

New stability results for inverse coefficient problems

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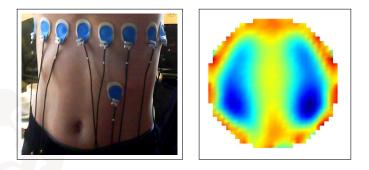
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Electrical impedance tomography



- Apply electric currents on subject's boundary
- Measure necessary voltages
- → Reconstruct conductivity inside subject



Calderón problem

Can we recover $\sigma \in L^\infty_+(\Omega)$ in

$$\nabla \cdot (\boldsymbol{\sigma} \nabla \boldsymbol{u}) = 0, \quad \boldsymbol{x} \in \boldsymbol{\Omega} \subset \mathbb{R}^d \qquad (1)$$

from all possible Dirichlet and Neumann boundary values

 $\{(u|_{\partial\Omega}, \sigma\partial_{\nu}u|_{\partial\Omega}) : u \text{ solves (1)}\}?$

Equivalent: Recover σ from Neumann-to-Dirichlet-Operator

 $\Lambda(\sigma): L^2_\diamond(\partial\Omega) \to L^2_\diamond(\partial\Omega), \quad g \mapsto u|_{\partial\Omega},$

where *u* solves (1) with $\sigma \partial_v u |_{\partial \Omega} = g$.

Challenges in idealized EIT



Mathematical idealization of EIT ~> Calderón problem

- infinitely many unknowns $\sigma \in L^{\infty}_{+}(\Omega)$
- infinitely many measurements $\Lambda(\sigma) \in \mathcal{L}(L^2_{\diamond}(\partial \Omega))$
- nonlinear forward map $\sigma \mapsto \Lambda(\sigma)$

Mathematical challenges

- Uniqueness? Does $\Lambda(\sigma)$ determine σ ?
- Stability? $\Lambda^{-1} \colon \Lambda(\sigma) \mapsto \sigma$ continuous?
- Convergence (local/global)? How to determine σ from $\Lambda(\sigma)$?

Consequences for practical EIT?



EIT in practice

In practice

- finitely many unknowns, σ pcw. const. on given resolution $\Omega = \bigcup_{i=1}^{m} \omega_i$
- finitely many measurements
- → Finite-dimensional inverse problem

Model for finitely many measurements:

Galerkin projection $P_{G_n}\Lambda(\sigma)P_{G_n}$. P_{G_n} : orthoprojection to

$$G_1 \subseteq G_2 \subseteq \ldots \subseteq L^2_\diamond(\partial \Omega), \quad \overline{\bigcup_{n \in \mathbb{N}} G_n} = L^2_\diamond(\partial \Omega)$$

Better: use realistic electrode model



For a fixed desired resolution:

- Do finitely many measurements uniquely determine σ?
 - (... and how many measurements/electrodes do we need?)
- Is the resulting finite-dimensional inverse problem stable? (... and how large is the stability constant / noise amplification?)
- Do the results hold for realistic electrode models?
 - (... and how can we derive globally convergent reconstruction algorithms?)

This talk: Affirmative answer to these challenges (... and some handwaving comments)

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Uniqueness and stability

Given desired resolution $\Omega = \bigcup_{i=1}^{m} \omega_i, b > a > 0$

 $\mathcal{F}_{[a,b]} = \{ \sigma \in L^{\infty}_{+}(\Omega) : a \leq \sigma(x) \leq b, \sigma \text{ pcw. const. on } \Omega \}$

and subspaces

$$G_1 \subseteq G_2 \subseteq \ldots \subseteq L^2_\diamond(\partial \Omega), \quad \overline{\bigcup_{n \in \mathbb{N}} G_n} = L^2_\diamond(\partial \Omega).$$

Theorem. (H., IP 2019) There exists $N \in \mathbb{N}$ and c > 0:

$$\|P_{G_n}(\Lambda(\sigma_1)-\Lambda(\sigma_2))P_{G_n}\| \geq c \|\sigma_1-\sigma_2\| \quad \forall \sigma_1,\sigma_2 \in \mathcal{F}_{[a,b]}, n \geq N.$$

Finitely many measurement uniquely determine σ at a given resolution if enough measurements are being used



Main tools for proof

Monotonicity lemma. (Kang/Seo/Sheen 1997, Ikehata 1998) For all $\sigma_1, \sigma_2 \in L^{\infty}_+(\Omega), g \in L^2_{\diamond}(\partial \Omega)$

$$\int_{\partial\Omega} g(\Lambda(\sigma_1) - \Lambda(\sigma_2))g \, \mathrm{d}s \ge \int_{\partial\Omega} g\Lambda'(\sigma_2)(\sigma_1 - \sigma_2)g \, \mathrm{d}s$$
$$= \int_{\Omega} (\sigma_2 - \sigma_1) |\nabla u_{\sigma_2}^g|^2 \, \mathrm{d}x$$

Localized potentials lemma. (H. 2008, H./Ullrich 2013) For pcw. anal. $\sigma \in L^{\infty}_{+}(\Omega)$, measurable $D_1, D_2 \subseteq \overline{\Omega}$, $\operatorname{int} D_1 \notin \operatorname{out}_{\partial\Omega} D_2$

$$\exists (g_k)_{k\in\mathbb{N}} \in L^2_{\diamond}(\partial\Omega): \ \int_{D_1} |\nabla u^{g_k}_{\sigma}|^2 \, \mathrm{d}x \to \infty, \quad \int_{D_2} |\nabla u^{g_k}_{\sigma}|^2 \, \mathrm{d}x \to 0.$$

(Closed outer hull $out_{\partial\Omega}D_2$: complement of all open sets connected to $\partial\Omega$ not intersecting D_2)

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Sketch of proof 1/2

By self-adjointness

$$\|\Lambda(\sigma_1) - \Lambda(\sigma_2)\| = \sup_{\|g\|=1} \left| \int_{\partial\Omega} g(\Lambda(\sigma_1) - \Lambda(\sigma_2)) g \, \mathrm{d}s \right|$$

By monotonicity

$$\begin{aligned} \left| \int_{\partial\Omega} g(\Lambda(\sigma_1) - \Lambda(\sigma_2)) g \, ds \right| \\ &= \max \left\{ \int_{\partial\Omega} g(\Lambda(\sigma_1) - \Lambda(\sigma_2)) g \, ds, \ \int_{\partial\Omega} g(\Lambda(\sigma_2) - \Lambda(\sigma_1)) g \, ds \right\} \\ &\geq \max \left\{ \int_{\partial\Omega} g\Lambda'(\sigma_2) (\sigma_1 - \sigma_2) g \, ds, \ \int_{\partial\Omega} g\Lambda'(\sigma_1) (\sigma_2 - \sigma_1) g \, ds \right\} \\ &= \|\sigma_1 - \sigma_2\| f\left(\sigma_1, \sigma_2, \frac{\sigma_1 - \sigma_2}{\|\sigma_1 - \sigma_2\|}, g\right) \end{aligned}$$

with $f(\tau_1, \tau_2, \kappa, g) \coloneqq \max \left\{ \int_{\partial \Omega} g \Lambda'(\tau_1) \kappa g \, \mathrm{d}s, - \int_{\partial \Omega} g \Lambda'(\tau_2)(\kappa) g \, \mathrm{d}s \right\}$



Sketch of proof 2/2

By last slide

$$\frac{\|\Lambda(\sigma_1) - \Lambda(\sigma_2)\|}{\|\sigma_1 - \sigma_2\|} \ge \sup_{\|g\|=1} f\left(\sigma_1, \sigma_2, \frac{\sigma_1 - \sigma_2}{\|\sigma_1 - \sigma_2\|}, g\right)$$
$$\ge \inf_{\tau_1, \tau_2, \kappa} \sup_{\|g\|=1} f(\tau_1, \tau_2, \kappa, g)$$

with infimum taken over compact set of all

 τ_1, τ_2, κ pcw. const., $\tau_1(x), \tau_2(x) \in [a, b], \|\kappa\| = 1$

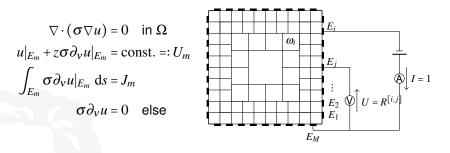
- f continuous → sup f l.s.c. → infimum is attained
- ► Localized potentials $\neg \forall \tau_1, \tau_2, \kappa : \exists g : f(\tau_1, \tau_2, \kappa, g) > 0$

$$\Rightarrow \exists c > 0: \|\Lambda(\sigma_1) - \Lambda(\sigma_2)\| \ge c \|\sigma_1 - \sigma_2\|.$$

(and Galerkin projection can be treated by another compactness argument)



Complete Electrode Model



Current-to-Voltage operator

$$R_M(\sigma): \mathbb{R}^M_\diamond \to \mathbb{R}^M_\diamond, \quad (J_1, \dots, J_M) \mapsto (U_1, \dots, U_M).$$

Uniqueness and stability (for enough electrodes)?

Uniqueness and Lipschitz-stability for fixed resolution

Assumptions:

- Increasing number of electrodes fulfills Hyvönen conditions
- *F*: finite-dimensional subset of pcw.-analytic functions
 (e.g., pcw. constant on fixed a-priori known partition)
- Known background conductivity: $\exists U$ nbr.hood of $\partial \Omega$, $\sigma_0 \in C^{\infty}$, so that $\sigma|_U = \sigma_0|_U$ for all $\sigma \in \mathcal{F}$
- A-prior known bounds

$$\mathcal{F}_{[a,b]} \coloneqq \{ \sigma \in \mathcal{F} \colon a \leq \sigma(x) \leq b \text{ for all } x \in \Omega \}$$

Theorem. (H, IP 2019) $\exists N \in \mathbb{N}, c > 0$:

$$\|R_M(\sigma_1)-R_M(\sigma_2)\|_{\mathcal{L}(\mathbb{R}^M_\diamond)} \ge c \|\sigma_1-\sigma_2\|_{L^\infty(\Omega)} \quad \forall \sigma_1,\sigma_2 \in \mathcal{F}_{[a,b]}, M \ge N.$$



EIT with fixed resolution is uniquely and stably solvable if enough electrodes are being used.

- EIT's ill-posedness due to inf.-dimens., not due to non-linearity
- Stability gets worse (exponentially) for finer resolution (For full NtD: Alessandrini/Vessella 2005, Rondi 2006)

Open questions:

- How many electrodes are required for a desired resolution?
- How good is the stability (error-amplification) in a given setting?
- Globally convergent solvers the discretized non-linear problem?
- Consequences of conductivity discretization?



Stability can be proven by monotonicity and localized potentials

Advantages:

- Simple: No analytic construction of special solutions required.
- Flexible: Method already applied to show stability for
 - Robin coefficient problem (H./Meftahi, SIAP 2019)
 - Deep learning approach to EIT (*Seo/Kim/Jargal/Lee/H. SIIMS, to appear*)
 - Fractional Calderón problem (*H./Lin, arXiv:1903.08771*)
- Constructive (possibly):
 - In Robin coeff. problem, Lipschitz constant for given resolution can be calculated by solving finitely many well-posed PDEs
 - Identifying necessary meas. for desired resol. seems in reach