^* \iota + \iota^* M_q \iota). \quad (2.5)$$

Here $\mathcal{N}(A)$ stands for the kernel of the linear operator A . Since $\iota^* \iota - \iota^* M_q \iota$ is compact and self-adjoint, Fredholm theory (cf., e.g., [22, Appendix D, Theorem 5]) yields that N_q is finite-dimensional, and that (2.3) is solvable if and only if

$$(v_f^F, w)_{H_0^s(\Omega)} = 0 \quad \text{for all } w \in \mathcal{N}(I - \iota^* \iota + \iota^* M_q \iota) = N_q,$$

which gives the condition (2.4).

Clearly $u^{(0)}$ is unique up to addition of a function in N_q , and $u^{(0)} + N_q$ depends linearly and continuously on $v_f^F \in H_0^s(\Omega)$. It easily follows that $u = u^{(0)} + u^{(F)}$ is unique up to addition of a function in N_q , and that $u + N_q \in H^s(\mathbb{R}^n)/N_q$ depends linearly and continuously on $F \in H(\Omega_e)$ and $f \in L^2(\Omega)$. \square

COROLLARY 2.2. *Let $H_q^s(\mathbb{R}^n) \subseteq H^s(\mathbb{R}^n)$ be the $H^s(\mathbb{R}^n)$ -orthogonal complement of N_q , and*

$$H_q(\Omega_e) := \{F \in H(\Omega_e) : \mathcal{B}_q(u^{(F)}, w) = 0 \quad \text{for all } w \in N_q\}.$$

Then the codimension of $H_q(\Omega_e)$ in $H(\Omega_e)$ is at most $\dim N_q$, and for all $F \in H_q(\Omega_e)$ there exists a unique solution $u \in H_q^s(\mathbb{R}^n)$ of the Dirichlet problem

$$(-\Delta)^s u + qu = 0 \quad \text{in } \Omega, \quad u|_{\Omega_e} = F, \quad (2.6)$$

and that the solution operator

$$S_q : H_q(\Omega_e) \rightarrow H_q^s(\mathbb{R}^n), \quad F \mapsto u, \quad \text{where } u \text{ solves (2.6),}$$

is linear and bounded.

Proof. We first show that $H_q(\Omega_e)$ is well-defined. If $u^{(F)}, \tilde{u}^{(F)} \in H^s(\mathbb{R}^n)$ both fulfill $u^{(F)}|_{\Omega_e} = F = \tilde{u}^{(F)}|_{\Omega_e}$, then $u^{(F)} - \tilde{u}^{(F)}|_{\Omega_e} \in H_0^s(\Omega)$ and thus it follows from the definition of N_q (2.5) and (2.1) that

$$\mathcal{B}_q(u^{(F)} - \tilde{u}^{(F)}, w) = 0 \quad \text{for all } w \in N_q.$$

Next, we show that the codimension of $H_q(\Omega_e)$ in $H(\Omega_e)$ is at most $d := \dim N_q$. Let $(w_1, \dots, w_d) \subset N_q$ be an orthonormal basis of N_q , and let $\gamma^- : H(\Omega_e) \rightarrow H^s(\mathbb{R}^n)$ be a linear right inverse of the Dirichlet trace operator $\gamma_{\Omega_e}^{(D)}$. Then, by linearity,

$$H_q(\Omega_e) = \{F \in H(\Omega_e) : \mathcal{B}_q(\gamma^- F, w_j) = 0 \quad \text{for all } j = 1, \dots, d\} = \mathcal{N}(\mathcal{A}),$$

with a linear operator

$$\mathcal{A} : H(\Omega_e) \rightarrow \mathbb{R}^d, \quad F \mapsto (\mathcal{B}_q(\gamma^- F, w_j))_{j=1, \dots, d}.$$

Hence, the codimension of $H_q(\Omega_e) = \mathcal{N}(\mathcal{A})$ is $\dim(\mathcal{R}(\mathcal{A})) \leq d$.

Finally, it follows from Lemma 2.1(b) that (2.6) possesses a solution $\tilde{u} \in H^s(\mathbb{R}^n)$ which is unique up to addition of a function in N_q . Hence,

$$u := \tilde{u} - \sum_{j=1}^d w_j(\tilde{u}, w_j)_{H^s(\mathbb{R}^n)} \in H_q^s(\mathbb{R}^n)$$

solves (2.6), and $H_q^s(\mathbb{R}^n)$ contains no other solutions of (2.6). Since $H_q^s(\mathbb{R}^n)$ is isomorphic to $H^s(\mathbb{R}^n)/N_q$, the continuity and linearity of the solution operator S_q also follow from Lemma 2.1(b). \square

2.3. Neumann traces and the Dirichlet-to-Neumann operator. We define the Neumann trace operator

$$\gamma_{\Omega_e}^{(N)} : H_{\Delta}^s(\mathbb{R}^n) := \{u \in H^s(\mathbb{R}^n) : \exists f \in L^2(\Omega) \text{ with } (-\Delta)^s u = f \text{ in } \Omega\} \rightarrow H(\Omega_e)^*$$

by setting

$$\left\langle \gamma_{\Omega_e}^{(N)} u, F \right\rangle := \int_{\mathbb{R}^n} (-\Delta)^{s/2} u \cdot (-\Delta)^{s/2} v^{(F)} dx - \int_{\Omega} (-\Delta)^s u \cdot v^{(F)} dx, \quad (2.7)$$

where $v^{(F)} \in H^s(\mathbb{R}^n)$ fulfills $v^{(F)}|_{\Omega_e} = F$, $H(\Omega_e)^*$ is the dual space of $H(\Omega_e)$, and throughout this paper $\langle \cdot, \cdot \rangle$ denotes the dual pairing on $H(\Omega_e)^* \times H(\Omega_e)$. Note that $\gamma_{\Omega_e}^{(N)} u$ is well-defined since the right hand side of (2.7) does not depend on the choice of $v^{(F)}$, and that $\gamma_{\Omega_e}^{(N)}$ is a bounded linear operator.

For the sake of readability, we also use the formal notation $(-\Delta)^s u|_{\Omega_e} := \gamma_{\Omega_e}^{(N)} u$ for the Neumann trace, which can be motivated by the following lemma, see also [39, Remark 2.4] and [32] for further justifications of this notation under additional smoothness conditions on u or Ω .

LEMMA 2.3. *Let $u \in H_{\Delta}^s(\mathbb{R}^n)$. If $\gamma_{\Omega_e}^{(N)} u \in L^2(\Omega)$ in the sense that there exists $g \in L^2(\Omega_e)$ with*

$$\left\langle \gamma_{\Omega_e}^{(N)} u, F \right\rangle = \int_{\Omega_e} g v^{(F)} dx \quad \text{for all } v^{(F)} \in H^s(\mathbb{R}^n) \text{ with } v^{(F)}|_{\Omega_e} = F,$$

then $g = (-\Delta)^s u$ in Ω_e (in the sense of distributions).

Proof. For all $\varphi \in C_c^\infty(\Omega_e) \subseteq H(\Omega_e)$ (cf. subsection 2.2), we have that

$$\begin{aligned} \int_{\Omega} g \varphi dx &= \left\langle \gamma_{\Omega_e}^{(N)} u, G \right\rangle = \int_{\mathbb{R}^n} (-\Delta)^{s/2} u \cdot (-\Delta)^{s/2} \varphi dx - \int_{\Omega} (-\Delta)^s u \cdot \varphi dx \\ &= \int_{\mathbb{R}^n} (-\Delta)^{s/2} u \cdot (-\Delta)^{s/2} \varphi dx = \langle (-\Delta)^s u, \varphi \rangle_{\mathcal{D}'(\Omega_e) \times \mathcal{D}(\Omega_e)}. \end{aligned}$$

□

Note also that if $u \in H^s(\mathbb{R}^n)$ solves $(-\Delta)^s u + qu = 0$ in Ω , then

$$\langle (-\Delta)^s u|_{\Omega_e}, G \rangle = \mathcal{B}_q(u, v^{(G)})$$

holds for all $G \in H(\Omega_e)$ and all $v^{(G)} \in H^s(\mathbb{R}^n)$ with $v^{(G)}|_{\Omega_e} = G$. Using Corollary 2.2, we can thus define the linear bounded DtN operator

$$\Lambda(q) : H_q(\Omega_e) \rightarrow H(\Omega_e)^*, \quad F \mapsto (-\Delta)^s u|_{\Omega_e}$$

where $u \in H_q^s(\mathbb{R}^n)$ solves

$$(-\Delta)^s u + qu = 0 \quad \text{in } \Omega, \quad u|_{\Omega_e} = F.$$

In view of the following sections, note that for $q_1, q_2 \in L^\infty(\Omega)$,

$$H_{q_1, q_2}(\Omega_e) := H_{q_1}(\Omega_e) \cap H_{q_2}(\Omega_e)$$

is a subspace of $H(\Omega_e)$ with codimension less than or equal to $\dim N_{q_1} + \dim N_{q_2}$, on which both $\Lambda(q_1)$ and $\Lambda(q_2)$ are defined. Hence, throughout this work, $\Lambda(q_1) - \Lambda(q_2)$ will always denote the linear bounded operator

$$\Lambda(q_1) - \Lambda(q_2) : H_{q_1, q_2}(\Omega_e) \rightarrow H(\Omega_e)^*.$$

The following relation between the DtN operator and the bilinear form will be useful.

LEMMA 2.4. *Let $q_1, q_2 \in L^\infty(\Omega)$, $F \in H_{q_1}(\Omega_e)$, $G \in H_{q_2}(\Omega_e)$, and let $u \in H_{q_1}^s(\mathbb{R}^n)$, $v \in H_{q_2}^s(\mathbb{R}^n)$ solve*

$$\begin{aligned} (-\Delta)^s u + q_1 u &= 0 & \text{in } \Omega, & \quad u|_{\Omega_e} = F, \\ (-\Delta)^s v + q_2 v &= 0 & \text{in } \Omega, & \quad v|_{\Omega_e} = G. \end{aligned}$$

Then

$$\langle \Lambda(q_1)F, F \rangle = \mathcal{B}_{q_1}(u, u) \quad \text{and} \quad \langle \Lambda(q_1)F, G \rangle = \mathcal{B}_{q_1}(u, v),$$

and under the additional restriction that $F, G \in H_{q_1, q_2}(\Omega_e)$ this also implies that

$$\langle (\Lambda(q_1) - \Lambda(q_2))F, G \rangle = \mathcal{B}_{q_1}(u, v) - \mathcal{B}_{q_2}(u, v) = \int_{\Omega} (q_1 - q_2)uv \, dx.$$

Proof. This immediately follows from the variational formulation in Lemma 2.1 and the definition of the Neumann trace. □

2.4. Unique continuation from open sets and Cauchy data. We recall the unique continuation result from Ghosh, Salo and Uhlmann [32]:

THEOREM 2.5. [32, Theorem 1.2] *Let $n \in \mathbb{N}$, and $0 < s < 1$. If $u \in H^r(\mathbb{R}^n)$ for some $r \in \mathbb{R}$, and both u and $(-\Delta)^s u$ vanish in the same arbitrary non-empty open set in \mathbb{R}^n , then $u \equiv 0$ in \mathbb{R}^n .*

We will make use of the following simple corollary.

COROLLARY 2.6. *Let $u \in H^s(\mathbb{R}^n)$ solve $(-\Delta)^s u + qu = f$ in Ω , with $f \in L^2(\Omega)$*

- (a) If u and f vanish in the same nonempty open set $\mathcal{O} \subset \Omega$, then $u \equiv 0$ in \mathbb{R}^n .
(b) If $u|_{\Omega_e} = 0$ and $(-\Delta)^s u|_{\Omega_e} = 0$, then $u \equiv 0$ in \mathbb{R}^n .

Proof. (a) follows since $u = 0$ in \mathcal{O} , and $(-\Delta)^s u + qu = 0$ in \mathcal{O} , implies $(-\Delta)^s u = 0$ in \mathcal{O} . For (b) note that $u|_{\Omega_e}$ and $(-\Delta)^s u|_{\Omega_e}$ are only formal notations for the Dirichlet and Neumann traces of u , but $u|_{\Omega_e} = 0$, and $(-\Delta)^s u|_{\Omega_e} = 0$ do imply that

$$u = 0 \quad \text{in } \Omega_e, \quad \text{and} \quad (-\Delta)^s u = 0 \quad \text{in } \Omega_e$$

in the sense of distributions by (2.2) and Lemma 2.3. Hence, both cases follow from Theorem 2.5. \square

REMARK 2.7. When $\frac{1}{4} \leq s < 1$, then the unique continuation property in Corollary 2.6(a) already holds under the weaker condition that u vanishes in a subset of Ω with positive measure, cf. [31, Proposition 5.1]. Moreover, based on such property, [31] shows global uniqueness for the fractional Schrödinger equation by a single measurement.

3. Monotonicity relations and localized potentials. In this section we derive monotonicity relations between $L^\infty(\Omega)$ potentials and their associated DtN operators, and show how to control the energy terms in the monotonicity relations with the technique of localized potentials.

3.1. Monotonicity relations. We characterize the monotonicity relations between DtN operators with an extended *Loewner order* that holds up to finite dimensional subspaces.

DEFINITION 3.1. Let H be a Hilbert space and $H_1, H_2 \subseteq H$ be two subspaces of finite codimension, and let $L_1 : H_1 \rightarrow H$, $L_2 : H_2 \rightarrow H$ be two linear bounded operators. For a number $d \in \mathbb{N}_0 := \mathbb{N} \cup \{0\}$ we write

$$L_1 \leq_d L_2$$

if there exists a subspace $W \subseteq H_{12} := H_1 \cap H_2$ with $\dim(W) \leq d$, and

$$\langle (L_2 - L_1)v, v \rangle \geq 0 \quad \text{for all } v \in W^\perp \subseteq H_{12}.$$

Here and in the following, we use the notation $W^\perp \subseteq H_{12}$ to indicate that the orthogonal complement is taken in H_{12} .

We write $L_1 \leq L_2$ if $L_1 \leq_0 L_2$, and $L_1 \leq_{fn} L_2$ if $L_1 \leq_d L_2$ for some $d \in \mathbb{N}_0$. We also write

$$L_1 \stackrel{fn}{=} L_2 \quad \text{if} \quad L_1 \leq_{fn} L_2, \quad \text{and} \quad L_2 \leq_{fn} L_1,$$

i.e. if there exists a finite dimensional subspace $W \subseteq H_{12}$ so that

$$\langle (L_2 - L_1)v, v \rangle = 0 \quad \text{for all } v \in W^\perp \subseteq H_{12}.$$

Note that if $H_1 = H_2 = H$ and L_1, L_2 are self-adjoint and compact, this is the same extended Loewner order as in [45].

Let us stress that the binary relation \leq_d is reflexive, but generally neither transitive, nor antisymmetric. Obviously, $L_1 \leq_{d_1} L_2$ and $L_2 \leq_{d_2} L_3$ imply that $L_1 \leq_d L_3$,

with $d = d_1 + d_2 + \text{codim}(H_2)$, so that \leq_{fin} is a reflexive and transitive relation, i.e., a preorder. Moreover, Corollaries 4.2 and 4.9 will show that \leq_{fin} is antisymmetric on the set of NtD operators and on their linearizations around a fixed non-resonant potential, so that on these sets, \leq_{fin} is a partial order.

For two potentials $q_1, q_2 \in L^\infty(\Omega)$ we write $q_1 \leq q_2$ if $q_1(x) \leq q_2(x)$ for almost everywhere (a.e.) $x \in \Omega$. We will show that increasing the potential q in this sense increases the DtN map $\Lambda(q)$ in the sense of the generalized Loewner order in Definition 3.1. Note that monotonicity relations in inverse coefficient problems go back to the works of Ikehata, Kang, Seo, and Sheen [50, 55], and they have been at the core of many reconstruction algorithms including the Factorization method and the Monotonicity method, cf. the list of references in the introduction. Extensions of monotonicity relations to subspaces of finite codimensions have first been studied in [45, 33], and we follow the general approach from there. A sharper bound on the dimension of the excluded subspaces has recently been obtained for the standard Helmholtz equation in [44].

DEFINITION 3.2. *For $q \in L^\infty(\Omega)$ let $d(q) \in \mathbb{N}_0$ denote the number of eigenvalues (counted with multiplicity) of the compact self-adjoint operator $\iota^* \iota - \iota^* M_q \iota$ that are greater than 1.*

THEOREM 3.3 (Monotonicity relations). *Let $q_1, q_2 \in L^\infty(\Omega)$. There exists a subspace $V \subseteq H_{q_1, q_2}(\Omega_e)$ with $\dim(V) \leq d(q_2)$ so that*

$$\langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \geq \int_{\Omega} (q_1 - q_2) |u_1|^2 dx \quad \text{for all } F \in V^\perp \subseteq H_{q_1, q_2}(\Omega_e), \quad (3.1)$$

where $u_1 \in H_{q_1}^s(\mathbb{R}^n)$ solves $(-\Delta)^s u_1 + q_1 u_1 = 0$ in Ω with $u_1|_{\Omega_e} = F$.

Hence

$$q_1 \geq q_2 \text{ a.e. in } \Omega \quad \text{implies that} \quad \Lambda(q_1) \geq_{d(q_2)} \Lambda(q_2).$$

Before we prove Theorem 3.3, let us also formulate a variant that will be useful for applying the idea of localized potentials in the next sections, remark on interchanging q_1 and q_2 , and discuss the dependence of $\dim(N_q)$ and $d(q)$ on q .

THEOREM 3.4. *Let $q_1, q_2 \in L^\infty(\Omega)$. There exists a subspace*

$$V_+ \subseteq H_{q_1, q_2}(\Omega_e) \quad \text{with} \quad \dim(V_+) \leq d(q_2) + \dim(N_{q_2}),$$

and a constant $\lambda > 0$, so that for all $F \in V_+^\perp \subseteq H_{q_1, q_2}(\Omega_e)$

$$\langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \geq \int_{\Omega} (q_1 - q_2) |u_1|^2 dx + \lambda \|u_1 - u_2\|_{H^s(\mathbb{R}^n)}^2 \quad (3.2)$$

and, for all $D \subseteq \Omega$ containing $\text{supp}(q_1 - q_2)$,

$$\|u_2\|_{L^2(D)} \leq c \|u_1\|_{L^2(D)}, \quad (3.3)$$

where $c := 1 + \frac{1}{\lambda} \|q_1 - q_2\|_{L^\infty(D)}$, and, for $j = 1, 2$, $u_j \in H_{q_j}^s(\mathbb{R}^n)$ solve

$$(-\Delta)^s u_j + q_j u_j = 0 \quad \text{in } \Omega, \quad u_j|_{\Omega_e} = F.$$

REMARK 3.5. *By interchanging q_1 and q_2 in Theorems 3.3 and 3.4, we also obtain that there exist subspaces*

$$V, V_+ \subseteq H_{q_1, q_2}(\Omega_e) \quad \text{with} \quad \dim(V) \leq d(q_1), \quad \text{and} \quad \dim(V_+) \leq d(q_1) + \dim(N_{q_1}),$$

and a constant $\lambda > 0$, so that

$$\langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \leq \int_{\Omega} (q_1 - q_2) |u_2|^2 dx \quad \text{for all } F \in V^\perp \subseteq H_{q_1, q_2}(\Omega_e),$$

and

$$\begin{aligned} \langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle &\leq \int_{\Omega} (q_1 - q_2) |u_2|^2 dx - \lambda \|u_1 - u_2\|_{H^s(\mathbb{R}^n)}^2, \\ \|u_1\|_{L^2(D)} &\leq c \|u_2\|_{L^2(D)}, \end{aligned}$$

for all $D \supseteq \text{supp}(q_1 - q_2)$, and all $F \in V_+^\perp \subseteq H_{q_1, q_2}(\Omega_e)$, where $c := 1 + \frac{1}{\lambda} \|q_1 - q_2\|_{L^\infty(D)}$, $u_1 = S_{q_1}(F)$, and $u_2 = S_{q_2}(F)$.

Combining Theorem 3.3 with its interchanged version, we obtain a subspace

$$V \subseteq H_{q_1, q_2}(\Omega_e) \quad \text{with} \quad \dim(V) \leq d(q_1) + d(q_2),$$

so that

$$\int_{\Omega} (q_1 - q_2) |u_1|^2 dx \leq \langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \leq \int_{\Omega} (q_1 - q_2) |u_2|^2 dx.$$

for all $F \in V^\perp \subseteq H_{q_1, q_2}(\Omega_e)$, $u_1 = S_{q_1}(F)$, and $u_2 = S_{q_2}(F)$.

Combining Theorem 3.4 with its interchanged version, we obtain a subspace

$$V_+ \subseteq H_{q_1, q_2}(\Omega_e) \quad \text{with} \quad \dim(V_+) \leq d(q_1) + d(q_2) + \dim(N_{q_1}) + \dim(N_{q_2}),$$

and constants $\lambda, c_1, c_2 > 0$, so that

$$\begin{aligned} \int_{\Omega} (q_1 - q_2) |u_1|^2 dx + \lambda \|u_1 - u_2\|_{H^s(\mathbb{R}^n)}^2 &\leq \langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \\ &\leq \int_{\Omega} (q_1 - q_2) |u_2|^2 dx - \lambda \|u_1 - u_2\|_{H^s(\mathbb{R}^n)}^2, \end{aligned}$$

and

$$c_1 \|u_1\|_{L^2(D)} \leq \|u_2\|_{L^2(D)} \leq c_2 \|u_1\|_{L^2(D)}$$

for all $D \supseteq \text{supp}(q_1 - q_2)$, and all $F \in V_+^\perp \subseteq H_{q_1, q_2}(\Omega_e)$, $u_1 = S_{q_1}(F)$, and $u_2 = S_{q_2}(F)$.

THEOREM 3.6. *Let $d(q)$ be given by Definition 3.2 and N_q be defined by (2.5).*

(a) *For $q_1, q_2 \in L^\infty(\Omega)$*

$$q_1 \leq q_2 \quad \text{implies} \quad d(q_1) \geq d(q_2).$$

(b) *For all $q_1 \in L^\infty(\Omega)$ there exists $\epsilon > 0$ so that*

$$\dim(N_{q_1}) \geq \dim(N_{q_2}) \quad \text{for all } q_2 \in L^\infty(\Omega) \text{ with } \|q_2 - q_1\|_{L^\infty(\Omega)} \leq \epsilon.$$

To prove Theorems 3.3, 3.4, and 3.6, we first show the following lemmas.

LEMMA 3.7. *Let $q_1, q_2 \in L^\infty(\Omega)$. Then, for all $F \in H_{q_1, q_2}(\Omega_e)$,*

$$\langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle + \int_{\Omega} (q_2 - q_1) |u_1|^2 dx = \mathcal{B}_{q_2}(u_2 - u_1, u_2 - u_1),$$

where $u_1 = S_{q_1}(F)$, and $u_2 = S_{q_2}(F)$.

Proof. Using lemma 2.4, the assertion follows from

$$\begin{aligned} \mathcal{B}_{q_2}(u_2 - u_1, u_2 - u_1) &= \mathcal{B}_{q_2}(u_2, u_2) - 2\mathcal{B}_{q_2}(u_2, u_1) + \mathcal{B}_{q_2}(u_1, u_1) \\ &= -\mathcal{B}_{q_2}(u_2, u_1) + \mathcal{B}_{q_2}(u_1, u_1) = -\mathcal{B}_{q_2}(u_2, u_1) + \mathcal{B}_{q_1}(u_1, u_1) + \int_{\Omega} (q_2 - q_1) |u_1|^2 dx \\ &= \langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle + \int_{\Omega} (q_2 - q_1) |u_1|^2 dx. \end{aligned}$$

□

LEMMA 3.8. *Let $q \in L^\infty(\Omega)$. Then there exists a subspace $W \subseteq H_0^s(\Omega)$ with $\dim(W) = d(q)$, and a constant $\lambda > 0$, so that*

$$\begin{aligned} \mathcal{B}_q(w, w) &\geq 0 && \text{for all } w \in W^\perp \subseteq H_0^s(\Omega), \text{ and} \\ \mathcal{B}_q(w, w) &\geq \lambda \|w\|_{H^s(\mathbb{R}^n)}^2 && \text{for all } w \in (W + N_q)^\perp \subseteq H_0^s(\Omega). \end{aligned}$$

Proof. Let W be the sum of eigenspaces of the compact self-adjoint operator $\iota^* \iota - \iota^* M_q \iota$ corresponding to eigenvalues larger than 1. Then

$$\mathcal{B}_q(w, w) = ((I - \iota^* \iota + \iota^* M_q \iota)w, w)_{H^s(\mathbb{R}^n)} \geq 0 \text{ for all } w \in W^\perp \subseteq H_0^s(\Omega).$$

Since $N_q = \mathcal{N}(I - \iota^* \iota + \iota^* M_q \iota)$ is the eigenspace of $\iota^* \iota - \iota^* M_q \iota$ corresponding to the eigenvalue 1, it also follows that

$$\mathcal{B}_q(w, w) \geq (1 - \mu) \|w\|_{H^s(\mathbb{R}^n)}^2 \text{ for all } w \in (W + N_q)^\perp \subseteq H_0^s(\Omega),$$

where μ is the largest eigenvalue of $\iota^* \iota - \iota^* M_q \iota$ smaller than 1. Hence, the assertion follows with $\lambda := 1 - \mu$. □

LEMMA 3.9. *Let $q_1, q_2 \in L^\infty(\Omega)$. There exists $\lambda > 0$ and subspaces*

$$V \subseteq V_+ \subseteq H_{q_1, q_2}(\Omega_e) \quad \text{with} \quad \dim(V) \leq d(q_2), \quad \dim(V_+) \leq d(q_2) + \dim(N_{q_2}),$$

so that

$$\mathcal{B}_{q_2}(u_2 - u_1, u_2 - u_1) \geq 0 \quad \text{for all } F \in V^\perp \subseteq H_{q_1, q_2}(\Omega_e), \quad (3.4)$$

$$\mathcal{B}_{q_2}(u_2 - u_1, u_2 - u_1) \geq \lambda \|u_2 - u_1\|_{H^s(\mathbb{R}^n)}^2 \quad \text{for all } F \in V_+^\perp \subseteq H_{q_1, q_2}(\Omega_e), \quad (3.5)$$

where $u_1 = S_{q_1}(F)$, and $u_2 = S_{q_2}(F)$.

Proof. The difference of the solution operators

$$S : H_{q_1, q_2}(\Omega_e) \rightarrow H_0^s(\Omega), \quad F \mapsto (S_{q_2} - S_{q_1})F = u_2 - u_1 \in H_0^s(\Omega),$$

is linear and bounded by Corollary 2.2. Using Lemma 3.8 with $q := q_2$ we obtain a subspace $W \subseteq H_0^s(\Omega)$ with $\dim(W) = d(q_2)$, so that (3.4) holds for all F with

$SF \in W^\perp$ which is equivalent to $F \in (S^*W)^\perp$. Also, by Lemma 3.8, (3.5) holds for all F with $SF \in (W + N_q)^\perp$ which is equivalent to $F \in (S^*(W + N_q))^\perp$. Hence, the assertion follows with $V := S^*W$, and $V_+ := S^*(W + N_q)$. \square

Proof of Theorem 3.3. This immediately follows using the Lemmas 3.7–3.9. \square

Proof of Theorem 3.4. The monotonicity relation (3.2) immediately follows using Lemmas 3.7–3.9. To prove (3.3), we use that

$$0 = \mathcal{B}_{q_1}(u_1, w) = \mathcal{B}_{q_2}(u_2, w) \quad \text{for all } w \in H_0^s(\Omega),$$

to conclude that for all $D \subseteq \Omega$ containing $\text{supp}(q_1 - q_2)$

$$\begin{aligned} \lambda \|u_2 - u_1\|_{H^s(\mathbb{R}^n)}^2 &\leq \mathcal{B}_{q_2}(u_2 - u_1, u_2 - u_1) = -\mathcal{B}_{q_2}(u_1, u_2 - u_1) \\ &= \mathcal{B}_{q_1}(u_1, u_2 - u_1) - \mathcal{B}_{q_2}(u_1, u_2 - u_1) = \int_{\Omega} (q_1 - q_2)u_1(u_2 - u_1) \, dx \\ &\leq \|q_1 - q_2\|_{L^\infty(D)} \|u_1\|_{L^2(D)} \|u_2 - u_1\|_{H^s(\mathbb{R}^n)}. \end{aligned}$$

Hence

$$\|u_2\|_{L^2(D)} - \|u_1\|_{L^2(D)} \leq \|u_2 - u_1\|_{L^2(D)} \leq \frac{1}{\lambda} \|q_1 - q_2\|_{L^\infty(D)} \|u_1\|_{L^2(D)},$$

which yields (3.3) with $c := 1 + \frac{1}{\lambda} \|q_1 - q_2\|_{L^\infty(D)}$. \square

Proof of Theorem 3.6. For $q_j \in L^\infty(\Omega)$, $j = 1, 2$, we denote the positive eigenvalues (counted with multiplicities) of the compact self-adjoint operator

$$\iota^* \iota - \iota^* M_{q_j} \iota : H_0^s(\Omega) \rightarrow H_0^s(\Omega), \quad \text{by } \lambda_1^{(j)} \geq \lambda_2^{(j)} \geq \lambda_3^{(j)} \geq \dots$$

(a) Let $q_1 \leq q_2$. Then for all $v \in H_0^s(\Omega)$

$$\begin{aligned} ((\iota^* \iota - \iota^* M_{q_1} \iota)v, v)_{H_0^s(\Omega)} &= \int_{\Omega} (1 - q_1)|v|^2 \, dx \geq \int_{\Omega} (1 - q_2)|v|^2 \, dx \\ &= ((\iota^* \iota - \iota^* M_{q_2} \iota)v, v)_{H_0^s(\Omega)}. \end{aligned}$$

Hence, it follows from the Courant-Fischer-Weyl min-max principle, (see, e.g., [63]) that

$$\begin{aligned} \lambda_k^{(1)} &= \max_{\substack{X \subset H_0^s(\Omega) \\ \dim(X)=k}} \min_{\substack{v \in X \\ \|v\|_{H_0^s(\Omega)}=1}} ((\iota^* \iota - \iota^* M_{q_1} \iota)v, v)_{H_0^s(\Omega)} \\ &\geq \max_{\substack{X \subset H_0^s(\Omega) \\ \dim(X)=k}} \min_{\substack{v \in X \\ \|v\|_{H_0^s(\Omega)}=1}} ((\iota^* \iota - \iota^* M_{q_2} \iota)v, v)_{H_0^s(\Omega)} = \lambda_k^{(2)}, \end{aligned}$$

for all $k \in \mathbb{N}$, which shows $d(q_1) \geq d(q_2)$.

(b) Let $q_1 \in L^\infty(\Omega)$. Since $N_{q_1} = \mathcal{N}(I - \iota^* \iota + \iota^* M_{q_1} \iota)$, exactly $\dim(N_{q_1})$ eigenvalues of $\iota^* \iota - \iota^* M_{q_1} \iota$ are identically one, so that

$$\dots \geq \lambda_{d(q_1)}^{(1)} > 1 = \lambda_{d(q_1)+1}^{(1)} = \dots = \lambda_{d(q_1)+\dim(N_{q_1})}^{(1)} > \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(1)} \geq \dots$$

Since $\lambda_{d(q_1)}^{(1)} - 1 > 0$ and $1 - \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(1)} > 0$, we can set

$$\epsilon := \frac{1}{2} \min \left\{ \lambda_{d(q_1)}^{(1)} - 1, 1 - \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(1)} \right\} > 0. \quad (3.6)$$

Then for all $q_2 \in L^\infty(\Omega)$ with $\|q_2 - q_1\|_{L^\infty(\Omega)} \leq \epsilon$, and all $v \in H_0^s(\Omega)$ with $\|v\|_{H_0^s(\Omega)} = 1$, we have that

$$\left| ((\iota^* \iota - \iota^* M_{q_1} \iota)v, v)_{H_0^s(\Omega)} - ((\iota^* \iota - \iota^* M_{q_2} \iota)v, v)_{H_0^s(\Omega)} \right| \leq \int_{\Omega} |q_1 - q_2| |v|^2 dx \leq \epsilon.$$

Hence, using the Courant-Fischer-Weyl min-max principle as in (a) again, we obtain that $\left| \lambda_k^{(1)} - \lambda_k^{(2)} \right| \leq \epsilon$ for all $k \in \mathbb{N}$. In particular, using the definition of ϵ in (3.6), $\left| \lambda_{d(q_1)}^{(1)} - \lambda_{d(q_1)}^{(2)} \right| \leq \epsilon$ yields that

$$\lambda_1^{(2)} \geq \dots \geq \lambda_{d(q_1)}^{(2)} \geq \lambda_{d(q_1)}^{(1)} - \epsilon > 1,$$

and $\left| \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(1)} - \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(2)} \right| \leq \epsilon$ yields that

$$1 > \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(1)} + \epsilon \geq \lambda_{d(q_1)+\dim(N_{q_1})+1}^{(2)} \geq \lambda_{d(q_1)+\dim(N_{q_1})+2}^{(2)} \geq \dots$$

It follows that only the eigenvalues $\lambda_{d(q_1)+1}^{(2)}, \dots, \lambda_{d(q_1)+\dim(N_{q_1})}^{(2)}$ of $\iota^* \iota - \iota^* M_{q_2} \iota$ could possibly be identically one, so that $\dim(N_{q_2}) \leq \dim(N_{q_1})$ is proven.

□

3.2. Localized potentials for the fractional Schrödinger equation.

In this subsection, we extend the localized potentials result that was derived in [39] for positive potentials to general $L^\infty(\Omega)$ -potentials and spaces of finite codimension. Moreover, we will show a new result on controlling two localized potentials simultaneously. We will prove the following two theorems.

THEOREM 3.10 (Localized potentials). *Let $q \in L^\infty(\Omega)$. For every measurable set $M \subseteq \Omega$ with positive measure, and every finite-dimensional subspace $V \subseteq H_q(\Omega_e)$ there exists a sequence $\{F^k\}_{k \in \mathbb{N}} \subseteq V^\perp \subseteq H_q(\Omega_e)$ so that the corresponding solutions $u^k \in H_q^s(\mathbb{R}^n)$ of*

$$(-\Delta)^s u + qu = 0 \quad \text{in } \Omega, \quad \text{with } u|_{\Omega_e} = F^k, \quad (3.7)$$

fulfill

$$\int_M |u^k|^2 dx \rightarrow \infty, \quad \text{and} \quad \int_{\Omega \setminus M} |u^k|^2 dx \rightarrow 0.$$

THEOREM 3.11 (Simultaneously localized potentials). *Let $q_1, q_2 \in L^\infty(\Omega)$, and let $\text{supp}(q_1 - q_2) \subseteq M$ where $M \subseteq \Omega$ is a measurable set with positive measure. For every finite-dimensional subspace $V \subseteq H_{q_1, q_2}(\Omega_e)$, there exists a sequence $\{F^k\}_{k \in \mathbb{N}} \subseteq V^\perp \subseteq H_{q_1, q_2}(\Omega_e)$ so that the corresponding solutions $u_1^k \in H_{q_1}^s(\mathbb{R}^n)$, $u_2^k \in H_{q_2}^s(\mathbb{R}^n)$, of*

$$\begin{aligned} (-\Delta)^s u_1^k + q_1 u_1^k &= 0 \quad \text{in } \Omega, \quad \text{with } u_1^k|_{\Omega_e} = F^k, \\ (-\Delta)^s u_2^k + q_2 u_2^k &= 0 \quad \text{in } \Omega, \quad \text{with } u_2^k|_{\Omega_e} = F^k, \end{aligned}$$

fulfill

$$\begin{aligned} \int_M |u_1^k|^2 dx &\rightarrow \infty, & \int_{\Omega \setminus M} |u_1^k|^2 dx &\rightarrow 0, \\ \int_M |u_2^k|^2 dx &\rightarrow \infty, & \int_{\Omega \setminus M} |u_2^k|^2 dx &\rightarrow 0. \end{aligned}$$

To prove Theorem 3.10 and 3.11, we follow the general line of reasoning developed by one of the authors in [29]. We formulate the energy terms as norms of operator evaluations and characterize their adjoints and the ranges of their adjoints using the unique continuation property in Section 2.4. We then prove the two theorems using a functional analytic relation between norms of operator evaluations and ranges of their adjoints.

We start by defining the so-called virtual measurement operators.

LEMMA 3.12. *For $q \in L^\infty(\Omega)$, a measurable set $M \subseteq \Omega$ with positive measure, and a subspace $H \subseteq H_q(\Omega_e)$ with finite codimension, we define the operator*

$$L_{M,q} : H \rightarrow L^2(M), \quad F \mapsto u|_M,$$

where $u \in H_q^s(\mathbb{R}^n)$ solves

$$(-\Delta)^s u + qu = 0 \quad \text{in } \Omega, \quad \text{with } u|_{\Omega_e} = F. \quad (3.8)$$

Furthermore, let $V_M := \{u|_M : u \in N_q\}$.

Then $L_{M,q}$ is a linear bounded operator, $\dim(V_M) < \infty$, and for all $g \in V_M^\perp \subseteq L^2(M)$ and $F \in H$

$$(L_{M,q}^* g, F)_{H(\Omega_e)} = -\langle (-\Delta)^s v|_{\Omega_e}, F \rangle, \quad (3.9)$$

where $v \in H_q^s(\mathbb{R}^n)$ solves $(-\Delta)^s v + qv = g\chi_M$ in Ω , and $v|_{\Omega_e} = 0$.

Proof. By Lemma 2.1 and Corollary 2.2, we have that $L_{M,q}$ is a linear bounded operator, $\dim(V_M) < \infty$, and for all $g \in V_M^\perp \subseteq L^2(M)$ there exists a solution $v \in H_q^s(\Omega)$ of $(-\Delta)^s v + qv = g\chi_M$ in Ω , and $v|_{\Omega_e} = 0$. Then $v \in H_0^s(\Omega)$ fulfills

$$\mathcal{B}_q(v, w) = \int_M gw \, dx \quad \text{for all } w \in H_0^s(\Omega).$$

For $F \in H$ let $u = u^{(0)} + u^{(F)}$ solve (3.8) as in Lemma 2.1. Then

$$\begin{aligned} (L_{M,q}^* g, F)_{H(\Omega_e)} &= \int_M g(L_{M,q} F) \, dx = \int_M gu \, dx = \int_M g(u^{(0)} + u^{(F)}) \, dx \\ &= \mathcal{B}_q(v, u^{(0)}) + \int_M gu^{(F)} \, dx = -\mathcal{B}_q(v, u^{(F)}) + \int_M gu^{(F)} \, dx \\ &= -\int_{\mathbb{R}^n} (-\Delta)^{s/2} v \cdot (-\Delta)^{s/2} u^{(F)} \, dx + \int_\Omega (-\Delta)^s v \cdot v^{(F)} \, dx \\ &= -\langle (-\Delta)^s v|_{\Omega_e}, F \rangle. \end{aligned}$$

□

We now proceed similarly to [45] to extend the functional analytic relation between the norms of two operators and the ranges of their adjoints from [29, Lemma 2.5], [23, Corollary 3.5] to spaces of finite codimension.

LEMMA 3.13. *Let X, Y and Z be Hilbert spaces, $A_1 : X \rightarrow Y$ and $A_2 : X \rightarrow Z$ be linear bounded operators, and let $N \subseteq X$ be a finite dimensional subspace. Then*

$$\mathcal{R}(A_1^*) \subseteq \mathcal{R}(A_2^*) + N \quad \text{if and only if} \quad \exists c > 0 : \|A_1 x\| \leq c \|A_2 x\| \quad \forall x \in N^\perp,$$

where $\mathcal{R}(A)$ denotes the range of the linear bounded operator A .

Proof. For both implications, we use that there exists an orthogonal projection operator $P_N : X \rightarrow X$ with

$$\mathcal{R}(P_N) = N, \quad \mathcal{N}(P_N) = \mathcal{R}(I - P_N) = N^\perp, \quad \text{and} \quad P_N^2 = P_N = P_N^*.$$

To show the first implication, let $\mathcal{R}(A_1^*) \subseteq \mathcal{R}(A_2^*) + N$. Using block operator matrix notation we then have that

$$\mathcal{R}(A_1^*) \subseteq \mathcal{R}(A_2^*) + \mathcal{R}(P_N) = \mathcal{R}\left(\begin{pmatrix} A_2^* & P_N \end{pmatrix}\right).$$

Hence, by [29, Lemma 2.5] there exists $c > 0$ so that

$$\|A_1x\|^2 \leq c^2 \left\| \begin{pmatrix} A_2 \\ P_N \end{pmatrix} x \right\|^2 = c^2 \|A_2x\|^2 + c^2 \|P_Nx\|^2 \quad \text{for all } x \in X,$$

and thus

$$\|A_1x\| \leq c \|A_2x\| \quad \text{for all } x \in \mathcal{N}(P_N) = N^\perp.$$

To show the converse implication, let $c > 0$ and $\|A_1x\| \leq c \|A_2x\|$ for all $x \in N^\perp$. Then

$$\|A_1(I - P_N)x\| \leq c \|A_2(I - P_N)x\| \quad \text{for all } x \in X,$$

so that [29, Lemma 2.5] yields that

$$\mathcal{R}((I - P_N)A_1^*) \subseteq \mathcal{R}((I - P_N)A_2^*).$$

Hence,

$$\mathcal{R}(A_1^*) \subseteq \mathcal{R}((I - P_N)A_1^*) + N \subseteq \mathcal{R}((I - P_N)A_2^*) + N \subseteq \mathcal{R}(A_2^*) + N.$$

□

For the application of Lemma 3.13, the following elementary (and purely algebraic) observation will also be useful.

LEMMA 3.14. *Let X and Y be vector spaces, let $A : X \rightarrow Y$ be linear, and let Y' be a subspace of Y . The following two statements are equivalent:*

- (a) *There exists a finite dimensional subspace $N \subseteq Y$ with $A(X) \subseteq Y' + N$.*
- (b) *There exists a subspace $X' \subseteq X$ with finite codimension so that $A(X') \subseteq Y'$.*

Moreover, for all subspaces $X' \subseteq X$ with finite codimension, there exists a finite dimensional subspace $N \subseteq Y$ with $A(X) \subseteq A(X') + N$, and $\dim(A(X')) = \infty$ holds if $\dim A(X) = \infty$.

Proof. Let $A(X) \subseteq Y' + N$, where Y' and N are subspaces of Y , and $\dim(N) < \infty$. Since any basis of N can be extended to a Hamel basis of $Y' + N$, there exists a linear projection

$$P : Y' + N \rightarrow N \quad \text{with} \quad \mathcal{R}(P) = N, \quad \text{and} \quad \mathcal{N}(P) \subseteq Y'.$$

Define $X' := \{x \in X : PAx = 0\} = \mathcal{N}(PA)$. Then

$$\text{codim}(X') = \dim(\mathcal{R}(PA)) \leq \dim(\mathcal{R}(P)) = \dim(N),$$

and by definition $A(X') \subseteq \mathcal{N}(P) \subseteq Y'$. This shows that (a) implies (b).

Clearly, (b) implies (a) by setting $N := A(X'')$ where X'' is a linear complement of X' in X .

Moreover, if X' is a subspace of finite codimension then (b) holds with $Y' = A(X')$, so that (a) implies the existence of a finite dimensional subspace $N \subseteq Y$ with $A(X) \subseteq A(X') + N$. Clearly, this also implies that $\dim(A(X')) = \infty$ if $\dim(A(X)) = \infty$. \square

Now, we are ready to prove Theorem 3.10 and Theorem 3.11.

Proof of Theorem 3.10. Let $q \in L^\infty(\Omega)$, $M \subseteq \Omega$ be a measurable set with positive measure, and $V \subseteq H_q(\Omega_e)$ be a finite-dimensional subspace. As in Lemma 3.12, we define the virtual measurement operators

$$\begin{aligned} L_{M,q} : H_q(\Omega_e) &\rightarrow L^2(M), & F &\mapsto u|_M, & \text{and} \\ L_{\Omega \setminus M, q} : H_q(\Omega_e) &\rightarrow L^2(\Omega \setminus M), & F &\mapsto u|_{\Omega \setminus M}, \end{aligned}$$

where $u \in H_q^s(\mathbb{R}^n)$ solves

$$(-\Delta)^s u + qu = 0 \quad \text{in } \Omega \quad \text{with } u|_{\Omega_e} = F.$$

Then the assertion follows if we can show that there exists a sequence $\{F^k\}_{k \in \mathbb{N}} \subseteq V^\perp \subseteq H_q(\Omega_e)$ so that

$$\|L_{M,q} F^k\|_{L^2(M)} \rightarrow \infty, \quad \text{and} \quad \|L_{\Omega \setminus M, q} F^k\|_{L^2(\Omega \setminus M)} \rightarrow 0.$$

By a simple normalization argument (cf., e.g., the proof of [39, Corollary 3.5]), it suffices to show that

$$\exists c > 0 : \|L_{M,q} F\|_{L^2(M)} \leq c \|L_{\Omega \setminus M, q} F\|_{L^2(\Omega \setminus M)} \quad \text{for all } F \in V^\perp \subseteq H_q(\Omega_e).$$

This follows from Lemma 3.13 if we can show that

$$\mathcal{R}(L_{M,q}^*) \not\subseteq \mathcal{R}(L_{\Omega \setminus M, q}^*) + V. \quad (3.10)$$

We prove this by contradiction and assume that $\mathcal{R}(L_{M,q}^*) \subseteq \mathcal{R}(L_{\Omega \setminus M, q}^*) + V$.

As in Lemma 3.12, define

$$V_M := \{u|_M : u \in N_q\}, \quad \text{and} \quad V_{\Omega \setminus M} := \{u|_{\Omega \setminus M} : u \in N_q\}.$$

Then V_M^\perp and $V_{\Omega \setminus M}^\perp$ have finite codimension in $L^2(M)$ and $L^2(\Omega \setminus M)$, respectively. Moreover, we define their subspaces

$$\begin{aligned} W_M &:= \left\{ g \in V_M^\perp : \langle (-\Delta)^s v_M|_{\Omega_e}, F \rangle = 0 \text{ for all } F \in H_q(\Omega_e)^\perp \right\}, \\ W_{\Omega \setminus M} &:= \left\{ g \in V_{\Omega \setminus M}^\perp : \langle (-\Delta)^s v_{\Omega \setminus M}|_{\Omega_e}, F \rangle = 0 \text{ for all } F \in H_q(\Omega_e)^\perp \right\}, \end{aligned}$$

where $v_M, v_{\Omega \setminus M} \in H_q^s(\mathbb{R}^n)$ are the solutions of

$$(-\Delta)^s v_M + qv_M = g_M \chi_M \quad \text{in } \Omega, \quad v_M|_{\Omega_e} = 0, \quad (3.11)$$

$$(-\Delta)^s v_{\Omega \setminus M} + qv_{\Omega \setminus M} = g_{\Omega \setminus M} \chi_{\Omega \setminus M} \quad \text{in } \Omega, \quad v_{\Omega \setminus M}|_{\Omega_e} = 0. \quad (3.12)$$

Then also W_M and $W_{\Omega \setminus M}$ are subspaces of $L^2(M)$, resp., $L^2(\Omega \setminus M)$, with finite codimension, since the conditions in their definitions are equivalent to a system of finitely many homogeneous linear equations.

From Lemma 3.14 we then obtain that

$$L_{M,q}^*(W_M) \subseteq \mathcal{R}(L_{M,q}^*) \subseteq \mathcal{R}(L_{\Omega \setminus M,q}^*) + V \subseteq L_{\Omega \setminus M,q}^*(W_{\Omega \setminus M}) + V',$$

with a finite-dimensional space V' . Moreover, using Lemma 3.14 again, there exists a subspace $W'_M \subseteq W_M$ with finite codimension in W_M and thus in $L^2(M)$, so that

$$L_{M,q}^*(W'_M) \subseteq L_{\Omega \setminus M,q}^*(W_{\Omega \setminus M}). \quad (3.13)$$

Let $g_M \in W'_M$. Then, by (3.13), there exists $g_{\Omega \setminus M} \in W_{\Omega \setminus M}$, so that the corresponding solutions $v_M, v_{\Omega \setminus M} \in H_q^s(\mathbb{R}^n)$ of (3.11) and (3.12) fulfill

$$-\langle (-\Delta)^s v_M|_{\Omega_e}, F \rangle = (L_{M,q}^* g_M, F)_{H(\Omega_e)} = -\langle (-\Delta)^s v_{\Omega \setminus M}|_{\Omega_e}, F \rangle \text{ for all } F \in H_q(\Omega_e),$$

where we have utilized (3.9). By definition of W_M and $W_{\Omega \setminus M}$, it also holds that

$$\langle (-\Delta)^s v_M|_{\Omega_e}, F \rangle = 0 = \langle (-\Delta)^s v_{\Omega \setminus M}|_{\Omega_e}, F \rangle \quad \text{for all } F \in H_q(\Omega_e)^\perp.$$

Hence $v := v_M - v_{\Omega \setminus M}$ fulfills

$$(-\Delta)^s v + qv = g_M \chi_M - g_{\Omega \setminus M} \chi_{\Omega \setminus M} \quad \text{in } \Omega$$

with vanishing Cauchy data $v|_{\Omega_e} = 0$ and $(-\Delta)^s v|_{\Omega_e} = 0$. From the unique continuation result in Corollary 2.6(b) it follows that $v \equiv 0$ in \mathbb{R}^n . But this yields $g_M = 0$, and since this arguments holds for all $g_M \in W'_M$, it follows that $W'_M = \{0\}$ which contradicts the fact that W'_M is a subspace of finite codimension in the infinite dimensional space $L^2(M)$. Hence, (3.10) and thus the assertion is proven. \square

Proof of Theorem 3.11. Let $q_1, q_2 \in L^\infty(\Omega)$, and let $\text{supp}(q_1 - q_2) \subseteq M$ where $M \subseteq \Omega$ is a measurable set with positive measure. We first note that it suffices to show that for all finite-dimensional subspaces $V \subseteq H_{q_1, q_2}(\Omega_e)$, there exists a sequence $\{F^k\}_{k \in \mathbb{N}} \subseteq V^\perp \subseteq H_{q_1, q_2}(\Omega_e)$ with

$$\int_M |u_1^k|^2 dx \rightarrow \infty, \quad \text{and} \quad \int_{\Omega \setminus M} (|u_1^k|^2 + |u_2^k|^2) dx \rightarrow 0, \quad (3.14)$$

since $\int_M |u_1^k|^2 dx \rightarrow \infty$ implies $\int_M |u_2^k|^2 dx \rightarrow \infty$ on a subspace of finite codimension in $H_{q_1, q_2}(\Omega_e)$ by Remark 3.5.

We define as in Lemma 3.12,

$$\begin{aligned} L_{M, q_1} &: H_{q_1, q_2}(\Omega_e) \rightarrow L^2(M), & F &\mapsto u_1|_M, \\ L_{\Omega \setminus M, q_1} &: H_{q_1, q_2}(\Omega_e) \rightarrow L^2(\Omega \setminus M), & F &\mapsto u_1|_{\Omega \setminus M}, \quad \text{and} \\ L_{\Omega \setminus M, q_2} &: H_{q_1, q_2}(\Omega_e) \rightarrow L^2(\Omega \setminus M), & F &\mapsto u_2|_{\Omega \setminus M}, \end{aligned}$$

where $u_j \in H_{q_j}^s(\mathbb{R}^n)$ solves (for $j = 1, 2$)

$$(-\Delta)^s u_j + q_j u_j = 0 \quad \text{in } \Omega \quad \text{with } u_j|_{\Omega_e} = F.$$

Thus (3.14) can be reformulated as

$$\|L_{M, q_1} F^k\|_{L^2(M)} \rightarrow \infty \quad \text{and} \quad \left\| \begin{pmatrix} L_{\Omega \setminus M, q_1} \\ L_{\Omega \setminus M, q_2} \end{pmatrix} F^k \right\|_{L^2(\Omega \setminus M) \times L^2(\Omega \setminus M)} \rightarrow 0.$$

Hence, using Lemma 3.13 as in the proof of Theorem 3.10, the assertion follows if we can show that

$$\mathcal{R}(L_{M,q_1}^*) \not\subseteq \mathcal{R}\left(\begin{pmatrix} L_{\Omega \setminus M,q_1}^* & L_{\Omega \setminus M,q_2}^* \end{pmatrix}\right) + V = \mathcal{R}(L_{\Omega \setminus M,q_1}^*) + \mathcal{R}(L_{\Omega \setminus M,q_2}^*) + V. \quad (3.15)$$

We argue by contradiction and assume that

$$\overline{\mathcal{R}(L_{M,q_1}^*)} \subseteq \mathcal{R}(L_{\Omega \setminus M,q_1}^*) + \mathcal{R}(L_{\Omega \setminus M,q_2}^*) + V.$$

As in the proof of Theorem 3.10, we define (for $j = 1, 2$)

$$V_{M,q_1} := \{u|_M : u \in N_{q_1}\}, \quad V_{\Omega \setminus M,q_j} := \{u|_{\Omega \setminus M} : u \in N_{q_j}\}.$$

and

$$\begin{aligned} W_{M,q_1} &:= \{g \in V_{M,q_1}^\perp : \langle (-\Delta)^s v_{M,q_1}|_{\Omega_e}, F \rangle = 0 \text{ for all } F \in H_{q_1,q_2}(\Omega_e)^\perp\}, \\ W_{\Omega \setminus M,q_j} &:= \{g \in V_{\Omega \setminus M,q_j}^\perp : \langle (-\Delta)^s v_{\Omega \setminus M,q_j}|_{\Omega_e}, F \rangle = 0 \text{ for all } F \in H_{q_1,q_2}(\Omega_e)^\perp\}, \end{aligned}$$

where $v_{M,q_1}, v_{\Omega \setminus M,q_j} \in H_q^s(\mathbb{R}^n)$ are the solutions of

$$(-\Delta)^s v_{M,q_1} + q_1 v_{M,q_1} = g_{M,q_1} \chi_M \quad \text{in } \Omega, \quad v_{M,q_1}|_{\Omega_e} = 0, \quad (3.16)$$

$$(-\Delta)^s v_{\Omega \setminus M,q_j} + q_j v_{\Omega \setminus M,q_j} = g_{\Omega \setminus M,q_j} \chi_{\Omega \setminus M} \quad \text{in } \Omega, \quad v_{\Omega \setminus M,q_j}|_{\Omega_e} = 0, \quad (3.17)$$

for $j = 1, 2$. Then, as in the proof of Theorem 3.10, we obtain using lemma 3.14 that

$$L_{M,q_1}^*(W'_{M,q_1}) \subseteq L_{\Omega \setminus M,q_1}^*(W_{\Omega \setminus M,q_1}) + L_{\Omega \setminus M,q_2}^*(W_{\Omega \setminus M,q_2}) \quad (3.18)$$

with a subspace $W'_{M,q_1} \subseteq W_{M,q_1}$ that has finite codimension in $L^2(M)$.

Let $g_{M,q_1} \in W'_{M,q_1}$. As in the proof of Theorem 3.10, it then follows from (3.18) and the definition of W_{M,q_1} , $W_{\Omega \setminus M,q_j}$, and $W_{\Omega \setminus M,q_2}$, that there exist $g_{q_j, \Omega \setminus M} \in W_{\Omega \setminus M,q_j}$ ($j = 1, 2$), so that the solutions v_{M,q_1} , $v_{\Omega \setminus M,q_1}$, and $v_{\Omega \setminus M,q_2}$ of (3.16) and (3.17) fulfill

$$(-\Delta)^s v_{q_1, M}|_{\Omega_e} = (-\Delta)^s v_{q_1, \Omega \setminus M}|_{\Omega_e} + (-\Delta)^s v_{q_2, \Omega \setminus M}|_{\Omega_e}.$$

It follows that $v := v_{q_1, \Omega \setminus M} + v_{q_2, \Omega \setminus M} - v_{q_1, M}$ solves

$$(-\Delta)^s v + q_1 v = g_{q_1, \Omega \setminus M} \chi_{\Omega \setminus M} + (q_1 - q_2) v_{q_2, \Omega \setminus M} + g_{q_2, \Omega \setminus M} \chi_{\Omega \setminus M} - g_{q_1, M} \chi_M$$

with zero Cauchy data. Hence, by Corollary 2.6(b), $v = 0$, and with $\text{supp}(q_1 - q_2) \subseteq M$ this also implies

$$(q_1 - q_2) v_{q_2, \Omega \setminus M} - g_{q_1, M} = 0.$$

Since $v_{q_2, \Omega \setminus M} \in H_0^s(\Omega)$, and the above arguments hold for all $g_{M,q_1} \in W'_{M,q_1}$, it follows that

$$W'_{M,q_1} \subseteq (M_{q_1} - M_{q_2})\iota(H_0^s(\Omega)) \subseteq L^2(\Omega).$$

Hence, the range of the compact operator $(M_{q_1} - M_{q_2})\iota$ would be a subspace of finite codimension in $L^2(\Omega)$ and thus closed. But the range of a compact operator can only be closed if it is finite dimensional (cf., e.g., [68, Theorem. 4.18]), so that this contradicts the infinite dimensionality of $L^2(\Omega)$. Thus, (3.15) is proven. \square

4. Converse monotonicity, uniqueness and inclusion detection. Using the localized potentials and monotonicity relations from the last section, we can now extend the results from [39] to the case of a general potential $q \in L^\infty(\Omega)$.

4.1. Converse monotonicity and the Calderón problem. We first derive an if-and-only-if monotonicity relation between the potential and the DtN operators.

THEOREM 4.1. *Let $n \in \mathbb{N}$, $\Omega \subset \mathbb{R}^n$ be a Lipschitz bounded open set and $s \in (0, 1)$. For any two potential $q_1, q_2 \in L^\infty(\Omega)$, we have*

$$q_1 \geq q_2 \quad \text{if and only if} \quad \Lambda(q_1) \geq_{d(q_2)} \Lambda(q_2) \quad \text{if and only if} \quad \Lambda(q_1) \geq_{\text{fin}} \Lambda(q_2), \quad (4.1)$$

where $d(q_2)$ is the integer given in Section 3.

Proof. Via Theorem 3.3, $q_1 \geq q_2$ implies $\Lambda(q_1) \geq_{d(q_2)} \Lambda(q_2)$, and clearly $\Lambda(q_1) \geq_{d(q_2)} \Lambda(q_2)$ implies $\Lambda(q_1) \geq_{\text{fin}} \Lambda(q_2)$. The assertion is proven if we can show that $\Lambda(q_1) \geq_{\text{fin}} \Lambda(q_2)$ implies $q_1 \geq q_2$ a.e. in Ω .

Let $\Lambda(q_1) \geq_{\text{fin}} \Lambda(q_2)$. Using this together with Remark 3.5 and that the intersection of subspaces with finite codimension still has finite codimension, we obtain a subspace $V \subseteq H_{q_1, q_2}(\Omega_e)$ so that

$$0 \leq \langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \leq \int_{\Omega} (q_1 - q_2) |u_2|^2 dx \quad \text{for all } F \in V^\perp \subseteq H_{q_1, q_2}(\Omega_e), \quad (4.2)$$

where $u_2 \in H_{q_2}^s(\mathbb{R}^n)$ solves

$$(-\Delta)^s u_2 + q_2 u_2 = 0 \text{ in } \Omega, \quad \text{and} \quad u_2|_{\Omega_e} = F. \quad (4.3)$$

To show that this implies $q_1 \geq q_2$ a.e. in Ω , we argue by contradiction and assume that there exists $\delta > 0$ and a positive measurable set $M \subset \Omega$ such that $q_2 - q_1 \geq \delta$ on M . Then utilizing the localized potentials from Theorem 3.10 we obtain a sequence $(F^k)_{k \in \mathbb{N}} \subset V^\perp \subseteq H_{q_1, q_2}(\Omega_e)$ where the corresponding solutions of (4.3) with $F = F^k$ solve

$$\int_M |u_2^k|^2 dx \rightarrow \infty, \quad \text{and} \quad \int_{\Omega \setminus M} |u_2^k|^2 dx \rightarrow 0.$$

But together with (4.2) this yields to the contradiction

$$0 \leq \int_{\Omega} (q_1 - q_2) |u_2^k|^2 dx \leq -\delta \int_M |u_2^k|^2 dx + \|q_1 - q_2\|_{L^\infty(\Omega)} \int_{\Omega \setminus M} |u_2^k|^2 dx \rightarrow -\infty,$$

which proves $q_1 \geq q_2$ a.e. in Ω . \square

COROLLARY 4.2. *Let $n \in \mathbb{N}$, $\Omega \subset \mathbb{R}^n$ be a bounded Lipschitz domain and $s \in (0, 1)$. For any two potentials $q_0, q_1 \in L_+^\infty(\Omega)$,*

$$q_0 = q_1 \text{ if and only if } \Lambda(q_0) \stackrel{\text{fin}}{=} \Lambda(q_1).$$

Proof. This follows immediately from Theorem 4.1. \square

4.2. A monotonicity-based reconstruction formula. In [39], we considered positive potentials $q \in L_+^\infty(\Omega)$, where $L_+^\infty(\Omega)$ denotes the set of all $L^\infty(\Omega)$ -functions with positive essential infima. We showed that $q \in L_+^\infty(\Omega)$ can be reconstructed from $\Lambda(q)$ by taking the supremum of all positive density one simple functions ψ with $\Lambda(\psi) \leq \Lambda(q)$. The space of density one simple functions is defined by

$$\Sigma := \left\{ \psi = \sum_{j=1}^m a_j \chi_{M_j} : a_j \in \mathbb{R}, M_j \subseteq \Omega \text{ is a density one set} \right\},$$

where we call a subset $M \subseteq \Omega$ a *density one set* if it is non-empty, measurable and has Lebesgue density 1 in all $x \in M$. Note that density one simple functions can be regarded as simple functions where function values that are only attained on a null set are replaced by zero, and that, by the Lebesgue's density theorem, every measurable set agrees almost everywhere with a density one set, so that every simple function agrees with a density one simple function almost everywhere. For our results, it is important to control the values on null sets since these values might still affect the supremum when the supremum is taken over uncountably many functions.

For general potentials we obtain the following reconstruction formula.

THEOREM 4.3. *Let $n \in \mathbb{N}$, $\Omega \subset \mathbb{R}^n$ be a bounded Lipschitz domain and $s \in (0, 1)$. A potential $q \in L^\infty(\Omega)$ is uniquely determined by $\Lambda(q)$ via the following formula*

$$\begin{aligned} q(x) &= \sup\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \leq_{f_n} \Lambda(q)\} + \inf\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \geq_{f_n} \Lambda(q)\} \\ &= \sup\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \leq_{d(\psi)} \Lambda(q)\} + \inf\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \geq_{d(q)} \Lambda(q)\} \end{aligned}$$

for a.e. $x \in \Omega$.

To prove Theorem 4.3, we first show the following lemma.

LEMMA 4.4. *For each function $q \in L^\infty(\Omega)$, and $x \in \Omega$ a.e., we have that*

$$\max\{q(x), 0\} = \sup\{\psi(x) : \psi \in \Sigma \text{ with } \psi \leq q\}.$$

Proof. Let $q \in L^\infty(\Omega)$. By the standard simple function approximation lemma, cf., e.g., [67], there exists a sequence $(\psi_k), k \in \mathbb{N}$ of simple functions with

$$q(x) - \frac{1}{k} \leq \psi_k(x) \leq q(x) \quad (4.4)$$

for all $k \in \mathbb{N}$ and $x \in \Omega$. Since every simple function agrees with a density one simple function almost everywhere, we can change the values of the countably many functions ψ_k on a null set, to obtain $\psi_k \in \Sigma$ for which (4.4) holds almost everywhere. Hence, for a.e. $x \in \Omega$,

$$q(x) = \lim_{k \rightarrow \infty} \psi_k(x) \leq \sup\{\psi(x) : \psi \in \Sigma, \psi \leq q\}.$$

Moreover, if $x \in \Omega$ then $\psi_x = -\|q\|_{L^\infty(\Omega)} \chi_{\Omega \setminus \{x\}}$ is a density one simple function fulfilling $\psi_x(x) = 0$ and $\psi_x(\xi) \leq q(\xi)$ for a.e. $\xi \in \Omega$, so that $\psi_x \leq q$. Hence,

$$0 \leq \sup\{\psi(x) : \psi \in \Sigma, \psi \leq q\} \quad \text{for a.e. } x \in \Omega.$$

It remains to show that

$$\max\{q(x), 0\} \geq \sup\{\psi(x) : \psi \in \Sigma, \psi \leq q\} \quad \text{for a.e. } x \in \Omega. \quad (4.5)$$

We argue as in the proof of [39, Lemma 4.4]. It suffices to show that for each $\delta > 0$ the set

$$M := \{x \in \Omega : \max\{q(x), 0\} + \delta < \sup\{\psi(x) : \psi \in \Sigma, \psi \leq q\}\} \quad (4.6)$$

is a null set. To prove this, assume that M is not a null set for some $\delta > 0$. By removing a null set from M , we can assume that M is a density one set. By using Lusin's theorem (see [67] for instance), all measurable function are approximately continuous at almost every point. Hence, M must contain a point \hat{x} in which the function $x \mapsto \max\{q(x), 0\}$ is approximately continuous, and thus the set

$$M' := \{x \in \Omega : \max\{q(x), 0\} \leq \max\{q(\hat{x}), 0\} + \delta/3\}$$

has density one in \hat{x} . Removing a null set, we can assume that M' is a density one set still containing \hat{x} .

Moreover, by the definition of M , there must exist a $\psi \in \Sigma$ with $\psi \leq q$ and

$$\max\{q(\hat{x}), 0\} + \frac{2}{3}\delta \leq \psi(\hat{x}).$$

This shows $\psi(\hat{x}) > 0$, so that, by [39, Lemma 4.3], there exists a density one set M'' containing \hat{x} , where $\psi(x) = \psi(\hat{x})$ for all $x \in M''$.

We thus have that for all $x \in M' \cap M''$

$$q(x) + \delta/3 \leq \max\{q(x), 0\} + \delta/3 \leq \max\{q(\hat{x}), 0\} + \frac{2}{3}\delta \leq \psi(\hat{x}) = \psi(x),$$

and $M' \cap M''$ possesses positive measure since M' and M'' are density one sets that both contain \hat{x} , cf., again, [39, Lemma 4.3]. But this contradicts that $\psi(x) \leq q(x)$ almost everywhere, and thus shows that M defined in (4.6) is a null set for all $\delta > 0$. It follows that (4.5) holds, so that the assertion is proven. \square

Proof of Theorem 4.3. Using lemma 4.4 and the if-and-only-if monotonicity relation in Theorem 4.1, we have that for all $q \in L^\infty(\Omega)$, and all $x \in \Omega$ a.e.,

$$\begin{aligned} q(x) &= \max\{q(x), 0\} - \max\{-q(x), 0\} \\ &= \sup\{\psi(x) : \psi \in \Sigma, \psi \leq q\} - \sup\{\psi(x) : \psi \in \Sigma, \psi \leq -q\} \\ &= \sup\{\psi(x) : \psi \in \Sigma, \psi \leq q\} + \inf\{\psi(x) : \psi \in \Sigma, \psi \geq q\} \\ &= \sup\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \leq_{\text{fin}} \Lambda(q)\} + \inf\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \geq_{\text{fin}} \Lambda(q)\} \\ &= \sup\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \leq_{d(\psi)} \Lambda(q)\} + \inf\{\psi(x) : \psi \in \Sigma, \Lambda(\psi) \geq_{d(q)} \Lambda(q)\}. \end{aligned}$$

This completes the proof. \square

4.3. The linearized Calderón problem. In this subsection, we will only consider $q \in L^\infty(\Omega)$ that fulfill the following assumption.

DEFINITION 4.5. *Let N_q be the set defined by (2.5), then we say that $q \in L^\infty(\Omega)$ is non-resonant, if $N_q = \{0\}$.*

This assumption is also called an *eigenvalue condition* in the literature, since it is equivalent to $\{0\}$ being not an Dirichlet eigenvalue of the fractional operator $(-\Delta)^s + q$ in Ω . Note that it implies that $H_q(\Omega_e) = H(\Omega_e)$, and $H_q^s(\mathbb{R}^n) = H^s(\mathbb{R}^n)$, i.e., that the Dirichlet problem is uniquely solvable for all Dirichlet data in $H(\Omega_e)$, cf. Corollary 2.2.

We start by showing that the non-resonant potentials are an open subset of $L^\infty(\Omega)$, on which the DtN operator is Fréchet differentiable.

LEMMA 4.6. *The set $\mathcal{O} = \{q \in L^\infty(\Omega) : N_q = \{0\}\}$ is an open subset of $L^\infty(\Omega)$. On this set, the DtN operator*

$$\Lambda : \mathcal{O} \subseteq L^\infty(\Omega) \rightarrow \mathcal{L}(H(\Omega_e), H(\Omega_e)^*), \quad q \mapsto \Lambda(q),$$

is Fréchet differentiable. For each $q \in \mathcal{O}$ its derivative is given by

$$\begin{aligned} \Lambda'(q) : L^\infty(\Omega) &\rightarrow \mathcal{L}(H(\Omega_e), H(\Omega_e)^*), \quad r \mapsto \Lambda'(q)r, \\ \langle (\Lambda'(q)r)F, G \rangle &:= \int_{\Omega} r S_q(F) S_q(G) \, dx \quad \text{for all } r \in L^\infty(\Omega), \, F, G \in H(\Omega_e), \end{aligned}$$

where $S_q : H(\Omega_e) \rightarrow H^s(\mathbb{R}^n)$, $F \mapsto u$, is the solution operator of the Dirichlet problem

$$(-\Delta)^s u + qu = 0 \text{ in } \Omega \quad \text{and} \quad u|_{\Omega_e} = F.$$

Proof. The fact that \mathcal{O} is open immediately follows from Theorem 3.6(b).

Let $q \in \mathcal{O} \subseteq L^\infty(\Omega)$. $\Lambda'(q)$ is a linear bounded operator since S_q is linear and bounded, cf. Corollary 2.2. For sufficiently small $r \in L^\infty(\Omega)$, we have that $q+r \in \mathcal{O}$, and it follows from Lemma 2.4 that

$$\langle (\Lambda(q+r) - \Lambda(q))F, F \rangle = \int_{\Omega} r S_{q+r}(F) S_q(F) \, dx.$$

With the operator formulation from the proof of Lemma 2.1, it is then easy to show that, for sufficiently small $r \in L^\infty(\Omega)$, there exists a constant $C > 0$ with

$$\|S_{q+r}(F) - S_q(F)\|_{H^s(\mathbb{R}^n)} \leq C \|r\|_{L^\infty(\Omega)} \|F\|_{H(\Omega_e)}.$$

Using that $\Lambda(q)$, $\Lambda(q+r)$, and $\Lambda'(q)r$ are symmetric operators, it now follows that

$$\begin{aligned} &\|\Lambda(q+r) - \Lambda(q) - \Lambda'(q)r\|_{\mathcal{L}(H(\Omega_e), H(\Omega_e)^*)} \\ &= \sup_{\|F\|_{H(\Omega_e)}=1} |\langle (\Lambda(q+r) - \Lambda(q) - \Lambda'(q)r)F, F \rangle| \\ &= \sup_{\|F\|_{H(\Omega_e)}=1} \left| \int_{\Omega} r (S_{q+r}(F) - S_q(F)) S_q(F) \, dx \right| \leq C \|r\|_{L^\infty(\Omega)}^2 \|S_q\|_{\mathcal{L}(H(\Omega_e), H^s(\mathbb{R}^n))}. \end{aligned}$$

which proves the assertion. \square

Using the Fréchet derivative from Lemma 4.6, the monotonicity relations in Theorem 3.3 and 3.4 can now be written as follows.

COROLLARY 4.7. *For all non-resonant $q_1, q_2 \in L^\infty(\Omega)$,*

$$\Lambda'(q_2)(q_1 - q_2) \geq_{d(q_1)} \Lambda(q_1) - \Lambda(q_2) \geq_{d(q_2)} \Lambda'(q_1)(q_1 - q_2),$$

and there exists $c > 0$ so that for all measurable $D \subseteq \Omega$ containing $\text{supp}(q_1 - q_2)$

$$c\Lambda'(q_2)\chi_D \geq_{d(q_1)} \Lambda'(q_1)\chi_D \geq_{d(q_2)} \frac{1}{c}\Lambda'(q_2)\chi_D.$$

Proof. Since $q_1, q_2 \in L^\infty(\Omega)$ are non-resonant, we have that $H_{q_1}(\Omega_e) = H_{q_2}(\Omega_e) = H_{q_1, q_2}(\Omega_e) = H(\Omega_e)$. It then follows from Theorem 3.3 and Lemma 4.6 that there exists a subspace $V \subseteq H(\Omega_e)$ with $\dim(V) \leq d(q_2)$ so that for all $F \in V^\perp$

$$\langle (\Lambda(q_1) - \Lambda(q_2)) F, F \rangle \geq \int_{\Omega} (q_1 - q_2) S_{q_1}(F)^2 dx = \langle (\Lambda'(q_1)(q_1 - q_2)) F, F \rangle,$$

which shows that $\Lambda(q_1) - \Lambda(q_2) \geq_{d(q_2)} \Lambda'(q_1)(q_1 - q_2)$.

Also, it follows from Theorem 3.4 and Lemma 4.6 that there exists a subspace $V_+ \subseteq H(\Omega_e)$ with $\dim(V_+) \leq d(q_2) + N(q_2) = d(q_2)$ and a constant $\tilde{c} > 0$, so that for all measurable $D \subseteq \Omega$ containing $\text{supp}(q_1 - q_2)$, and all $F \in V_+^\perp$,

$$\langle (\Lambda'(q_2)\chi_D) F, F \rangle = \|S_{q_2}(F)\|_{L^2(D)}^2 \leq \tilde{c}^2 \|S_{q_1}(F)\|_{L^2(D)}^2 = \tilde{c}^2 \langle (\Lambda'(q_1)\chi_D) F, F \rangle,$$

which shows $\Lambda'(q_2)\chi_D \leq_{d(q_2)} c\Lambda'(q_1)\chi_D$ with $c := \tilde{c}^2$.

The other assertions follow by interchanging q_1 and q_2 . \square

We also have an if-and-only if monotonicity result for the linearized DtN-operators.

THEOREM 4.8. *Let $n \in \mathbb{N}$, $\Omega \subset \mathbb{R}^n$ be a Lipschitz bounded open set and $s \in (0, 1)$. Then for all non-resonant $q \in L^\infty(\Omega)$ and $r_1, r_2 \in L^\infty(\Omega)$,*

$$r_1 \leq r_2 \quad \text{if and only if} \quad \Lambda'(q)r_1 \leq \Lambda'(q)r_2 \quad \text{if and only if} \quad \Lambda'(q)r_1 \leq_{\text{fin}} \Lambda'(q)r_2.$$

Proof. If $r_1 \leq r_2$ then $\Lambda'(q)r_1 \leq \Lambda'(q)r_2$ follows immediately from the characterization of $\Lambda'(q)$ in Lemma 4.6. (Note that this holds on the whole space $H(\Omega_e)$, and not just on a subspace of finite codimension).

Clearly, $\Lambda'(q)r_1 \leq \Lambda'(q)r_2$ implies $\Lambda'(q)r_1 \leq_{\text{fin}} \Lambda'(q)r_2$, and the implication from $\Lambda'(q)r_1 \leq_{\text{fin}} \Lambda'(q)r_2$ to $r_1 \leq r_2$ follows from the same localized potentials argument as in the proof of Theorem 4.1. \square

This implies uniqueness of the linearized fractional Calderón problem:

COROLLARY 4.9. *Let $n \in \mathbb{N}$, $\Omega \subset \mathbb{R}^n$ be a Lipschitz bounded open set and $s \in (0, 1)$. For all non-resonant $q \in L^\infty(\Omega)$, the Fréchet derivative $\Lambda'(q)$ is injective, i.e.*

$$\Lambda'(q)r \stackrel{\text{fin}}{=} 0 \quad \text{if and only if} \quad \Lambda'(q)r = 0 \quad \text{if and only if} \quad r = 0.$$

Proof. This follows immediately from Theorem 4.8. \square

4.4. Inclusion detection by linearized monotonicity tests. In this section we will study the inclusion detection (or shape reconstruction) problem of determining regions where a non-resonant potential $q \in L^\infty(\Omega)$ changes from a known non-resonant reference potentials $q_0 \in L^\infty(\Omega)$, i.e., we aim to reconstruct the support $q - q_0$ by comparing $\Lambda(q)$ with $\Lambda(q_0)$. q_0 may describe a background coefficient, and q denotes the coefficient function in the presence of anomalies or scatterers.

We will generalize the results in [39] and show that the support of $q - q_0$ can be reconstructed with *linearized monotonicity tests* [47, 25]. These linearized tests only

utilize the solution of the fractional Schrödinger equation with the reference coefficient function $q_0 \in L^\infty(\Omega)$. They do not require any other special solutions of the equation.

In all of the following let $n \in \mathbb{N}$, $\Omega \subset \mathbb{R}^n$ be a Lipschitz bounded open set, $s \in (0, 1)$, and let $q_0, q \in L^\infty(\Omega)$ be non-resonant.

For a measurable subset $M \subseteq \Omega$, we introduce the testing operator $\mathcal{T}_M : H(\Omega_e) \rightarrow H(\Omega_e)^*$ by setting $\mathcal{T}_M := \Lambda'(q_0)\chi_M$. i.e.,

$$\langle \mathcal{T}_M F, G \rangle := \int_M S_{q_0}(F)S_{q_0}(G)dx \quad \text{for all } F, G \in H(\Omega_e), \quad (4.7)$$

where $S_{q_0} : H(\Omega_e) \rightarrow H^s(\mathbb{R}^n)$ denotes the solution operator as in as in Lemma 4.6.

The following theorem shows that we can find the support of $q - q_0$ by shrinking closed sets, cf. [47, 28].

THEOREM 4.10. *For each closed subset $C \subseteq \Omega$,*

$$\begin{aligned} \text{supp}(q - q_0) &\subseteq C, \\ \text{if and only if} \quad &\exists \alpha > 0 : -\alpha \mathcal{T}_C \leq_{d(q_0)+d(q)} \Lambda(q) - \Lambda(q_0) \leq_{d(q)} \alpha \mathcal{T}_C, \\ \text{if and only if} \quad &\exists \alpha > 0 : -\alpha \mathcal{T}_C \leq_{\text{fin}} \Lambda(q) - \Lambda(q_0) \leq_{\text{fin}} \alpha \mathcal{T}_C. \end{aligned}$$

Hence,

$$\begin{aligned} &\text{supp}(q - q_0) \\ &= \bigcap \{C \subseteq \Omega \text{ closed} : \exists \alpha > 0 : -\alpha \mathcal{T}_C \leq_{d(q_0)+d(q)} \Lambda(q) - \Lambda(q_0) \leq_{d(q)} \alpha \mathcal{T}_C\} \\ &= \bigcap \{C \subseteq \Omega \text{ closed} : \exists \alpha > 0 : -\alpha \mathcal{T}_C \leq_{\text{fin}} \Lambda(q) - \Lambda(q_0) \leq_{\text{fin}} \alpha \mathcal{T}_C\}. \end{aligned}$$

Proof.

- (a) Let $\text{supp}(q - q_0) \subseteq C$. Then, by Corollary 4.7, there exists a constant $c > 0$ with

$$\mathcal{T}_C = \Lambda'(q_0)\chi_C \geq_{d(q)} c\Lambda'(q)\chi_C.$$

Moreover, $\text{supp}(q - q_0) \subseteq C$ implies that for sufficiently large $\alpha > 0$

$$-\alpha c\chi_C \leq q - q_0 \leq \alpha c\chi_C.$$

Using Corollary 4.7 and Theorem 4.8, we thus obtain

$$\begin{aligned} \Lambda(q) \leq_{d(q)} \Lambda(q_0) + \Lambda'(q_0)(q - q_0) &\leq \Lambda(q_0) + \Lambda'(q_0)\alpha c\chi_C = \Lambda(q_0) + \alpha \mathcal{T}_C, \\ \Lambda(q) \geq_{d(q_0)} \Lambda(q_0) + \Lambda'(q)(q - q_0) &\geq \Lambda(q_0) - \alpha c\Lambda'(q)\chi_C \geq_{d(q)} \Lambda(q_0) - \alpha \mathcal{T}_C. \end{aligned}$$

- (b) We will now show that

$$\exists \alpha > 0 : -\alpha \mathcal{T}_C \leq_{\text{fin}} \Lambda(q) - \Lambda(q_0) \leq_{\text{fin}} \alpha \mathcal{T}_C \quad (4.8)$$

implies $\text{supp}(q - q_0) \subseteq C$.

Let $\alpha > 0$ fulfill (4.8). Then we obtain from the first inequality in (4.8) with Corollary 4.7

$$\Lambda'(q_0)(-\alpha c\chi_C) = -\alpha \mathcal{T}_C \leq_{\text{fin}} \Lambda(q) - \Lambda(q_0) \leq_{\text{fin}} \Lambda'(q_0)(q - q_0),$$

so that Theorem 4.8 yields that

$$-\alpha\chi_C \leq q - q_0. \quad (4.9)$$

It remains to show that the second inequality in (4.8) implies that

$$q - q_0 \leq 0 \quad \text{on } \Omega \setminus C. \quad (4.10)$$

We argue by contradiction and assume that (4.10) is not true. Then there exists $\delta > 0$, and a measurable subset $M \subseteq \Omega \setminus C$ with positive measure so that $q - q_0 \geq \delta$ on M .

We now use an idea from [41] to rewrite energy terms by repeated application of the monotonicity relation, and define

$$\tilde{q} := q + \delta\chi_M - \alpha\chi_C + (q_0 - q)\chi_{\Omega \setminus (M \cup C)} = \begin{cases} q + \delta & \text{in } M, \\ q - \alpha & \text{in } C, \\ q_0 & \text{in } \Omega \setminus (M \cup C), \end{cases}$$

and note that

$$q - q_0 \geq \delta\chi_M - \alpha\chi_C \geq \delta\chi_M - \alpha\chi_C + (q_0 - q)\chi_{\Omega \setminus (C \cup M)} = \tilde{q} - q.$$

Using Theorem 3.3 and Remark 3.5, there exists a finite dimensional subspace $V \subseteq H_{\tilde{q}}(\Omega_e)$ so that for all $F \in V^\perp \subseteq H_{\tilde{q}}(\Omega_e)$

$$\begin{aligned} \langle (\Lambda(q) - \Lambda(q_0)) F, F \rangle &\geq \int_{\Omega} (q - q_0) |u_q|^2 dx \geq \int_{\Omega} (\tilde{q} - q) |u_q|^2 dx \\ &\geq \langle (\Lambda(\tilde{q}) - \Lambda(q)) F, F \rangle \geq \int_{\Omega} (\tilde{q} - q) |u_{\tilde{q}}|^2 dx \quad (4.11) \\ &\geq \delta \int_M |u_{\tilde{q}}|^2 dx - \alpha \int_{\Omega \setminus M} |u_{\tilde{q}}|^2 dx, \end{aligned}$$

where $u_q = S_q(F)$, $u_{\tilde{q}} = S_{\tilde{q}}(F)$, and, for the last inequality, we assumed without loss of generality that $\alpha > 0$ is larger than $\|q - q_0\|_{L^\infty(\Omega)}$. For the last argument, note that the inequalities in (4.11) each hold on possibly different subspaces of finite codimension in $H_{\tilde{q}}(\Omega_e)$, so that V is obtained by taking the orthogonal complement of the intersection of all these spaces.

We also define

$$\tilde{q}_0 := \begin{cases} q - \alpha & \text{in } C, \\ q_0 & \text{in } \Omega \setminus C. \end{cases}$$

Since $\text{supp}(\tilde{q}_0 - q_0) \subseteq C$, we can apply Theorem 3.4 to obtain a finite dimensional subspace $V' \subseteq H_{\tilde{q}_0, q_0}(\Omega_e) = H_{\tilde{q}_0}(\Omega_e)$ (note that q_0 is non-resonant), and a constant $c > 0$, so that for all $F \in V'^\perp \subseteq H_{\tilde{q}_0}(\Omega_e)$

$$\langle \mathcal{T}_C F, F \rangle = \int_C |u_{q_0}|^2 dx \leq c \int_C |u_{\tilde{q}_0}|^2 dx,$$

where $u_{q_0} = S_{q_0}(F)$, $u_{\tilde{q}_0} = S_{\tilde{q}_0}(F)$. Hence, the second inequality in (4.8) implies that

$$c \int_C |u_{\tilde{q}_0}|^2 dx \geq \delta \int_M |u_{\tilde{q}}|^2 dx - \alpha \int_{\Omega \setminus M} |u_{\tilde{q}}|^2 dx \quad (4.12)$$

for all $F \in W^\perp \subseteq H_{\tilde{q}_0, \tilde{q}}(\Omega_e)$, where $W \subseteq H_{\tilde{q}_0, \tilde{q}}(\Omega_e)$ is a finite dimensional subspace. But $\text{supp}(\tilde{q} - \tilde{q}_0) \subseteq M$, so that the result on simultaneously localized potentials in Theorem 3.11 (with Theorem 3.11 applied to the herein constructed subspace W) yields the existence of a sequence $\{F^k\}_{k \in \mathbb{N}} \subseteq W^\perp \subseteq H_{\tilde{q}_0, \tilde{q}}(\Omega_e)$, so that the corresponding solutions $u_{\tilde{q}_0}^k = S_{\tilde{q}_0}(F^k)$, $u_{\tilde{q}}^k = S_{\tilde{q}}(F^k)$, fulfill

$$\int_M |u_{\tilde{q}}^k|^2 dx \rightarrow \infty, \quad \int_{\Omega \setminus M} |u_{\tilde{q}_0}^k|^2 dx \rightarrow 0, \quad \text{and} \quad \int_{\Omega \setminus M} |u_{\tilde{q}}^k|^2 dx \rightarrow 0,$$

which contradicts (4.12) since $C \subseteq \Omega \setminus M$. Hence, (4.10) and thus the assertion is proven.

□

We also extend the simpler results for the definite case, where either $q \geq q_0$ or $q \leq q_0$ holds almost everywhere in Ω , from [39] to general (but non-resonant) $L^\infty(\Omega)$ -potentials. We will show that it suffices to test open balls to reconstruct the inner support (for $q \geq q_0$), resp., a set between the support of $q - q_0$ and its inner support (for $q \leq q_0$), where, as in [47, Section 2.2], the inner support $\text{inn supp}(r)$ of a measurable function $r : \Omega \rightarrow \mathbb{R}$ is defined as the union of all open sets U on which the essential infimum of $|\kappa|$ is positive.

THEOREM 4.11.

(a) Let $q \leq q_0$. For every open set $B \subseteq \Omega$ and every $\alpha > 0$

$$q \leq q_0 - \alpha \chi_B \quad \text{implies} \quad \Lambda(q) \leq_{d(q)} \Lambda(q_0) - \alpha \mathcal{T}_B, \quad (4.13)$$

$$\Lambda(q) \leq_{\text{fin}} \Lambda(q_0) - \alpha \mathcal{T}_B \quad \text{implies} \quad B \subseteq \text{supp}(q - q_0). \quad (4.14)$$

Hence,

$$\begin{aligned} & \text{inn supp}(q - q_0) \\ & \subseteq \bigcup \{B \subseteq \Omega \text{ open ball} : \exists \alpha > 0 : \Lambda(q) \leq_{d(q)} \Lambda(q_0) - \alpha \mathcal{T}_B\} \\ & \subseteq \bigcup \{B \subseteq \Omega \text{ open ball} : \exists \alpha > 0 : \Lambda(q) \leq_{\text{fin}} \Lambda(q_0) - \alpha \mathcal{T}_B\} \\ & \subseteq \text{supp}(q - q_0). \end{aligned}$$

(b) Let $q \geq q_0$. For every open set $B \subseteq \Omega$ and every $\alpha > 0$

$$q \geq q_0 + \alpha \chi_B \quad \text{implies} \quad \exists \tilde{\alpha} > 0 : \Lambda(q) \geq_{\text{fin}} \Lambda(q_0) + \tilde{\alpha} \mathcal{T}_B, \quad (4.15)$$

$$\Lambda(q) \geq_{\text{fin}} \Lambda(q_0) + \alpha \mathcal{T}_B \quad \text{implies} \quad q \geq q_0 + \alpha \chi_B. \quad (4.16)$$

Hence,

$$\text{inn supp}(q - q_0) = \bigcup \{B \subseteq \Omega \text{ open ball} : \exists \alpha > 0 : \Lambda(q) \geq_{\text{fin}} \Lambda(q_0) + \alpha \mathcal{T}_B\}.$$

Proof.

(a) If $q_1 \leq q_0 - \alpha \chi_B$, then we obtain using Theorem 4.8, and Corollary 4.7 that

$$\Lambda(q) - \Lambda(q_0) \leq_{d(q)} \Lambda'(q_0)(q - q_0) \leq -\alpha \Lambda'(q_0) \chi_B = -\alpha \mathcal{T}_B,$$

so that (4.13) is proven. On the other hand, if $\Lambda(q) \leq_{\text{fin}} \Lambda(q_0) - \alpha\mathcal{T}_B$ then we obtain from Theorem 4.8, and Corollary 4.7, that there exists $c > 0$ with

$$\begin{aligned} \alpha\Lambda'(q_0)\chi_B &= \alpha\mathcal{T}_B \leq_{\text{fin}} \Lambda(q_0) - \Lambda(q) \leq_{\text{fin}} \Lambda'(q)(q_0 - q) \\ &\leq \|q_0 - q\|_{L^\infty(\Omega)} \Lambda'(q) \chi_{\text{supp}(q-q_0)} \\ &\leq_{\text{fin}} c \|q_0 - q\|_{L^\infty(\Omega)} \Lambda'(q_0) \chi_{\text{supp}(q-q_0)}, \end{aligned}$$

and that this implies

$$\alpha\chi_B \leq c \|q_0 - q\|_{L^\infty(\Omega)} \chi_{\text{supp}(q-q_0)},$$

so that (4.14) is proven.

- (b) Let $q \geq q_0 + \alpha\chi_B$. By Theorem 3.3, there exists a subspace $V \subseteq H_{q_0 + \alpha\chi_B}(\Omega_e)$ with $\dim(V) \leq d(q_0 + \alpha\chi_B)$ so that

$$\langle \Lambda(q)F, F \rangle \geq \langle \Lambda(q_0 + \alpha\chi_B)F, F \rangle \quad \text{for all } F \in V^\perp \subseteq H_{q_0 + \alpha\chi_B}(\Omega_e).$$

Moreover, by Theorem 3.4 there also exists a subspace $V' \subseteq H_{q_0 + \alpha\chi_B}(\Omega_e)$ with $\dim(V') \leq d(q_0)$ and a constant $c > 0$ so that

$$\begin{aligned} \langle (\Lambda(q_0 + \alpha\chi_B) - \Lambda(q_0))F, F \rangle &\geq \alpha \int_B |u_{q_0 + \alpha\chi_B}|^2 dx \\ &\geq \alpha c \int_B |u_{q_0}|^2 dx = c\alpha \langle \mathcal{T}_B F, F \rangle \end{aligned}$$

for all $F \in V'^\perp \subseteq H_{q_0 + \alpha\chi_B}$, where $u_{q_0 + \alpha\chi_B} = S_{q_0 + \alpha\chi_B}(F)$, and $u_{q_0} = S_{q_0}(F)$. Hence

$$\langle (\Lambda(q) - \Lambda(q_0))F, F \rangle \geq c\alpha\mathcal{T}_B,$$

holds for all $F \in (V + V')^\perp \subseteq H_{q_0 + \alpha\chi_B}(\Omega_e)$, which is a subspace of codimension $\dim(N_{q_0 + \alpha\chi_B})$ in $H(\Omega_e)$. Hence,

$$\Lambda(q) \geq_d \Lambda(q_0) + c\alpha\mathcal{T}_B \quad \text{with } d = d(q) + d(q_0 + \alpha\chi_B) + \dim(N_{q_0 + \alpha\chi_B}),$$

which shows (4.15). On the other hand, $\Lambda(q) \geq_{\text{fin}} \Lambda(q_0) + \alpha\mathcal{T}_B$ implies by Corollary 4.7

$$\alpha\Lambda'(q_0)\chi_B = \alpha\mathcal{T}_B \leq_{\text{fin}} \Lambda(q) - \Lambda(q_0) \leq_{\text{fin}} \Lambda'(q_0)(q - q_0),$$

so that it follows from Theorem 4.8 that

$$\alpha\chi_B \leq q - q_0,$$

which proves (4.15). \square

5. Uniqueness and Lipschitz stability for the fractional Calderón problem with finitely many measurements. In this section let $\mathcal{Q} \subseteq L^\infty(\Omega)$ be a finite dimensional subspace and, with a fixed constant $a > 0$, let

$$\mathcal{Q}_{[-a, a]} := \{q \in \mathcal{Q} : \|q\|_{L^\infty(\Omega)} \leq a\}.$$

We will show that a sufficiently high number of measurements of the DtN operator uniquely determines a potential in $\mathcal{Q}_{[-a,a]}$ and prove a Lipschitz stability result.

To formulate our result, we denote the orthogonal projection operators from $H(\Omega_e)$ to a subspace H by P_H , i.e. P_H is the linear operator with

$$P_H : H(\Omega_e) \rightarrow H, \quad P_H F := \begin{cases} F & \text{if } F \in H, \\ 0 & \text{if } F \in H^\perp \subseteq H(\Omega_e). \end{cases}$$

$P'_H : H^* \rightarrow H(\Omega_e)^*$ denotes the dual operator of P_H . For possibly resonant potentials $q_1, q_2 \in L^\infty(\Omega)$, the subspace H might contain non-admissible Dirichlet boundary values, so we also require the orthogonal projection $P_{q_1 q_2} := P_{H_{q_1, q_2}(\Omega_e)}$.

THEOREM 5.1. *For each sequence of subspaces*

$$H_1 \subseteq H_2 \subseteq H_3 \subseteq \dots \subseteq H(\Omega_e), \quad \text{with} \quad \overline{\bigcup_{l \in \mathbb{N}} H_l} = H(\Omega_e),$$

there exists $k \in \mathbb{N}$, and $c > 0$, so that

$$\|P'_{H_l} P'_{q_1 q_2} (\Lambda(q_2) - \Lambda(q_1)) P_{q_1 q_2} P_{H_l}\|_{\mathcal{L}(H(\Omega_e), H(\Omega_e)^*)} \geq \frac{1}{c} \|q_2 - q_1\|_{L^\infty(\Omega)} \quad (5.1)$$

for all $q_1, q_2 \in \mathcal{Q}_{[-a,a]}$ and all $l \geq k$.

Before we prove Theorem 5.1, let us briefly remark on its implications for some special cases.

REMARK 5.2. *Theorem 5.1 implies that there exists $c > 0$ so that*

$$\|\Lambda(q_2) - \Lambda(q_1)\|_{\mathcal{L}(H_{q_1, q_2}(\Omega_e), H(\Omega_e)^*)} \geq c \|q_2 - q_1\|_{L^\infty(\Omega)} \quad \text{for all } q_1, q_2 \in \mathcal{Q}_{[-a,a]}.$$

If $\{F_1, F_2, \dots\} \subseteq H(\Omega_e)$ is a set of Dirichlet values whose linear span is dense in $H(\Omega_e)$, then Theorem 5.1 implies that there exists $k \in \mathbb{N}$, so that every non-resonant $q \in \mathcal{Q}_{[-a,a]}$ is uniquely determined by the finitely many entries of the matrix

$$A(q) = (\langle \Lambda(q) F_i, F_j \rangle)_{i,j=1,\dots,k} \in \mathbb{R}^{k \times k}.$$

Moreover, if $\{F_1, F_2, \dots\}$ is an orthonormal (Schauder) basis of $H(\Omega_e)$, then there exists $k \in \mathbb{N}$, and $c > 0$, so that

$$\|A(q_2) - A(q_1)\|_2 \geq c \|q_2 - q_1\|_{L^\infty(\Omega)} \quad \text{for all non-resonant } q \in \mathcal{Q}_{[-a,a]},$$

where $\|A\|_2$ is the spectral norm of the matrix $A \in \mathbb{R}^{k \times k}$.

The general outline of the proof of Theorem 5.1 is as follows. In Lemma 5.3, we will derive a number of subsets $M_1, \dots, M_m \subseteq \Omega$, on which normalized potential differences can be estimated from above or below. Then we define for each of these sets a special potential $\hat{q}_j \in L^\infty(\Omega)$, which is large on M_j and small on $\Omega \setminus M_j$, and show (in Lemma 5.4) that certain energy terms for the solutions for an arbitrary $q \in L^\infty(\Omega)$ can always be estimated by solutions corresponding to these special potentials \hat{q}_j . Lemma 5.5 gives a bound on the maximal codimension of the subspaces arising from resonances, and Lemma 5.6 shows the existence of sufficiently many (depending on the maximal codimension) Dirichlet boundary values \hat{F}_{ij} to control the energy terms arising from the special potentials \hat{q}_j . The constant $c > 0$ of the Lipschitz stability

estimate (5.1) and the subspace index $k \in \mathbb{N}$ for Theorem 5.1, will be defined in Lemma 5.6 via the maximal norm of the finitely many Dirichlet values \widehat{F}_{ij} , and the possibility of sufficiently well approximating \widehat{F}_{ij} in H_k . Finally, we prove that Theorem 5.1 holds with these constants $c > 0$ and $k \in \mathbb{N}$.

Let us stress that this construction (the sets M_1, \dots, M_m , the finitely many special potentials \widehat{q}_j , the dimension bounds, the finitely many special Dirichlet data \widehat{F}_{ij} , and thus the constant $c > 0$ of (5.1), and the subspace index $k \in \mathbb{N}$) do only depend on the a-priori data $\mathcal{Q}_{[a,b]}$ and $\Omega \subseteq \mathbb{R}^n$.

To motivate the first lemma, let us note that a piecewise constant function on some partition of Ω with $L^\infty(\Omega)$ -norm equal to 1, must be either 1 or -1 on at least one of the subsets of the partition, which is a useful property for applying monotonicity estimates, cf., e.g., [41]. The following lemma generalizes this property to our arbitrary finite-dimensional subspace $\mathcal{Q} \subset L^\infty(\Omega)$.

LEMMA 5.3. *Let $\mathcal{Q}_1 := \{r \in \mathcal{Q} : \|r\|_{L^\infty(\Omega)} = 1\}$. There exists a family of measurable subsets M_1, \dots, M_m , $m \in \mathbb{N}$, with positive measure, so that for all $r \in \mathcal{Q}_1$, there exists $j \in \{1, \dots, m\}$ with either $r|_{M_j} \geq \frac{1}{2}$, or $r|_{M_j} \leq -\frac{1}{2}$. Hence, either*

$$r \geq \frac{1}{2}\chi_{M_j} - \chi_{\Omega \setminus M_j}, \quad \text{or} \quad -r \geq \frac{1}{2}\chi_{M_j} - \chi_{\Omega \setminus M_j}.$$

Proof. We argue by compactness. For $r \in \mathcal{Q}_1$, $\|r\|_{L^\infty(\Omega)} = 1$ implies that at least one of the sets $r^{-1}([\frac{1}{2}, \frac{3}{2}])$ or $r^{-1}([\frac{3}{2}, -\frac{1}{2}])$ must be of positive measure. In the first case we define

$$M_r := r^{-1}([\frac{1}{2}, \frac{3}{2}]), \quad \mathcal{O}_r := \{\tilde{r} \in L^\infty(\Omega) : \|\tilde{r}|_{M_r} - \chi_{M_r}\|_{L^\infty(M_r)} < \frac{1}{2}\},$$

and otherwise we define

$$M_r := r^{-1}([\frac{3}{2}, -\frac{1}{2}]), \quad \mathcal{O}_r := \{\tilde{r} \in L^\infty(\Omega) : \|\tilde{r}|_{M_r} + \chi_{M_r}\|_{L^\infty(M_r)} < \frac{1}{2}\}.$$

Then M_r has positive measure, \mathcal{O}_r is an open subset of $L^\infty(\Omega)$, and $r \in \mathcal{O}_r$ implies that

$$\mathcal{Q}_1 \subseteq \bigcup_{r \in \mathcal{Q}_1} \mathcal{O}_r.$$

By compactness, there exist $r_1, \dots, r_m \in \mathcal{Q}_1$ with $\mathcal{Q}_1 \subseteq \bigcup_{j=1, \dots, m} \mathcal{O}_{r_j}$, so that the assertion follows with $M_j := M_{r_j}$, $j = 1, \dots, m$. \square

We now use the idea from the constructive Lipschitz stability proof in [41, Section 5] to replace general potentials from $\mathcal{Q}_{[-a,a]}$ by a finite number of special potentials.

LEMMA 5.4. *With the constant $a > 0$ and the sets M_1, \dots, M_m from Lemma 5.3, we define*

$$\widehat{q}_j \in L^\infty(\Omega) \quad \text{by} \quad \widehat{q}_j := 2a\chi_{M_j} - 7a\chi_{\Omega \setminus M_j}, \quad j = 1, \dots, m.$$

If $q \in \mathcal{Q}_{[-a,a]}$ and $r \in \mathcal{Q}_1$ fulfills $r \geq \frac{1}{2}\chi_{M_j} - \chi_{\Omega \setminus M_j}$ with $j \in \{1, \dots, m\}$, then there exists a subspace $V \subseteq H_{q, \widehat{q}_j}(\Omega_e)$ with $\dim V \leq d(q) + d(\widehat{q}_j)$, so that

$$\int_{\Omega} r |S_q(F)|^2 dx \geq \int_{\Omega} \left(\frac{1}{6}\chi_{M_j} - \frac{4}{3}\chi_{\Omega \setminus M_j} \right) |S_{\widehat{q}_j}(F)|^2 dx \quad \text{for all } F \in V^\perp \subseteq H_{q, \widehat{q}_j}(\Omega_e).$$

Proof. Let $q \in \mathcal{Q}_{[-a,a]}$ and $r \in \mathcal{Q}_1$ fulfill $r \geq \frac{1}{2}\chi_{M_j} - \chi_{\Omega \setminus M_j}$ with $j \in \{1, \dots, m\}$. Then we obtain from Remark 3.5 a subspace $V \subseteq H_{q, \hat{q}_j}(\Omega_e)$ with $\dim V \leq d(q) + d(\hat{q}_j)$, so that for all $F \in V^\perp \subseteq H_{q, \hat{q}_j}(\Omega_e)$

$$\int_{\Omega} (\hat{q}_j - q) |S_{\hat{q}_j}(F)|^2 dx \leq \langle (\Lambda(\hat{q}_j) - \Lambda(q)) F, F \rangle \leq \int_{\Omega} (\hat{q}_j - q) |S_q(F)|^2 dx.$$

Observe that

$$a\chi_{M_j} - 8a\chi_{\Omega \setminus M_j} \leq \hat{q}_j - q \leq 3a\chi_{M_j} - 6a\chi_{\Omega \setminus M_j},$$

then it follows for all $F \in V^\perp \subseteq H_{q, \hat{q}_j}(\Omega_e)$

$$\begin{aligned} & \int_{\Omega} r |S_q(F)|^2 dx \geq \int_{\Omega} \left(\frac{1}{2}\chi_{M_j} - \chi_{\Omega \setminus M_j} \right) |S_q(F)|^2 dx \\ &= \frac{1}{6a} \int_{\Omega} (3a\chi_{M_j} - 6a\chi_{\Omega \setminus M_j}) |S_q(F)|^2 dx \geq \frac{1}{6a} \int_{\Omega} (\hat{q}_j - q) |S_q(F)|^2 dx \\ &\geq \frac{1}{6a} \int_{\Omega} (\hat{q}_j - q) |S_{\hat{q}_j}(F)|^2 dx \geq \frac{1}{6a} \int_{\Omega} (a\chi_{M_j} - 8a\chi_{\Omega \setminus M_j}) |S_{\hat{q}_j}(F)|^2 dx \\ &= \int_{\Omega} \left(\frac{1}{6}\chi_{M_j} - \frac{4}{3}\chi_{\Omega \setminus M_j} \right) |S_{\hat{q}_j}(F)|^2 dx. \end{aligned}$$

□

The next lemma shows that the codimension of the subspaces where the DtN operators are defined, and the subspaces where the monotonicity relations hold, can be uniformly bounded in $\mathcal{Q}_{[-a,a]}$.

LEMMA 5.5. *There exists numbers $d, N \in \mathbb{N}$, so that*

$$\dim(N_q) \leq N \quad \text{and} \quad d(q) \leq d \quad \text{for all } q \in \mathcal{Q}_{[-a,a]},$$

where N_q is defined by (2.5) and $d(q)$ is given by Definition 3.2.

Proof. The first assertion follows from Theorem 3.6(b) with a standard compactness argument. The second assertion follows from Theorem 3.6(a) with $d := d(-a)$, where $d(-a)$ is the number defined in Definition 3.2 for $q \equiv -a$. □

Our last lemma demonstrates how to control the energy terms in Lemma 5.4, and defines the Lipschitz stability constant $c > 0$ and the subspace index $k \in \mathbb{N}$, with which the assertion of Theorem 5.1 holds.

LEMMA 5.6. *Let $d, N \in \mathbb{N}$ be the numbers given in Lemma 5.5, then we have*

(a) *For all $j \in \{1, \dots, m\}$, there exist Dirichlet data $\hat{F}_{i,j} \in H_{\hat{q}_j}(\Omega_e)$ with*

$$\int_{\Omega} \left(\frac{1}{6}\chi_{M_j} - \frac{4}{3}\chi_{\Omega \setminus M_j} \right) |S_{\hat{q}_j}(\hat{F}_{i,j})|^2 dx \geq 2, \quad (5.2)$$

$$\int_{\Omega} \left(\frac{1}{6}\chi_{M_j} - \frac{4}{3}\chi_{\Omega \setminus M_j} \right) S_{\hat{q}_j}(\hat{F}_{i,j}) S_{\hat{q}_j}(\hat{F}_{i',j}) dx = 0, \quad (5.3)$$

$$\left(\hat{F}_{i,j}, \hat{F}_{i',j} \right)_{H(\Omega_e)} = 0, \quad (5.4)$$

for all $i, i' = 1, \dots, 3d + 2N + 1$ with $i' \neq i$. We set

$$c := 2 \max \left\{ \|\hat{F}_{i,j}\|_{H(\Omega_e)}^2 : j = 1, \dots, m, i = 1, \dots, 3d + 2N + 1 \right\}.$$

(b) For $\delta := \frac{1}{3d+2n+2}$, and for each sequence of subspaces

$$H_1 \subseteq H_2 \subseteq H_3 \subseteq \dots \subseteq H(\Omega_e), \quad \text{with} \quad \overline{\bigcup_{l \in \mathbb{N}} H_l} = H(\Omega_e),$$

there exists $k \in \mathbb{N}$, and $F_{i,j} \in H_k \cap H_{\hat{q}_j}(\Omega_e)$, so that

$$\int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) |S_{\hat{q}_j}(F_{i,j})|^2 dx \geq 2 - \delta, \quad (5.5)$$

$$\left| \int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) S_{\hat{q}_j}(F_{i,j}) S_{\hat{q}_j}(F_{i',j}) dx \right| \leq \delta, \quad (5.6)$$

$$\left| (F_{i,j}, F_{i',j})_{H(\Omega_e)} \right| \leq \frac{c}{2} \delta, \quad (5.7)$$

and $\|F_{i,j}\|_{H(\Omega_e)}^2 \leq (1 + \delta) \frac{c}{2}$ for all $j = 1, \dots, m$, and all $i, i' = 1, \dots, 3d + 2N + 1$ with $i' \neq i$.

(c) For all $j = 1, \dots, m$, all subspaces $V \subseteq H_{\hat{q}_j}(\Omega_e)$, with $\dim V \leq 3d + 2N$, contain an element $F_j \in V^\perp \cap H_k$ with

$$\int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) |S_{\hat{q}_j}(F_j)|^2 dx \geq 1, \quad \text{and} \quad \|F_j\|_{H(\Omega_e)}^2 \leq c.$$

Proof. Let $j \in \{1, \dots, m\}$.

(a) Theorem 3.10 yields that every subspace V^\perp of finite codimension in $H_{\hat{q}_j}(\Omega_e)$ contains F that fulfill the property (5.2). Hence, for $i = 1$, we can apply Theorem 3.10 on $H_{\hat{q}_j}(\Omega_e)$ to obtain $\hat{F}_{1,j}$, and for $i > 1$, we obtain $\hat{F}_{i,j}$ by applying Theorem 3.10 on the subspace

$$\left\{ \hat{F} \in H_{\hat{q}_j}(\Omega_e) : \int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) S_{\hat{q}_j}(\hat{F}) S_{\hat{q}_j}(\hat{F}_{i',j}) dx = 0, \right. \\ \left. \text{and } (\hat{F}, \hat{F}_{i',j})_{H(\Omega_e)} = 0 \text{ for all } i' \in \{1, \dots, i-1\} \right\},$$

which is obviously of finite codimension in $H_{\hat{q}_j}(\Omega_e)$, and this shows (5.3) and (5.4).

(b) From the finite codimension of $H_{\hat{q}_j}(\Omega_e)$ in $H(\Omega_e)$, we obtain that $\bigcup_{l \in \mathbb{N}} H_l \cap H_{\hat{q}_j}(\Omega_e)$ is dense in $H_{\hat{q}_j}(\Omega_e)$. Hence, the assertion (b) follows from the continuity of the solution operator $S_{\hat{q}_j}$.

(c) Since $V \subseteq H_{\hat{q}_j}(\Omega_e)$ has $\dim V \leq 3d + 2N$, there exists a non-trivial linear combination

$$0 \neq F_j := \sum_{i=1}^{3d+2N+1} \lambda_i F_{i,j} \in V^\perp \cap H_k, \quad \text{with coefficient } \lambda_i \in \mathbb{R},$$

where we normalize the coefficients so that $\sum_{i=1}^{3d+2N+1} |\lambda_i|^2 = 1$ and $k \in \mathbb{N}$ is the same number given as in (b). Then,

$$\sum_{i,i'=1}^{3d+2N+1} |\lambda_i| |\lambda_{i'}| \leq 3d + 2N + 1.$$

By using (5.5), (5.6) and (5.7), a simple calculation shows that

$$\int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) |S_{\widehat{q}_j}(F_j)|^2 dx \geq 2 - (3d + 2N + 2)\delta = 1,$$

$$\|F_j\|_{H(\Omega_e)}^2 \leq (1 + (3d + 2N + 2)\delta) \frac{c}{2} = c.$$

This completes the proof. \square

Now, we can prove Theorem 5.1.

Proof of Theorem 5.1. Let $q_1, q_2 \in \mathcal{Q}_{[-a, a]}$ with $q_1 \neq q_2$, and set $r := \frac{q_2 - q_1}{\|q_2 - q_1\|_{L^\infty(\Omega)}}$. Then, by Lemma 5.3, there exist $j \in \{1, \dots, m\}$ with either

$$(a) \ r \geq \frac{1}{2} \chi_{M_j} - \chi_{\Omega \setminus M_j}, \quad \text{or} \quad (b) \ -r \geq \frac{1}{2} \chi_{M_j} - \chi_{\Omega \setminus M_j}.$$

In case (a), Theorem 3.3 yields that there exists a subspace $V' \subseteq H_{q_1, q_2}(\Omega_e)$ of dimension $d(q_1)$, so that

$$\frac{\langle (\Lambda(q_2) - \Lambda(q_1)) F, F \rangle}{\|q_2 - q_1\|_{L^\infty(\Omega)}} \geq \int_{\Omega} r |S_{q_2}(F)|^2 dx \quad \text{for all } F \in (V')^\perp \subseteq H_{q_1, q_2}(\Omega_e).$$

Also, Lemma 5.4 yields a subspace $V'' \subseteq H_{q_2, \widehat{q}_j}(\Omega_e)$ with $\dim V'' \leq d(q_2) + d(\widehat{q}_j)$, so that

$$\int_{\Omega} r |S_{q_2}(F)|^2 dx \geq \int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) |S_{\widehat{q}_j}(F)|^2 dx \quad \forall F \in (V'')^\perp \subseteq H_{q_2, \widehat{q}_j}(\Omega_e).$$

Then $V := V' + V'' + H_{q_1}(\Omega_e)^\perp + H_{q_2}(\Omega_e)^\perp$ is a subspace with $\dim V \leq 3d + 2N$, and, by Lemma 5.6(c), there exists $F_j \in V^\perp \cap H_k$ with $\|F_j\|_{H(\Omega_e)}^2 \leq c$, and

$$\int_{\Omega} r |S_{q_2}(F_j)|^2 dx \geq \int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) |S_{\widehat{q}_j}(F_j)|^2 dx \geq 1.$$

Since $F_j \in V^\perp \cap H_k$, and the definition of V implies that $V^\perp \subseteq H_{\widehat{q}_j}(\Omega_e)$ is a subspace of $H_{q_1, q_2}(\Omega_e)$, we have that $P_{q_1 q_2} P_{H_l} F_j = P_{q_1 q_2} F_j = F_j$ for all $l \geq k$. Hence, it follows from the self-adjointness of $P'_{H_l} P'_{q_1 q_2} (\Lambda(q_2) - \Lambda(q_1)) P_{q_1 q_2} P_{H_l}$ that for all $l \geq k$,

$$\begin{aligned} & \frac{\|P'_{H_l} P'_{q_1 q_2} (\Lambda(q_2) - \Lambda(q_1)) P_{q_1 q_2} P_{H_l}\|_{\mathcal{L}(H(\Omega_e), H(\Omega_e)^*)}}{\|q_2 - q_1\|_{L^\infty(\Omega)}} \\ &= \sup_{0 \neq F \in H(\Omega_e)} \frac{| \langle (\Lambda(q_2) - \Lambda(q_1)) P_{q_1 q_2} P_{H_l} F, P_{q_1 q_2} P_{H_l} F \rangle |}{\|q_2 - q_1\|_{L^\infty(\Omega)} \|F\|_{H(\Omega_e)}^2} \\ &\geq \frac{| \langle (\Lambda(q_2) - \Lambda(q_1)) F_j, F_j \rangle |}{\|q_2 - q_1\|_{L^\infty(\Omega)} \|F_j\|_{H(\Omega_e)}^2} \geq \frac{1}{\|F_j\|_{H(\Omega_e)}^2} \int_{\Omega} r |S_{q_2}(F_j)|^2 dx \\ &\geq \frac{1}{\|F_j\|_{H(\Omega_e)}^2} \int_{\Omega} \left(\frac{1}{6} \chi_{M_j} - \frac{4}{3} \chi_{\Omega \setminus M_j} \right) |S_{\widehat{q}_j}(F_j)|^2 dx \geq \frac{1}{\|F_j\|_{H(\Omega_e)}^2} \geq \frac{1}{c}. \end{aligned}$$

In case (b), Theorem 3.3 yields that there exists a subspace $V' \subseteq H_{q_1, q_2}(\Omega_e)$ with dimension $d(q_2)$, so that

$$\frac{\langle (\Lambda(q_2) - \Lambda(q_1)) F, F \rangle}{\|q_2 - q_1\|_{L^\infty(\Omega)}} \geq - \int_{\Omega} r |S_{q_1}(F)|^2 dx \quad \text{for all } F \in (V')^\perp \subseteq H_{q_1, q_2}(\Omega_e),$$

so that the assertion follows analogously by using Lemma 5.4 with $-r$ instead of r . \square

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REFERENCES

- [1] G. Alessandrini. Singular solutions of elliptic equations and the determination of conductivity by boundary measurements. *J. Differential Equations*, 84(2):252–272, 1990.
- [2] G. Alessandrini, E. Beretta, and S. Vessella. Determining linear cracks by boundary measurements: Lipschitz stability. *SIAM Journal on Mathematical Analysis*, 27(2):361–375, 1996.
- [3] G. Alessandrini, M. V. de Hoop, R. Gaburro, and E. Sincich. Lipschitz stability for a piecewise linear Schrödinger potential from local Cauchy data. *Asymptotic Analysis*, 108(3):115–149, 2018.
- [4] G. Alessandrini, V. Maarten, R. Gaburro, and E. Sincich. Lipschitz stability for the electrostatic inverse boundary value problem with piecewise linear conductivities. *Journal de Mathématiques Pures et Appliquées*, 107(5):638–664, 2017.
- [5] G. Alessandrini and S. Vessella. Lipschitz stability for the inverse conductivity problem. *Advances in Applied Mathematics*, 35(2):207–241, 2005.
- [6] L. Arnold and B. Harrach. Unique shape detection in transient eddy current problems. *Inverse Problems*, 29(9):095004, 2013.
- [7] V. Bacchelli and S. Vessella. Lipschitz stability for a stationary 2D inverse problem with unknown polygonal boundary. *Inverse problems*, 22(5):1627, 2006.
- [8] A. Barth, B. Harrach, N. Hyvönen, and L. Mustonen. Detecting stochastic inclusions in electrical impedance tomography. *Inverse Problems*, 33(11):115012, 2017.
- [9] L. Beilina, M. Cristofol, S. Li, and M. Yamamoto. Lipschitz stability for an inverse hyperbolic problem of determining two coefficients by a finite number of observations. *Inverse Problems*, 34(1):015001, 2017.
- [10] M. Bellassoued, D. Jellali, and M. Yamamoto. Lipschitz stability for a hyperbolic inverse problem by finite local boundary data. *Applicable Analysis*, 85(10):1219–1243, 2006.
- [11] M. Bellassoued and M. Yamamoto. Lipschitz stability in determining density and two Lamé coefficients. *Journal of mathematical analysis and applications*, 329(2):1240–1259, 2007.
- [12] E. Beretta, M. V. de Hoop, E. Francini, S. Vessella, and J. Zhai. Uniqueness and Lipschitz stability of an inverse boundary value problem for time-harmonic elastic waves. *Inverse Problems*, 33(3):035013, 2017.
- [13] E. Beretta, M. V. De Hoop, and L. Qiu. Lipschitz stability of an inverse boundary value problem for a Schrödinger-type equation. *SIAM Journal on Mathematical Analysis*, 45(2):679–699, 2013.
- [14] E. Beretta and E. Francini. Lipschitz stability for the electrical impedance tomography problem: the complex case. *Communications in Partial Differential Equations*, 36(10):1723–1749, 2011.
- [15] T. Brander, B. Harrach, M. Kar, and M. Salo. Monotonicity and enclosure methods for the p -Laplace equation. *SIAM J. Appl. Math.*, 78(2):742–758, 2018.
- [16] X. Cao, Y.-H. Lin, and H. Liu. Simultaneously recovering potentials and embedded obstacles for anisotropic fractional Schrödinger operators. *Inverse Problems and Imaging*, 13(1):197–210, 2019.
- [17] X. Cao and H. Liu. Determining a fractional Helmholtz system with unknown source and medium parameter. *arXiv preprint arXiv:1803.09538*, 2018.
- [18] M. Cekić, Y.-H. Lin, and A. Rüländ. The Calderón problem for the fractional Schrödinger equation with drift. *arXiv preprint arXiv:1810.04211*, 2018.
- [19] J. Cheng, V. Isakov, M. Yamamoto, Q. Zhou, et al. Lipschitz stability in the lateral Cauchy problem for elasticity system. *Journal of Mathematics of Kyoto University*, 43(3):475–501, 2003.
- [20] E. Di Nezza, G. Palatucci, and E. Valdinoci. Hitchhiker’s guide to the fractional Sobolev spaces. *Bulletin des Sciences Mathématiques*, 136(5):521–573, 2012.
- [21] S. Eberle, B. Harrach, H. Meftahi, and T. Rezgui. Lipschitz stability estimate and reconstruction of Lamé parameters in linear elasticity. *arXiv preprint arXiv:1906.02194*, 2019.

- [22] L. C. Evans. *Partial differential equations*, volume 19 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, second edition, 2010.
- [23] F. Frühaufer, B. Gebauer, and O. Scherzer. Detecting interfaces in a parabolic-elliptic problem from surface measurements. *SIAM Journal on Numerical Analysis*, 45(2):810–836, 2007.
- [24] T. Furuya. The monotonicity based method for the inverse crack scattering problem. *arXiv preprint arXiv:1904.03655*, 2019.
- [25] H. Garde. Comparison of linear and non-linear monotonicity-based shape reconstruction using exact matrix characterizations. *Inverse Problems in Science and Engineering*, pages 1–18, 2017.
- [26] H. Garde. Reconstruction of piecewise constant layered conductivities in electrical impedance tomography. *arXiv preprint arXiv:1904.07775*, 2019.
- [27] H. Garde and S. Staboulis. Convergence and regularization for monotonicity-based shape reconstruction in electrical impedance tomography. *Numerische Mathematik*, 135(4):1221–1251, 2017.
- [28] H. Garde and S. Staboulis. The regularized monotonicity method: Detecting irregular indefinite inclusions. *Inverse Probl. Imaging*, 13(1):93–116, 2019.
- [29] B. Gebauer. Localized potentials in electrical impedance tomography. *Inverse Probl. Imaging*, 2(2):251–269, 2008.
- [30] T. Ghosh, Y.-H. Lin, and J. Xiao. The Calderón problem for variable coefficients nonlocal elliptic operators. *Communications in Partial Differential Equations*, 42(12):1923–1961, 2017.
- [31] T. Ghosh, A. Rüland, M. Salo, and G. Uhlmann. Uniqueness and reconstruction for the fractional Calderón problem with a single measurement. *arXiv preprint arXiv:1801.04449*, 2018.
- [32] T. Ghosh, M. Salo, and G. Uhlmann. The Calderón problem for the fractional Schrödinger equation. *Anal. PDE*, to appear.
- [33] R. Griesmaier and B. Harrach. Monotonicity in inverse medium scattering on unbounded domains. *SIAM J. Appl. Math.*, 78(5):2533–2557, 2018.
- [34] B. Harrach. On uniqueness in diffuse optical tomography. *Inverse Problems*, 25:055010 (14pp), 2009.
- [35] B. Harrach. Simultaneous determination of the diffusion and absorption coefficient from boundary data. *Inverse Probl. Imaging*, 6(4):663–679, 2012.
- [36] B. Harrach. On stable invertibility and global Newton convergence for convex monotonic functions. *arXiv preprint arXiv:1907.02759*, 2019.
- [37] B. Harrach. Uniqueness and Lipschitz stability in electrical impedance tomography with finitely many electrodes. *Inverse Problems*, 35(2):024005, 2019.
- [38] B. Harrach, E. Lee, and M. Ullrich. Combining frequency-difference and ultrasound modulated electrical impedance tomography. *Inverse Problems*, 31(9):095003, 2015.
- [39] B. Harrach and Y.-H. Lin. Monotonicity-based inversion of the fractional Schrödinger equation I. Positive potentials. *SIAM J. Math. Anal.*, to appear.
- [40] B. Harrach, Y.-H. Lin, and H. Liu. On localizing and concentrating electromagnetic fields. *SIAM J. Appl. Math.*, 78(5):2558–2574, 2018.
- [41] B. Harrach and H. Meftahi. Global uniqueness and Lipschitz-stability for the inverse Robin transmission problem. *SIAM J. Appl. Math.*, 79(2):525–550, 2019.
- [42] B. Harrach and M. N. Minh. Enhancing residual-based techniques with shape reconstruction features in electrical impedance tomography. *Inverse Problems*, 32(12):125002, 2016.
- [43] B. Harrach and M. N. Minh. Monotonicity-based regularization for phantom experiment data in electrical impedance tomography. In *New Trends in Parameter Identification for Mathematical Models*, pages 107–120. Springer, 2018.
- [44] B. Harrach, V. Pohjola, and M. Salo. Dimension bounds in monotonicity methods for the Helmholtz equation. *SIAM J. Math. Anal.*, to appear.
- [45] B. Harrach, V. Pohjola, and M. Salo. Monotonicity and local uniqueness for the Helmholtz equation. *Anal. PDE*, to appear.
- [46] B. Harrach and J. K. Seo. Exact shape-reconstruction by one-step linearization in electrical impedance tomography. *SIAM Journal on Mathematical Analysis*, 42(4):1505–1518, 2010.
- [47] B. Harrach and M. Ullrich. Monotonicity-based shape reconstruction in electrical impedance tomography. *SIAM Journal on Mathematical Analysis*, 45(6):3382–3403, 2013.
- [48] B. Harrach and M. Ullrich. Resolution guarantees in electrical impedance tomography. *IEEE Trans. Med. Imaging*, 34:1513–1521, 2015.
- [49] B. Harrach and M. Ullrich. Local uniqueness for an inverse boundary value problem with partial data. *Proceedings of the American Mathematical Society*, 145(3):1087–1095, 2017.
- [50] M. Ikehata. Size estimation of inclusion. *Journal of Inverse and Ill-Posed Problems*, 6(2):127–

- 140, 1998.
- [51] M. Ikehata. Identification of the shape of the inclusion having essentially bounded conductivity. *Journal of Inverse and Ill-Posed Problems*, 7(6):533–540, 1999.
 - [52] O. Y. Imanuvilov and M. Yamamoto. Lipschitz stability in inverse parabolic problems by the Carleman estimate. *Inverse problems*, 14(5):1229, 1998.
 - [53] O. Y. Imanuvilov and M. Yamamoto. Global Lipschitz stability in an inverse hyperbolic problem by interior observations. *Inverse problems*, 17(4):717, 2001.
 - [54] V. Isakov. On uniqueness of recovery of a discontinuous conductivity coefficient. *Comm. Pure Appl. Math.*, 41(7):865–877, 1988.
 - [55] H. Kang, J. K. Seo, and D. Sheen. The inverse conductivity problem with one measurement: stability and estimation of size. *SIAM Journal on Mathematical Analysis*, 28(6):1389–1405, 1997.
 - [56] M. A. Kazemi and M. V. Klibanov. Stability estimates for ill-posed Cauchy problems involving hyperbolic equations and inequalities. *Applicable Analysis*, 50(1-2):93–102, 1993.
 - [57] M. V. Klibanov and S. E. Pamyatnykh. Lipschitz stability of a non-standard problem for the non-stationary transport equation via a Carleman estimate. *Inverse Problems*, 22(3):881, 2006.
 - [58] M. V. Klibanov and M. Yamamoto. Lipschitz stability of an inverse problem for an acoustic equation. *Applicable Analysis*, 85(05):515–538, 2006.
 - [59] R. V. Kohn and M. Vogelius. Determining conductivity by boundary measurements. *Communications on Pure and Applied Mathematics*, 37(3):289–298, 1984.
 - [60] R. V. Kohn and M. Vogelius. Determining conductivity by boundary measurements II. Interior results. *Communications on Pure and Applied Mathematics*, 38(5):643–667, 1985.
 - [61] R.-Y. Lai and Y.-H. Lin. Global uniqueness for the fractional semilinear Schrödinger equation. *Proceedings of the American Mathematical Society*, 2018.
 - [62] R.-Y. Lai, Y.-H. Lin, and A. Rüländ. The Calderón problem for a space-time fractional parabolic equation. *arXiv preprint arXiv:1905.08719*, 2019.
 - [63] P. D. Lax. *Functional analysis*. Pure and Applied Mathematics (New York). Wiley-Interscience [John Wiley & Sons], New York, 2002.
 - [64] A. Maffucci, A. Vento, S. Ventre, and A. Tamburrino. A novel technique for evaluating the effective permittivity of inhomogeneous interconnects based on the monotonicity property. *IEEE Transactions on Components, Packaging and Manufacturing Technology*, 6(9):1417–1427, 2016.
 - [65] P. G. Meléndez. Lipschitz stability in an inverse problem for the main coefficient of a Kuramoto–Sivashinsky type equation. *Journal of Mathematical Analysis and Applications*, 408(1):275–290, 2013.
 - [66] G. Palatucci, O. Savin, and E. Valdinoci. Local and global minimizers for a variational energy involving a fractional norm. *Annali di Matematica Pura ed Applicata*, 192(4):673–718, 2013.
 - [67] H. L. Royden and P. Fitzpatrick. *Real analysis*, volume 32. Macmillan New York, 1988.
 - [68] W. Rudin. *Functional analysis*. International Series in Pure and Applied Mathematics. McGraw-Hill, Inc., New York, second edition, 1991.
 - [69] A. Rüländ and M. Salo. The fractional Calderón problem: lower regularity and stability. *arXiv preprint arXiv:1708.06294*, 2017.
 - [70] A. Rüländ and M. Salo. Exponential instability in the fractional Calderón problem. *Inverse Problems*, 34(4):045003, 2018.
 - [71] A. Rüländ and E. Sincich. Lipschitz stability for the finite dimensional fractional Calderón problem with finite cauchy data. *arXiv preprint arXiv:1805.00866*, 2018.
 - [72] J. K. Seo, K. C. Kim, A. Jargal, K. Lee, and B. Harrach. A learning-based method for solving ill-posed nonlinear inverse problems: a simulation study of lung EIT. *SIAM J. Imaging Sci.*, to appear.
 - [73] E. Sincich. Lipschitz stability for the inverse Robin problem. *Inverse Problems*, 23(3):1311, 2007.
 - [74] Z. Su, L. Udpa, G. Giovenco, S. Ventre, and A. Tamburrino. Monotonicity principle in pulsed eddy current testing and its application to defect sizing. In *Applied Computational Electromagnetics Society Symposium-Italy (ACES), 2017 International*, pages 1–2. IEEE, 2017.
 - [75] A. Tamburrino and G. Rubinacci. A new non-iterative inversion method for electrical resistance tomography. *Inverse Problems*, 18(6):1809, 2002.
 - [76] A. Tamburrino, Z. Sua, S. Ventre, L. Udpa, and S. S. Udpa. Monotonicity based imaging method in time domain eddy current testing. *Electromagnetic Nondestructive Evaluation (XIX)*, 41:1, 2016.
 - [77] S. Ventre, A. Maffucci, F. Caire, N. Le Lostec, A. Perrotta, G. Rubinacci, B. Sartre, A. Vento,

- and A. Tamburrino. Design of a real-time eddy current tomography system. *IEEE Transactions on Magnetics*, 53(3):1–8, 2017.
- [78] G. Yuan and M. Yamamoto. Lipschitz stability in inverse problems for a Kirchhoff plate equation. *Asymptotic Analysis*, 53(1, 2):29–60, 2007.
- [79] G. Yuan and M. Yamamoto. Lipschitz stability in the determination of the principal part of a parabolic equation. *ESAIM: Control, Optimisation and Calculus of Variations*, 15(3):525–554, 2009.
- [80] L. Zhou, B. Harrach, and J. K. Seo. Monotonicity-based electrical impedance tomography for lung imaging. *Inverse Problems*, 34(4):045005, 2018.