MONOTONICITY AND ENCLOSURE METHODS FOR THE p-LAPLACE EQUATION*

TOMMI BRANDER†, BASTIAN HARRACH‡, MANAS KAR§, AND MIKKO SALO¶

Abstract. We show that the convex hull of a monotone perturbation of a homogeneous background conductivity in the p-conductivity equation is determined by knowledge of the nonlinear Dirichlet–Neumann operator. We give two independent proofs: one is based on the monotonicity method and the other on the enclosure method. Our results are constructive and require no jump or smoothness properties on the conductivity perturbation or its support.

Key words. p-Laplace equation, inclusion detection, monotonicity method, enclosure method

AMS subject classification. 35R30

DOI. 10.1137/17M1128599

1. Introduction. We consider the shape reconstruction problem for nonlinear partial differential equations of p-Laplace type. More precisely, we are interested in identifying the shape and location of an unknown inclusion from boundary voltage to current measurements. Assume $\Omega \subset \mathbb{R}^n, n \geq 2$, to be a bounded open set with conductivity $\sigma \in L_+^\infty(\Omega)$, where

$$L_+^\infty(\Omega) = \{f \in L^\infty(\Omega) \, ; \, \text{ess inf} \, f > 0\}.$$

We assume that the conductivity is constant (taken to be 1 for simplicity) outside an unknown inclusion D, so that the inclusion is the set

$$D = \text{supp}(\sigma - 1).$$

We study the problem of detecting inclusions from boundary measurements for the p-conductivity equation

$$
(1.1) \qquad \begin{cases} \text{div}(\sigma(x) \, |\nabla u|^{p-2} \, \nabla u) = 0 \text{ in } \Omega, \\ u = f \text{ on } \partial\Omega, \end{cases}
$$

where the exponent p is in the range $1 < p < \infty$. Given a Dirichlet boundary condition $f \in W^{1,p}(\Omega)/W_0^{1,p}(\Omega)$, the forward problem (1.1) is well-posed in $W^{1,p}(\Omega)$, and the weak solution minimizes the energy functional

$$E_\sigma(v) := \int_\Omega \sigma \, |\nabla v|^p \, dx$$

†Department of Applied Mathematics and Computer Science, Technical University of Denmark, Copenhagen, Denmark (tobr@dtu.dk).

‡Department of Mathematics, Goethe University Frankfurt, Frankfurt 60325, Germany (harrach@ math.uni-frankfurt.de).

§NCTS–National Center for Theoretical Sciences, National Taiwan University, Taipei City 106, Taiwan (manas.kar@ncts.ntu.edu.tw).

¶Department of Mathematics and Statistics, University of Jyväskylä, Jyväskylä 40014, Finland (mikko.j.salo@jyu.fi).

over all $v \in W^{1,p}(\Omega)$ with $v - f \in W_0^{1,p}(\Omega)$ (see, e.g., [35, 5]).

The boundary voltage to current map, also called the nonlinear Dirichlet–Neumann (DN) map, is the map

$$\Lambda_\sigma : X \to X',$$

defined formally by

$$\Lambda_\sigma(f) := \sigma \, |\nabla u|^{p-2} \, \nabla u \cdot \nu|_{\partial\Omega},$$

where $X := W^{1,p}(\Omega)/W_0^{1,p}(\Omega)$, X' is the dual of X, and ν is the unit outward normal to $\partial\Omega$. See section 2.2 for the precise weak definition of Λ_σ.

In the special case $p = 2$, (1.1) is the well-known conductivity equation appearing in Calderón's inverse problem [7]. This problem has been studied extensively in the last 35 years; see the survey [42] for recent results. Our problem is an analogue of the standard Calderón problem for nonlinear p-Laplace–type equations. These appear as models in various physical phenomena, e.g., nonlinear dielectrics, plastic moulding, electro-rheological and thermo-rheological fluids, fluids governed by a power law, viscous flows in glaciology, or plasticity. Also, the 0- and 1-Laplacians have applications in ultrasound modulated electrical impedance tomography (UMEIT) and current density imaging (CDI). See the references in [6] and [4, section 2.1] for more details.

Our purpose is to detect the shape and location of the inclusion D from boundary measurements for the p-conductivity equation (1.1), as encoded by the DN map Λ_σ. More precisely, we will be able to reconstruct the essential convex hull of D, assuming that $\sigma - 1$ does not change sign. Such a result was proved in [6] for inclusions having some regularity and sharp jumps at the interface by extending the enclosure method of Ikehata [29, 27] to this nonlinear model. In this paper we remove all regularity and interface jump assumptions, and also show that monotonicity-based shape reconstruction methods [39, 38, 23] work in the nonlinear case and allow us to find the convex hull of the inclusion.

The following theorem is the main result of this paper. It does not require any regularity or jump properties for the inclusion, and we obtain this result using both the monotonicity and the enclosure method. See section 2 for the definition of the essential convex hull.

THEOREM 1.1. *Suppose $\Omega \subset \mathbb{R}^n$ is a bounded and open set, $n \geq 2$, and $1 < p < \infty$. Consider a conductivity $\sigma \in L_+^\infty(\Omega)$ such that either $\sigma \geq 1$ almost everywhere or $\sigma \leq 1$ almost everywhere. Then we can recover the essential convex hull of $\operatorname{supp}(\sigma-1)$ from the DN map Λ_σ, and we can also determine whether σ is less than or greater than 1 almost everywhere.*

In addition, we give an alternative proof for the boundary determination result of Salo and Zhong [35, Theorem 1.1]. We use an enclosure-type method to recover a continuous conductivity on the boundary of a strictly convex domain, whereas [35] considers domains with C^1 boundary with no convexity assumptions.

THEOREM 1.2. *Let $\Omega \subset \mathbb{R}^n$ be bounded, open, and strictly convex. Let σ_1, σ_2 be positive continuous functions on $\overline{\Omega}$. If $\Lambda_{\sigma_1} = \Lambda_{\sigma_2}$, then $\sigma_1|_{\partial\Omega} = \sigma_2|_{\partial\Omega}$.*

We now give some additional details and introductory remarks on the monotonicity and the enclosure methods. The monotonicity method is based on the monotonicity relation

$$(1.2) \qquad\qquad \sigma_1 \geq \sigma_0 \implies \Lambda_{\sigma_1} \geq \Lambda_{\sigma_0}.$$

On the left-hand side of (1.2), $\sigma_1 \geq \sigma_0$ is to be understood pointwise almost everywhere. In the case $p = 2$, the DN maps are linear, and $\Lambda_{\sigma_1} \geq \Lambda_{\sigma_0}$ can be understood in the sense of operator definiteness (the Loewner partial order). Monotonicity relations such as (1.2) allow us to constructively determine inclusions by choosing a small test set B and a contrast level α and by checking whether the DN map for the test conductivity $1 + \alpha\chi_B$ is larger or smaller than Λ_σ. This approach has been proposed and numerically tested by Tamburrino and Rubinacci [38, 39]. To show that test sets B outside the true inclusion will not give false positive results in the monotonicity method, one requires a nontrivial converse of the implication (1.2), which has been shown by Harrach and Ullrich [23] using the concept of localized potentials [13]. Monotonicity-based reconstruction methods have been intensively studied and successfully applied to simulated and real measurement data for the $p = 2$ case [17, 19, 10, 11, 12, 20] and to other imaging modalities [18, 37, 40, 31, 41]. Monotonicity-based arguments have also been used to prove theoretical uniqueness results [15, 22, 16, 1, 24, 2, 25, 21].

In this work, we utilize that for $p \neq 2$, the monotonicity relation (1.2) is still valid when $\Lambda_{\sigma_1} \geq \Lambda_{\sigma_0}$ is understood in the sense of a preorder defined by the associated quadratic forms; see Lemma 2.9 and the beginning of section 3. Using special Wolff solutions we can then prove a converse of the implication (1.2), showing that the union of measure theoretic interiors of balls $B \in \mathcal{B}$ marked by the monotonicity method generates the essential convex hull of the inclusion,

$$\text{conv}(\text{supp}(|\sigma - 1|)) = \text{conv}\left(\bigcup_{B \in \mathcal{B}} B^\circ\right).$$

See Theorem 3.1 and Corollary 3.2. Reconstructing a piecewise real-analytic conductivity with the monotonicity method requires the ability to focus the energy of the solution to an arbitrary subset, as done, for example, by Ikehata [28] and Harrach [13]. These tools, and in particular the Runge approximation property, are not yet known for the p-Laplace equation.

The enclosure method uses specific exponential solutions, or complex geometrical optics (CGO)-type solutions, for the conductivity equation as test functions to detect the convex hull of the inclusion D. This method is based on analyzing the behavior of the indicator function, defined via the DN map and the exponential solutions, to see whether or not the level set of the phase function touches the boundary of the inclusion. The indicator function is given by

$$I(t, \tau) = \tau^{-p} \int_{\partial\Omega} (\Lambda_\sigma - \Lambda_1)(f_\tau) f_\tau \, dS,$$

where t and τ are parameters and f_τ are boundary values of the exponential Wolff solutions given in section 2. Now consider a half-space $H \subset \mathbb{R}^n$ so that the energy of the Wolff solution concentrates in H when the parameter τ becomes large. Then we observe the following facts for the indicator function. For large τ, if H does not meet D in a set of positive measure, then

$$I(t, \tau) \to 0,$$

and if H meets D in a set of positive measure, then we have

$$|I(t, \tau)| \to \infty.$$

Many earlier works on the enclosure method consider hyperplanes which barely touch the inclusion D, whereas in this work we consider the situations where, roughly, either $H \cap D$ has positive measure or dist $(H, D) > 0$.

We prove our main result, Theorem 1.1, with both the monotonicity and the enclosure methods. The enclosure method for linear equations allows one to identify various shapes by using CGO solutions with different phase functions. See, for instance, [26, 30, 36], which concern the use of linear phase functions, and [33], where spherical phase functions are used to reconstruct nonconvex parts of the obstacle. The work [32] uses the enclosure method with CGO solutions with polynomial phases, where the energy is concentrated inside a cone, and in this case it is possible to approximate the exact shape of certain types of obstacles. In [9], the authors propose a set of CGO solutions for the linear Schrödinger equation with all possible different phases. So, using all of these different phases, it might be possible to approximate the obstacle up to some obstructions. Therefore, the enclosure method heavily depends on the CGO solutions and their blow-up properties on a specific region, depending on the phase functions we choose.

On the other hand, monotonicity-based shape reconstruction methods depend on constructing solutions which blow up in suitable regions. Harrach [13] was able to produce such solutions for the linear conductivity equation. These solutions are known as localized potentials, and they can be used to determine the exact shape of the inclusion also in certain cases where the inclusion has indefinite sign. However, the existence of such solutions is proved by linear functional analysis (this involves a duality argument and the unique continuation principle), and it is not known whether localized potentials exist in the present nonlinear model. We will replace the localized potentials by Wolff-type solutions for the p-Laplace equation. These solutions will have very large energy on one side of a given hyperplane, with very small energy on the other side. We will use this phenomenon to establish a version of the monotonicity method in the nonlinear case.

In the present work we implement these two methods to reconstruct the essential convex hull of the inclusion. This is the first application of monotonicity-based shape reconstruction methods to nonlinear equations; see Theorem 3.1. We reach the same conclusion with the enclosure method; see Theorem 4.1. In contrast to many previous works, we do not assume any regularity assumptions on the boundary of the inclusion or any jump conditions for the conductivity.

Unlike for the linear case, comparatively little is known for inverse problems related to the p-Laplace equation. The first boundary determination result is due to Salo and Zhong [35], and boundary determination for the normal derivative of the conductivity was shown by Brander [3]. Recently, under monotonicity assumptions on the conductivity, an interior uniqueness result has been given by Guo, Kar, and Salo [14]. This result is not constructive, but it is not restricted to constant background conductivity. Brander, Kar, and Salo [6] detect the convex hull of an inclusion D when the conductivity σ satisfies $\sigma = 1$ in $\Omega \setminus \overline{D}$ and $\sigma \geq 1 + \varepsilon > 1$ in D or $\sigma < 1 - \varepsilon$ in D. In particular, the conductivity σ has a jump discontinuity along the interface ∂D. The work [5] considers the inclusion detection problem for models with zero or infinite conductivity. For an introduction to the p-Calderón problem, see the thesis of Brander [4].

This paper is organized as follows. In section 2, we discuss the essential convex hull and Wolff solutions required for the main results. Section 3 introduces the monotonicity-based shape reconstruction method for the p-Laplace equation. In section 4, we justify the enclosure method under the present assumptions. Finally, we

establish the boundary determination result for strictly convex domains in section 5.

2. Notation and preliminaries. Throughout this work let $\Omega \subset \mathbb{R}^n$ be open and bounded and suppose that the conductivity satisfies $\sigma \in L_+^\infty(\Omega; \mathbb{R})$, where $L_+^\infty(\Omega; \mathbb{R}) := \{f \in L^\infty(\Omega; \mathbb{R}); \operatorname{ess\,inf} f > 0\}$. We can assume that Ω is connected without loss of generality.

2.1. Essential support and convexity. We summarize some basic facts on the (essential) support and its (closed essential) convex hull. Here and in the following, measure theoretical terms (such as measurable, almost everywhere, or null set) are always used with respect to the n-dimensional Lebesgue measure, which we denote by m. For elementary facts concerning convexity and half-spaces see, for example, [34].

DEFINITION 2.1 (essential support). *The (essential) support* supp f *of a measurable function* $f : \Omega \to \mathbb{R}$ *is the complement of the union of all open sets* $O \subseteq \Omega$ *where* $f|_O = 0$ *almost everywhere.*

DEFINITION 2.2 (measure theoretic interior). *The measure theoretic interior* A° *of a measurable set* A *is the set of all points* $x \in A$ *with density 1, i.e.,*

$$\lim_{\epsilon \to 0} \frac{m(A \cap B_\epsilon(x))}{m(B_\epsilon(x))} = 1.$$

DEFINITION 2.3 (closed half-space). *For* $\rho \in \mathbb{R}^n$ *and* $t \in \mathbb{R}$, *we define the closed half-space*

$$H_{\rho,t} := \{x \in \mathbb{R}^n : x \cdot \rho \le t\}.$$

DEFINITION 2.4 (closed essential convex hull). *For a measurable set* $A \subset \mathbb{R}^n$, *the (closed essential) convex hull* conv A *is the intersection of all closed half-spaces* $H_{\rho,t}$ ($\rho \in \mathbb{R}^n$, $t \in \mathbb{R}$) *such that*

(2.1) $$m(A \setminus H_{\rho,t}) = 0.$$

DEFINITION 2.5 (essential convex support function). *If* $A \subseteq \mathbb{R}^n$ *is a bounded measurable set with positive measure, then we define its essential convex support function* $h_A : S^{n-1} \to \mathbb{R}$ *of* A *by*

(2.2) $$h_A(\rho) = \inf\{t \in \mathbb{R} : m(A \setminus H_{\rho,t}) = 0\}.$$

Since A has positive measure, the infimum is taken over a nonempty set. The boundedness of A ensures that the set $\{t \in \mathbb{R} : m(A \setminus H_{\rho,t}) = 0\}$ is bounded from below. We have

$$t \ge h_A(\rho) \implies m(A \setminus \{x \cdot \rho \le t\}) = 0,$$
$$t < h_A(\rho) \implies m(A \setminus \{x \cdot \rho \le t\}) > 0.$$

LEMMA 2.6.
(a) *For a measurable set* A, conv A *is a closed and convex set.*
(b) *If* A *is measurable, then the Lebesgue density of* A *is zero for all points* $x \in A \setminus \operatorname{conv}(A)$. *In particular,*

$$A^\circ \subseteq \operatorname{conv}(A) \quad and \quad A \subseteq \operatorname{conv}(A) \cup N \quad with\ a\ null\ set\ N.$$

(c) *For measurable sets* $A, B \subseteq \mathbb{R}^n$ *and a null set* $N \subseteq \mathbb{R}^n$,

$$B \subseteq \operatorname{conv}(A) \cup N \quad implies \quad \operatorname{conv}(B) \subseteq \operatorname{conv}(A).$$

(d) *For a measurable function $f : \Omega \to \mathbb{R}$, conv(supp f) is the intersection of all closed half-spaces $H_{\rho,t}$ with $f|_{\Omega \setminus H_{\rho,t}} = 0$ almost everywhere.*

Proof.

(a) The convex hull conv A is an intersection of closed and convex half-spaces and thus is closed and convex.

(b) Let $x \in A \setminus \text{conv}(A)$. By Definition 2.4 there exists a closed half-space $H_{\rho,t}$ with $m(A \setminus H_{\rho,t}) = 0$ and $x \notin H_{\rho,t}$. Since $H_{\rho,t}$ is closed, $B_\epsilon(x) \cap H_{\rho,t} = \emptyset$ for all sufficiently small balls $B_\epsilon(x)$, so that

$$m(A \cap B_\epsilon(x)) \leq m((H_{\rho,t} \cap B_\epsilon(x)) \cup (A \setminus H_{\rho,t}))$$
$$\leq m(H_{\rho,t} \cap B_\epsilon(x)) + m(A \setminus H_{\rho,t}) = 0.$$

This shows that the Lebesgue density of A is zero in x.

In particular, $A^\circ \subseteq \text{conv}(A)$ and, by the Lebesgue density theorem, $m(A \setminus \text{conv } A) = 0$.

(c) Let $B \subseteq \text{conv}(A) \cup N$. Then B is a subset of $H_{\rho,t} \cup N$ for every half-space $H_{\rho,t}$ with $m(A \setminus H_{\rho,t}) = 0$. Hence, $m(B \setminus H_{\rho,t}) = 0$ for all half-spaces $H_{\rho,t}$ with $m(A \setminus H_{\rho,t}) = 0$; i.e., the set of half-spaces with $m(B \setminus H_{\rho,t}) = 0$ is a superset of those with $m(A \setminus H_{\rho,t}) = 0$. Thus the intersection of all half-spaces with $m(B \setminus H_{\rho,t}) = 0$ is a subset of the intersection of all half-spaces with $m(A \setminus H_{\rho,t}) = 0$, which shows $\text{conv}(B) \subseteq \text{conv}(A)$.

(d) By definition, $\text{conv}(\text{supp}(f))$ is the intersection of all closed half-spaces H for which $\text{supp}(f) \setminus H$ is a null set. Hence, it suffices to show that a closed half-space H fulfills $m(\text{supp}(f) \setminus H) = 0$ if and only if $f|_{\Omega \setminus H} = 0$ almost everywhere.

Let O be the union of all open sets on which f is zero almost everywhere. Then

$$\text{supp}(f) \setminus H = (\Omega \setminus O) \setminus H = \Omega \setminus (H \cup O) = (\Omega \setminus H) \setminus O.$$

If $m(\text{supp}(f) \setminus H) = 0$, then $\Omega \setminus H$ is a subset of $O \cup N$ with some null set N, so $f|_{\Omega \setminus H} = 0$ almost everywhere. On the other hand, if $f|_{\Omega \setminus H} = 0$ almost everywhere, then $\Omega \setminus H \subseteq O \cup N$ where N is a null set, and thus $(\Omega \setminus H) \setminus O = \text{supp}(f) \setminus H$ is a null set. $\qquad\square$

2.2. Wolff solutions and monotonicity for the nonlinear DN map. Let Λ_σ be the (nonlinear) Dirichlet–Neumann (DN) map for the weighted p-Laplace equation $(1 < p < \infty)$, or p-conductivity equation, i.e.,

$$\Lambda_\sigma : \ X \to X'$$

defined by

$$(\Lambda_\sigma(f), g) := \int_\Omega \sigma |\nabla u_\sigma^f|^{p-2} \nabla u_\sigma^f \cdot \nabla v^g \mathrm{d}x, \qquad f, g \in X,$$

where $X := W^{1,p}(\Omega)/W_0^{1,p}(\Omega)$, $v^g \in W^{1,p}(\Omega)$ is any representative of the quotient space element $g \in X$, and u_σ^f is the weak solution of the weighted p-Laplacian with boundary value f; i.e., u_σ^f is the unique minimizer in $f + W_0^{1,p}(\Omega)$ of the p-Dirichlet energy functional

$$u \mapsto E_\sigma(u) := \int_\Omega \sigma |\nabla u|^p \mathrm{d}x.$$

The map Λ_σ is well-defined by this definition [5, section 3.2].

We will use a certain exponential solution for the p-Laplace equation, which is real valued and periodic in one direction and exponentially behaving in the other direction. This specific solution is the Wolff solution, which was introduced by Wolff [43] and has applications in inverse problems, including boundary determination [3, 35] and inclusion detection [5, 6]. We explicitly write the solutions as follows (see [6, Lemma 3.1]).

LEMMA 2.7 (Wolff solutions). *Let $\rho, \rho^\perp \in \mathbb{R}^n$ satisfy $|\rho| = |\rho^\perp| = 1$ and $\rho \cdot \rho^\perp = 0$. Define $h \colon \mathbb{R}^n \to \mathbb{R}$ by $h(x) = e^{-\rho \cdot x} w(\rho^\perp \cdot x)$, where the function w satisfies the differential equation*

$$w''(s) + V(w, w')w = 0 \tag{2.3}$$

with

$$V(w, w') = \frac{(2p - 3)\,(w')^2 + (p - 1)w^2}{(p - 1)\,(w')^2 + w^2}. \tag{2.4}$$

The function h is then p-harmonic.

Given any initial conditions $(a_0, b_0) \in \mathbb{R}^2 \setminus \{(0, 0)\}$, there exists a solution $w \in C^\infty(\mathbb{R})$ to the differential equation (2.3) which is periodic with period $\lambda_p > 0$ and satisfies the initial conditions $(w(0), w'(0)) = (a_0, b_0)$ and $\int_0^{\lambda_p} w(s)\,ds = 0$. Furthermore, there exist constants c and C depending on a_0, b_0, and p such that for all $s \in \mathbb{R}$ we have

$$C > (w(s))^2 + (w'(s))^2 > c > 0. \tag{2.5}$$

For a large parameter $\tau \in \mathbb{R}$ and a fixed constant $t \in \mathbb{R}$, we now define the Wolff-type solutions $u_\tau \colon \mathbb{R}^n \to \mathbb{R}$ by

$$u_\tau(x) = e^{\tau(x \cdot \rho - t)} w\left(\tau x \cdot \rho^\perp\right). \tag{2.6}$$

The upper and lower bounds for the Wolff solution are due to [6, Lemma 3.1 and equations (3.5) and (3.6)].

LEMMA 2.8. *Let $\Omega \subset \mathbb{R}^n$ be a bounded open set. There exist $c, C > 0$ so that, for each $\rho \in S^{n-1}$, $\tau > 0$, and $t \in \mathbb{R}$, the function (2.6) solves $\mathrm{div}(|\nabla u|^{p-2}\nabla u) = 0$ in $\mathbb{R}^n$, satisfies $u_\tau|_\Omega \in C^\infty(\overline{\Omega}) \subseteq W^{1,p}(\Omega)$, and where*

$$c\tau e^{\tau(x \cdot \rho - t)} \le |\nabla u_\tau(x)| \le C\tau e^{\tau(x \cdot \rho - t)}, \quad x \in \Omega.$$

We will also use the monotonicity inequality [6, Lemma 2.1].

LEMMA 2.9 (monotonicity inequality). *Let $\sigma_0, \sigma_1 \in L^\infty_+(\Omega)$, $1 < p < \infty$. For every $f \in W^{1,p}(\Omega)$ and corresponding minimizer $u_0 \in f + W^{1,p}_0(\Omega)$ of the energy E_{σ_0}, it holds that*

$$
\begin{aligned}
(p - 1) &\int_\Omega \frac{\sigma_0}{\sigma_1^{1/(p-1)}} \left(\sigma_1^{\frac{1}{p-1}} - \sigma_0^{\frac{1}{p-1}}\right) |\nabla u_0|^p \mathrm{d}x \\
&\le ((\Lambda_{\sigma_1} - \Lambda_{\sigma_0})f, f) \le \int_\Omega (\sigma_1 - \sigma_0)|\nabla u_0|^p \mathrm{d}x.
\end{aligned}
\tag{2.7}
$$

3. Monotonicity method for the p-Laplacian. As above, let $\sigma \in L^\infty_+(\Omega)$. We make the global monotonicity assumption that either $\sigma(x) \geq 1$ holds (almost everywhere) in Ω or $\sigma(x) \leq 1$ holds (almost everywhere) in Ω.

We will show that the essential convex hull of the support of $|\sigma - 1|$ is uniquely determined by the DN operator Λ_σ and that it can be recovered by monotonicity tests.

The monotonicity tests are based on comparing Λ_σ to certain nonlinear test operators in the following sense. For two (possibly nonlinear) operators $\Lambda_1, \Lambda_2 : X \to X'$ we write that

$$\Lambda_1 \geq \Lambda_2 \quad \text{if} \quad ((\Lambda_1(f) - \Lambda_2(f)), f) \geq 0 \quad \text{for all } f \in X.$$

Note that this defines a reflexive and transitive relation (a preorder), but not necessarily a partial order since antisymmetry may fail.

Let Λ_0 denote the DN operator for the homogeneous conductivity $\sigma_0 = 1$. For a measurable set $B \subset \Omega$, we also introduce the operator $\Lambda_B : X \to X'$ by

$$(\Lambda_B(f), g) := \int_B |\nabla u_0^f|^{p-2} \nabla u_0^f \cdot \nabla u_0^g \, \mathrm{d}x, \qquad f, g \in X,$$

where u_0^h is the unique minimizer in $h + W_0^{1,p}(\Omega)$ of the p-Dirichlet energy functional E_{σ_0} with $\sigma_0 = 1$.

THEOREM 3.1. *Let $B \subseteq \Omega$ be a measurable set with positive measure. For every constant $\alpha > 0$,*
(a) *$\sigma|_B \geq 1 + \alpha$ implies $\Lambda_0 + \tilde{\alpha}\Lambda_B \leq \Lambda_\sigma$,*
(b) *$\sigma|_B \leq 1 - \alpha$ implies $\Lambda_0 - \alpha\Lambda_B \geq \Lambda_\sigma$,*
(c) *$\Lambda_0 + \alpha\Lambda_B \leq \Lambda_\sigma$ implies $B \subseteq \mathrm{conv}(\mathrm{supp}(|\sigma - 1|))$ up to null sets,*
(d) *$\Lambda_0 - \alpha\Lambda_B \geq \Lambda_\sigma$ implies $B \subseteq \mathrm{conv}(\mathrm{supp}(|\sigma - 1|))$ up to null sets,*
where in (a) *$\tilde{\alpha} := (p-1)(1 - (1+\alpha)^{-\frac{1}{p-1}}) > 0$, and in* (c) *and* (d), *the notion "$B \subseteq A$ up to null sets" means that $m(B \setminus A) = 0$.*

Before we give the proof of Theorem 3.1, we show that this implies that Λ_σ uniquely determines the convex hull of the support of $|\sigma - 1|$, and we comment on a possible numerical implementation.

COROLLARY 3.2. *Let either $\sigma(x) \geq 1$ hold (almost everywhere) in Ω or $\sigma(x) \leq 1$ hold (almost everywhere) in Ω. Then*

$$\mathrm{conv}(\mathrm{supp}(|\sigma - 1|)) = \mathrm{conv}\left(\bigcup_{B \in \mathcal{B}} B^\circ\right),$$

where $\mathcal{B}$ is the family of all measurable sets $B \subseteq \Omega$ for which there exists an $\alpha > 0$ such that either $\Lambda_0 + \alpha\Lambda_B \leq \Lambda_\sigma$ or $\Lambda_0 - \alpha\Lambda_B \geq \Lambda_\sigma$.

Proof. By Theorem 3.1(c),(d), for every $B \in \mathcal{B}$ there exists a null set N with

(3.1) $$B \subseteq \mathrm{conv}(\mathrm{supp}(|\sigma - 1|)) \cup N.$$

Lemma 2.6 then implies $B^\circ \subseteq \mathrm{conv}(B) \subseteq \mathrm{conv}(\mathrm{supp}(|\sigma - 1|))$, so that

(3.2) $$\bigcup_{B \in \mathcal{B}} B^\circ \subseteq \mathrm{conv}(\mathrm{supp}(|\sigma - 1|)),$$

and again by Lemma 2.6,

$$(3.3) \qquad \text{conv}(\text{supp}(|\sigma - 1|)) \supseteq \text{conv}\left(\bigcup_{B \in \mathcal{B}} B^\circ\right).$$

To show "$\subseteq$" we assume that

$$\text{conv}(\text{supp}(|\sigma - 1|)) \setminus \text{conv}\left(\bigcup_{B \in \mathcal{B}} B^\circ\right) \neq \emptyset.$$

By Definition 2.4, there would then exist a closed half-space H with

$$(3.4) \qquad m\left(\left(\bigcup_{B \in \mathcal{B}} B^\circ\right) \setminus H\right) = 0 \quad \text{and} \quad \text{conv}(\text{supp}(|\sigma - 1|)) \setminus H \neq \emptyset.$$

With Lemma 2.6 this would imply that every closed half-space $\tilde{H}$ with $|\sigma - 1| \, |_{\Omega \setminus \tilde{H}} = 0$ almost everywhere must have an intersection of positive measure with the complement of H. Hence $|\sigma - 1|$ could not be zero almost everywhere on $\Omega \setminus H$. But then there would exist a measurable set $B \subseteq \Omega \setminus H$ with $B = B^\circ$, $m(B) = m(B \setminus H) > 0$ and a constant $\alpha > 0$ with $|\sigma - 1| \geq \alpha$ in B. Theorem 3.1(a),(b) would then imply that $B \in \mathcal{B}$, which would contradict the first part of (3.4). □

Remark 3.3. Theorem 3.1 can be implemented in a numerical reconstruction algorithm as in [17, 19, 20]. We choose a finite pixel partition $\Omega = \bigcup_{i=1}^{m} P_i$ with (not necessarily disjoint) measurable subsets $P_i \subseteq \Omega$. For each pixel P_i one can numerically approximate the maximal number $\alpha_i \geq 0$ for which $\Lambda_0 + \tilde{\alpha}\Lambda_B \leq \Lambda_\sigma$ (in the case $\sigma \geq 1$), resp., $\Lambda_0 - \alpha\Lambda_B \geq \Lambda_\sigma$ (in the case $\sigma \leq 1$). Theorem 3.1 then yields that α_i will be nonzero only on pixels inside the convex hull of $\text{supp}(|\sigma - 1|)$. Moreover, if $|\sigma - 1| \geq c > 0$ on P_i, then $\alpha_i \geq (p-1)(1 - (1+c)^{-\frac{1}{p-1}})$ (in the case $\sigma \geq 1$), resp., $\alpha_i \geq c$ (in the case $\sigma \leq 1$). In that sense, a plot of α_i will show the convex hull of an inclusion up to the pixel partition and also give an estimate of the contrast of the inclusion. Note, however, that the calculation of α_i will be considerably more involved than in the linear case, where it could be calculated from a simple eigenvalue calculation; cf. [17, section 4].

Proof of Theorem 3.1.

(a) Note that by our global monotonicity assumption, $\sigma|_B \geq 1 + \alpha$ implies that we are in the case that $\sigma \geq 1$ almost everywhere in Ω. Hence we obtain from the monotonicity lemma, Lemma 2.9, that

$$\begin{aligned}
((\Lambda_\sigma - \Lambda_0)(f), f) &\geq (p-1)\int_\Omega \frac{1}{\sigma^{1/(p-1)}}\left(\sigma^{\frac{1}{p-1}} - 1\right)|\nabla u_0|^p \mathrm{d}x \\
&\geq (p-1)\int_B \left(1 - \sigma^{-\frac{1}{p-1}}\right)|\nabla u_0|^p \mathrm{d}x \\
&\geq (p-1)\left(1 - (1+\alpha)^{-\frac{1}{p-1}}\right)(\Lambda_B(f), f).
\end{aligned}$$

(b) If $\sigma|_B \leq 1 - \alpha$, then we are in the case that $\sigma \leq 1$ almost everywhere in Ω, and we obtain from the monotonicity lemma, Lemma 2.9, that

$$\begin{aligned}
((\Lambda_\sigma - \Lambda_0)(f), f) &\leq \int_\Omega (\sigma - 1)|\nabla u_0|^p \mathrm{d}x \leq \int_B (\sigma - 1)|\nabla u_0|^p \mathrm{d}x \\
&\leq -\alpha(\Lambda_B(f), f).
\end{aligned}$$

(c) We set $D := \mathrm{conv}(\mathrm{supp}(|\sigma - 1|))$ and assume that $B \setminus D$ is not a null set. We will prove that

$$(3.5) \qquad\qquad \Lambda_0 + \alpha\Lambda_B \nleq \Lambda_\sigma.$$

Since shrinking B will lead to a smaller Λ_B, it suffices to prove (3.5) with B replaced by a subset of B. Hence, by replacing B with $B \setminus D$, we can assume without loss of generality that

$$m(B) > 0 \quad \text{and} \quad B \cap D = \emptyset.$$

Moreover, it follows from elementary measure theory or the Lebesgue density theorem that there exists a point $x \in B$ such that $m(B \cap B_r(x)) > 0$ for all sufficiently small open balls $B_r(x)$. By replacing B with $B \cap B_r(x)$ for a sufficiently small ball, we can therefore assume without loss of generality that

$$m(B) > 0 \quad \text{and} \quad \mathrm{conv}(B) \cap D = \emptyset.$$

Using the Hahn–Banach (or hyperplane) separation theorem (see, for example, [34, Corollary 11.4.2]) we then obtain a vector $\rho \in \mathbb{R}^n$, $|\rho| = 1$, and numbers $t \in \mathbb{R}$, $\epsilon > 0$ such that

$$
\begin{aligned}
x \cdot \rho &< t & &\text{for all } x \in D, \\
x \cdot \rho &> t + \epsilon & &\text{for all } x \in B.
\end{aligned}
$$

With the Wolff solutions from Lemma 2.8, it follows that there exist constants $c, C > 0$ so that for each $\tau > 0$ there exists a solution $u_{0,\tau}$ of the homogeneous p-Laplace equation with

$$c\tau e^{\tau(x \cdot \rho - t)} \le |\nabla u_{0,\tau}(x)| \le C\tau e^{\tau(x \cdot \rho - t)} \quad \text{for all } x \in \Omega.$$

With $f_\tau := u_{0,\tau}|_{\partial\Omega}$ it follows that

$$(\alpha\Lambda_B(f_\tau), f_\tau) = \alpha\int_B |\nabla u_{0,\tau}|^p \mathrm{d}x \ge \alpha m(B)c^p\tau^p e^{p\tau\epsilon},$$

and using the monotonicity lemma, Lemma 2.9, we obtain

$$(3.6) \qquad
\begin{aligned}
((\Lambda_\sigma - \Lambda_0)(f_\tau), f_\tau) &\le \int_\Omega (\sigma - 1)|\nabla u_{0,\tau}|^p \mathrm{d}x \\
&\le \|\sigma - 1\|_{L^\infty(\Omega)} C^p\tau^p m(D).
\end{aligned}
$$

For large enough τ we have that

$$(3.7) \qquad\qquad \|\sigma - 1\|_{L^\infty(\Omega)} C^p m(D) < \alpha m(B)c^p e^{p\tau\epsilon}$$

and thus

$$(\alpha\Lambda_B(f_\tau), f_\tau) > ((\Lambda_\sigma - \Lambda_0)(f_\tau), f_\tau),$$

which proves (3.5).

(d) As in (c) we set $D := \operatorname{conv}(\operatorname{supp}(|\sigma-1|))$ and assume without loss of generality that $m(B) > 0$ and $\operatorname{conv}(B) \cap D = \emptyset$. With the same Wolff solutions as in (c) we obtain from the monotonicity lemma, Lemma 2.9, that

$$((\Lambda_\sigma - \Lambda_0)(f_\tau), f_\tau) \geq (p-1) \int_\Omega \left(1 - \sigma^{-\frac{1}{p-1}}\right) |\nabla u_{0,\tau}|^p \mathrm{d}x$$

$$\geq -(p-1) \left\|\sigma^{-1}\right\|_{L^\infty(\Omega)}^{\frac{1}{p-1}} \int_D |\nabla u_{0,\tau}|^p \mathrm{d}x$$

$$\geq -(p-1) \left\|\sigma^{-1}\right\|_{L^\infty(\Omega)}^{\frac{1}{p-1}} m(D) C^p \tau^p,$$

so that for large enough $\tau > 0$,

$$(-\alpha\Lambda_B(f_\tau), f_\tau) < ((\Lambda_\sigma - \Lambda_0)(f_\tau), f_\tau),$$

which shows that $\Lambda_\sigma - \Lambda_0 \nleq -\alpha\Lambda_B$. $\square$

4. Enclosure method. We define the indicator function $I_\rho(t, \tau)$ as

$$I_\rho(t, \tau) = \tau^{-p} \langle (\Lambda_\sigma - \Lambda_1) (f_\tau), f_\tau \rangle,$$

where $\rho \in \mathbb{R}^n$ is a unit vector, $t \in \mathbb{R}$, $\tau > 0$, and $f_\tau = u_\tau|_{\partial\Omega}$ where u_τ are the Wolff solutions given by (2.6).

The indicator function for any ρ, t, and τ is determined by Λ_σ. Thus the next theorem implies that Λ_σ determines the essential convex hull of the inclusion D, where we write

$$D = \operatorname{supp}(\sigma - 1).$$

THEOREM 4.1. *Suppose that $\sigma \in L_+^\infty(\Omega)$ and that either $\sigma \geq 1$ or $\sigma \leq 1$ almost everywhere. Then*

$$\operatorname{conv} D = \bigcap_{\rho \in \mathbb{R}^n; |\rho|=1} \overline{H}_\rho,$$

where

$$H_\rho = \{x \in \mathbb{R}^n ; |I_\rho(x \cdot \rho, \tau)| \to \infty \text{ as } \tau \to \infty\}.$$

Before the proof, we give a simple lemma. First recall that h_D is the convex support function of the set D, as defined in section 2.

LEMMA 4.2. *Suppose that $\sigma \in L_+^\infty(\Omega)$ and that either $\sigma \geq 1$ or $\sigma \leq 1$ almost everywhere. Let also $t < h_D(\rho)$. Then there exists a set $S \subset D$ which satisfies the following conditions:*

1. *$m(S) > 0$.*
2. *There is $\varepsilon_1 > 0$ such that for all $x \in S$ we have $x \cdot \rho > t + \varepsilon_1$.*
3. *There is $\varepsilon_2 > 0$ such that for all $x \in S$ we have $\sigma(x) > 1 + \varepsilon_2$ or $\sigma(x) + \varepsilon_2 < 1$.*

Proof. Choose ε_1 so that $0 < \varepsilon_1 < h_D(\rho) - t$, and define

$$\tilde{S} = \{x \in D ; x \cdot \rho > t + \varepsilon_1\}.$$

Then items 1 and 2 are true for $\tilde{S}$ by the definition of the essential convex hull and h_D. To also satisfy item 3, we observe that $\tilde{S}$ is a set of positive measure and that $\sigma > 1$ (or $\sigma < 1$) in $\tilde{S}$. Since

$$\tilde{S} = \bigcup_{k \in \mathbb{Z}_+} \left\{x \in \tilde{S}; |\sigma - 1| > 1/k\right\},$$

and $\tilde{S}$ has positive measure, at least one of the sets $\left\{x \in \tilde{S}; |\sigma - 1| > 1/k\right\}$ must have positive measure. We choose it as the set S. $\qquad\square$

Proof of Theorem 4.1. First, if $\sigma = 1$ almost everywhere, the indicator function vanishes, and the claim is true. Let us suppose this is not the case. Then $D = \operatorname{supp}(\sigma - 1)$ is nonempty, and we consider a fixed direction $\rho \in \mathbb{R}^n$, $|\rho| = 1$.

Suppose first that $\sigma \geq 1$ almost everywhere. If $t > h_D(\rho)$, then we have by the monotonicity inequality (Lemma 2.9) and Lemma 2.8 that

$$0 \leq I_\rho(t, \tau) = \tau^{-p} \left\langle (\Lambda_\sigma - \Lambda_1)(f_\tau), f_\tau \right\rangle \leq \tau^{-p} \int_\Omega (\sigma - 1) |\nabla u_\tau|^p \, \mathrm{d}x$$

$$= \tau^{-p} \int_D (\sigma - 1) |\nabla u_\tau|^p \, \mathrm{d}x$$

$$\leq C \int_D (\sigma - 1) e^{p\tau(x \cdot \rho - t)} \mathrm{d}x$$

$$\leq C \|\sigma - 1\|_{L^\infty(\Omega)} \int_{D \cap \{x \cdot \rho \leq h_D\}} e^{p\tau(x \cdot \rho - t)} \mathrm{d}x$$

$$\leq C \|\sigma - 1\|_{L^\infty(\Omega)} \int_{D \cap \{x \cdot \rho \leq h_D\}} e^{p\tau(h_D - t)} \mathrm{d}x$$

$$\leq C \|\sigma - 1\|_{L^\infty(\Omega)} \, m(D) e^{p\tau(h_D - t)} \to 0$$

as $\tau \to \infty$. Note that we used the fact that $m(D \setminus \{x \cdot \rho \leq h_D\}) = 0$.

Next, suppose $t < h_D(\rho)$. By the monotonicity inequality, Lemma 2.9, we get

$$I_\rho(t, \tau) = \tau^{-p} \left\langle (\Lambda_\sigma - \Lambda_1)(f_\tau), f_\tau \right\rangle$$

$$\geq (p - 1) \tau^{-p} \int_D \left(1 - \sigma^{-1/(p-1)}\right) |\nabla u_\tau|^p \mathrm{d}x$$

$$\geq (p - 1) \int_D \left(1 - \sigma^{-1/(p-1)}\right) e^{p\tau(x \cdot \rho - t)} \mathrm{d}x.$$

By Lemma 4.2, there is a set $S \subseteq D$ with positive measure so that $x \cdot \rho > t + \varepsilon_1$ and that $1 - \sigma^{-1/(p-1)} > \varepsilon_3$ on S for some $\varepsilon_1, \varepsilon_3 > 0$. Thus

$$\int_D \left(1 - \sigma^{-1/(p-1)}\right) e^{p\tau(x \cdot \rho - t)} \mathrm{d}x \geq \int_S \varepsilon_3 e^{p\tau\varepsilon_1} \mathrm{d}x.$$

This shows that $I_\rho(t, \tau) \to \infty$ as $\tau \to \infty$.

If instead $\sigma \leq 1$ almost everywhere, then the previous proof works with minor changes and gives that $I_\rho(t, \tau) \to 0$ for $t > h_D(\rho)$ and $I_\rho(t, \tau) \to -\infty$ for $t < h_D(\rho)$. Also note that, by definition, the set H_ρ has the following form:

$$H_\rho = \left\{x \in \mathbb{R}^n ; |I_\rho(x \cdot \rho, \tau)| \to \infty \text{ as } \tau \to \infty\right\}.$$

Therefore, we have established that

$$\{x \in \mathbb{R}^n; x \cdot \rho < h_D(\rho)\} \subseteq H_\rho \subseteq \{x \in \mathbb{R}^n; x \cdot \rho \leq h_D(\rho)\},$$

and thus

$$\overline{H}_\rho = \{x \in \mathbb{R}^n; x \cdot \rho \leq h_D(\rho)\}.$$

By the definition of h_D, one has $m\left(D \setminus \overline{H}_\rho\right) = 0$ for each direction ρ. This implies, by definition of the essential convex hull, that

$$\operatorname{conv} D \subseteq \bigcap_{\rho \in \mathbb{R}^n; |\rho|=1} \overline{H}_\rho.$$

To see the other direction, consider some $x_0 \in \bigcap_{\rho \in \mathbb{R}^n; |\rho|=1} \overline{H}_\rho$, and an arbitrary closed hyperplane H satisfying $m\left(D \setminus H\right) = 0$. There exists $\rho_0 \in \mathbb{R}^n, |\rho_0| = 1$, and $t_0 \in \mathbb{R}$ such that

$$H = \{x \in \mathbb{R}^n; x \cdot \rho_0 \le t_0\}.$$

By definition of h_D we then have $t_0 \ge h_D(\rho_0)$, so that $\overline{H}_{\rho_0} \subseteq H$ and consequently $x_0 \in H$. □

Remark 4.3. The situation where the conductivity takes values both less than one and greater than one in sets of positive measure is similar to that in [6, section 5].

Remark 4.4. The behavior of the indicator function when $t = h_D$ is tricky. By the same proof as in [6, Lemma 4.6] we have $I_\rho(h_D, \tau) \le C$ (note that we have a different power of τ in the indicator function). The lower bound uses Lipschitz regularity and a jump condition on the boundary of the inclusion, so it seems unlikely to work in the present case.

5. Boundary determination. In this section we give a boundary determination result for the p-Calderón problem. This result was proved in [35, Theorem 1.1] for domains with C^1 boundary. We assume that the domain is strictly convex instead.

THEOREM 5.1. *Let* $\Omega \subset \mathbb{R}^n$, $n \ge 2$, *be a bounded open set so that* $\overline{\Omega}$ *is strictly convex. Suppose* $\sigma \in C(\overline{\Omega})$ *is a positive function. Then the nonlinear DN map* Λ_σ *determines* $\sigma|_{\partial\Omega}$.

Proof. Let $x_0 \in \partial\Omega$ be an arbitrary boundary point. Since $\overline{\Omega}$ is convex, the supporting hyperplane theorem of Minkowski (see, for example, [34, Theorem 11.6]) implies that there is a closed half-space H with

$$\overline{\Omega} \subset H, \quad x_0 \in \partial H.$$

By strict convexity one has $\overline{\Omega} \cap \partial H = \{x_0\}$ (for if $\overline{\Omega} \cap \partial H$ were to contain another point y_0, then the line segment between x_0 and y_0 would lie in $\overline{\Omega} \cap \partial H$ and thus also in $\partial\Omega$, contradicting strict convexity). The half-space H may be written as

$$H = \{x \in \mathbb{R}^n \,; \rho \cdot x \le t_0\}$$

for some unit vector $\rho \in \mathbb{R}^n$, which will be fixed from now on, and some $t_0 \in \mathbb{R}$.

Let $\gamma \in \mathbb{R}_+$ be a constant, which we will use as a test conductivity. Define the indicator function

$$I_\gamma(t, \tau) = \tau^{-p} \langle (\Lambda_\sigma - \Lambda_\gamma)(f_\tau), f_\tau \rangle,$$

where $t \in \mathbb{R}$, $\tau > 0$, and $f_\tau = u_\tau|_{\partial\Omega}$ is the boundary value of the Wolff solution u_τ (which depends on ρ) that solves $\operatorname{div}(\gamma |\nabla u|^{p-2} \nabla u) = 0$ in Ω.

For any fixed $t < t_0$, we are going to show that

(5.1) $$\gamma < \sigma(x_0) \implies I_\gamma(t, \tau) \to +\infty \text{ as } \tau \to \infty,$$

(5.2) $$\gamma > \sigma(x_0) \implies I_\gamma(t, \tau) \to -\infty \text{ as } \tau \to \infty.$$

Then $\sigma(x_0) = \inf \{\gamma > 0 \,;\, \lim_{\tau \to \infty} I_\gamma(t,\tau) = -\infty\}$, which shows that $\sigma(x_0)$ can be determined from Λ_σ, thus concluding the proof.

Suppose that $\gamma \neq \sigma(x_0)$. By continuity of σ, there are $\varepsilon, r > 0$ such that

$$|\sigma(x) - \gamma| \geq \varepsilon, \qquad x \in B(x_0, r) \cap \overline{\Omega}.$$

For $t \in \mathbb{R}$ consider the set

$$S_t = \{x \in \mathbb{R}^n \,;\, \rho \cdot x > t\} \cap \overline{\Omega}.$$

We claim that there exists $\delta > 0$ so that

$$S_{t_0 - \delta} \subset B(x_0, r) \cap \overline{\Omega}.$$

If not, then for any j there is $x_j \in \overline{\Omega}$ with $\rho \cdot x_j > t_0 - 1/j$ but $x_j \notin B(x_0, r)$. The set $\overline{\Omega}$ is compact, and hence some subsequence of (x_j) converges to some $x_\infty \in \overline{\Omega}$ with $\rho \cdot x_\infty \geq t_0$. But we saw earlier that $\overline{\Omega} \subset H$ and $\overline{\Omega} \cap \partial H = \{x_0\}$ where $H = \{\rho \cdot x \leq t_0\}$, showing that $x_\infty = x_0$, which is a contradiction.

Now choose t with $t_0 - \delta < t < t_0$. Then the set S_t is nonempty with positive measure, and one has

$$|\sigma(x) - \gamma| \geq \varepsilon, \quad x \in S_t.$$

Let $s \in \{+1, -1\}$ be the sign of $\sigma - \gamma$ in S_t. In the monotonicity inequality, Lemma 2.9, we write

$$F_+(a,b) = (p-1) \frac{a}{b^{1/(p-1)}} \left(b^{1/(p-1)} - a^{1/(p-1)} \right),$$

$$F_-(a,b) = a - b$$

and estimate

$$sI_\gamma(t,\tau) = s\tau^{-p} \langle (\Lambda_\sigma - \Lambda_\gamma) f_\tau, f_\tau \rangle \geq s\tau^{-p} \int_\Omega F_s(\sigma(x), \gamma) |\nabla u_\tau|^p \, dx$$

$$= s\tau^{-p} \left(\int_{S_t} F_s(\sigma(x), \gamma) |\nabla u_\tau|^p \, dx + \int_{\Omega \setminus S_t} F_s(\sigma(x), \gamma) |\nabla u_\tau|^p \, dx \right)$$

$$\geq s \left(c \int_{S_t} F_s(\sigma(x), \gamma) e^{p\tau(x \cdot \rho - t)} dx - C \int_{\Omega \setminus S_t} F_s(\sigma(x), \gamma) e^{p\tau(x \cdot \rho - t)} dx \right)$$

$$\geq c \int_{S_t} e^{p\tau(x \cdot \rho - t)} dx - C \int_{\Omega \setminus S_t} e^{p\tau(x \cdot \rho - t)} dx.$$

If $t < t' < t_0$, then $S_{t'}$ has positive measure and one has

$$\int_{S_t} e^{p\tau(x \cdot \rho - t)} dx \geq \int_{S_{t'}} e^{p\tau(x \cdot \rho - t)} dx \geq e^{p\tau(t' - t)} m(S_{t'}) \to \infty$$

as $\tau \to \infty$, and have

$$\int_{\Omega \setminus S_t} e^{p\tau(x \cdot \rho - t)} dx \leq \int_{\Omega \setminus S_t} dx \leq C.$$

Thus $I_\gamma \to +\infty$ when $\sigma(x_0) > \gamma$ and $I_\gamma \to -\infty$ when the opposite inequality holds. This proves (5.1)–(5.2) under the condition $t_0 - \delta < t < t_0$, but the result holds also for $t \leq t_0 - \delta$ since

$$I_\gamma(t, \tau) = I_\gamma(t_0 - \delta/2, \tau) e^{p\tau(t_0 - \delta/2 - t)},$$

which follows by a short computation. $\qquad \square$

Remark 5.2. It is clear that the same proof also gives the following result: if $\Omega \subset \mathbb{R}^n$ is a bounded open set, if Ω is strictly convex at $x_0 \in \partial\Omega$ in the sense that there is a closed half-space H with

$$\overline{\Omega} \subset H, \quad \overline{\Omega} \cap \partial H = \{x_0\},$$

and if $\sigma \in L_+^\infty(\Omega)$ is continuous near x_0, then the knowledge of Λ_σ determines $\sigma(x_0)$.

Remark 5.3. The proof gives the following simple algorithm for boundary determination in strictly convex domains:

1. First, take an arbitrary point $x_0 \in \partial\Omega$. Our aim is to find the value of σ at x_0.
2. Select a direction ρ such that $\rho \cdot x_0 = \max_{x \in \overline{\Omega}} \rho \cdot x = t_0$. This is possible by strict convexity of the set Ω.
3. Fix a parameter t so that $t < t_0$, which corresponds to taking a hyperplane that intersects Ω.
4. Choose a small $\gamma_0 \in \mathbb{R}_+$ and a large $\gamma^0 \in \mathbb{R}_+$ such that

$$I_{\gamma_0} \to +\infty \text{ and } I_{\gamma^0} \to -\infty \text{ as } \tau \to \infty.$$

5. Suppose γ_{j-1} and γ^{j-1} have been determined. Define $\kappa_j = \frac{\gamma_{j-1}+\gamma^{j-1}}{2}$ and calculate I_{κ_j}.
6. If $I_{\kappa_j} \to +\infty$, take $\gamma_j = \kappa_j$. If $I_{\kappa_j} \to -\infty$, take $\gamma^j = \kappa_j$. Otherwise $\kappa_j = \sigma(x_0)$, in which case we are done.
7. Repeat steps 5 and 6.
8. We have

$$\lim_{j\to\infty} \kappa_j = \lim_{j\to\infty} \gamma_j = \lim_{j\to\infty} \gamma^j = \sigma(x_0).$$

REFERENCES

[1] L. ARNOLD AND B. HARRACH, *Unique shape detection in transient eddy current problems*, Inverse Problems, 29 (2013), 095004.

[2] A. BARTH, B. HARRACH, N. HYVÖNEN, AND L. MUSTONEN, *Detecting stochastic inclusions in electrical impedance tomography*, Inverse Problems, 33 (2017), 115012.

[3] T. BRANDER, *Calderón problem for the p-Laplacian: First order derivative of conductivity on the boundary*, Proc. Amer. Math. Soc., 144 (2016), pp. 177–189, https://doi.org/10.1090/proc/12681.

[4] T. BRANDER, *Calderón's Problem for p-Laplace Type Equations*, Ph.D. thesis, Department of Mathematics and Statistics, University of Jyväskylä, Jyväskylä, Finland, 2016. Available online at http://urn.fi/URN:ISBN:978-951-39-6576-1.

[5] T. BRANDER, J. ILMAVIRTA, AND M. KAR, *Superconductive and insulating inclusions for linear and non-linear conductivity equations*, Inverse Probl. Imaging, 12 (2018), pp. 91–123, https://doi.org/10.3934/ipi.2018004.

[6] T. BRANDER, M. KAR, AND M. SALO, *Enclosure method for the p-Laplace equation*, Inverse Problems, 31 (2015), 045001.

[7] A. P. CALDERÓN, *On an inverse boundary value problem*, in Seminar on Numerical Analysis and Its Applications to Continuum Physics, W. Meyer and M. Raupp, eds., Sociedade Brasileira de Matematica, 1980, pp. 65–73.

[8] A. P. CALDERÓN, *On an inverse boundary problem*, Comput. Appl. Math., 25 (2006), pp. 133–138.

[9] D. DOS SANTOS FERREIRA, C. E. KENIG, M. SALO, AND G. UHLMANN, *Limiting Carleman weights and anisotropic inverse problems*, Invent. Math., 178 (2009), pp. 119–171, https://doi.org/10.1007/s00222-009-0196-4.

[10] H. GARDE, *Comparison of linear and non-linear monotonicity-based shape reconstruction using exact matrix characterizations*, Inverse Probl. Sci. Eng., (2017), pp. 1–18.

[11] H. Garde and S. Staboulis, *Convergence and regularization for monotonicity-based shape reconstruction in electrical impedance tomography*, Numer. Math., 135 (2017), pp. 1221–1251.

[12] H. Garde and S. Staboulis, *The Regularized Monotonicity Method: Detecting Irregular Indefinite Inclusions*, preprint, https://arxiv.org/abs/1705.07372, 2017.

[13] B. Gebauer, *Localized potentials in electrical impedance tomography*, Inverse Probl. Imaging, 2 (2008), pp. 251–269, https://doi.org/10.3934/ipi.2008.2.251.

[14] C.-Y. Guo, M. Kar, and M. Salo, *Inverse problems for p-Laplace type equations under monotonicity assumptions*, Rend. Istit. Mat. Univ. Trieste, 48 (2016), pp. 79–99.

[15] B. Harrach, *On uniqueness in diffuse optical tomography*, Inverse Problems, 25 (2009), 055010, https://doi.org/10.1088/0266-5611/25/5/055010.

[16] B. Harrach, *Simultaneous determination of the diffusion and absorption coefficient from boundary data*, Inverse Probl. Imaging, 6 (2012), pp. 663–679, https://doi.org/10.3934/ipi.2012.6.663.

[17] B. Harrach, *Interpolation of missing electrode data in electrical impedance tomography*, Inverse Problems, 31 (2015), 115008.

[18] B. Harrach, E. Lee, and M. Ullrich, *Combining frequency-difference and ultrasound modulated electrical impedance tomography*, Inverse Problems, 31 (2015), 095003.

[19] B. Harrach and M. N. Minh, *Enhancing residual-based techniques with shape reconstruction features in electrical impedance tomography*, Inverse Problems, 32 (2016), 125002.

[20] B. Harrach and M. N. Minh, *Monotonicity-Based Regularization for Phantom Experiment Data in Electrical Impedance Tomography*, preprint, https://arxiv.org/abs/1610.05718, 2016.

[21] B. Harrach, V. Pohjola, and M. Salo, *Monotonicity and Local Uniqueness for the Helmholtz Equation*, preprint, https://arxiv.org/abs/1709.08756, 2017.

[22] B. Harrach and J. K. Seo, *Exact shape-reconstruction by one-step linearization in electrical impedance tomography*, SIAM J. Math. Anal., 42 (2010), pp. 1505–1518, https://doi.org/10.1137/090773970.

[23] B. Harrach and M. Ullrich, *Monotonicity-based shape reconstruction in electrical impedance tomography*, SIAM J. Math. Anal., 45 (2013), pp. 3382–3403, https://doi.org/10.1137/120886984.

[24] B. Harrach and M. Ullrich, *Resolution guarantees in electrical impedance tomography*, IEEE Trans. Med. Imaging, 34 (2015), pp. 1513–1521.

[25] B. Harrach and M. Ullrich, *Local uniqueness for an inverse boundary value problem with partial data*, Proc. Amer. Math. Soc., 145 (2017), pp. 1087–1095.

[26] M. Ikehata, *Enclosing a polygonal cavity in a two-dimensional bounded domain from Cauchy data*, Inverse Problems, 15 (1999), pp. 1231–1241, http://stacks.iop.org/0266-5611/15/i=5/a=308.

[27] M. Ikehata, *On reconstruction in the inverse conductivity problem with one measurement*, Inverse Problems, 16 (2000), pp. 785–793, https://doi.org/10.1088/0266-5611/16/3/314.

[28] M. Ikehata, *The probe method and its applications*, in Inverse Problems and Related Topics, Chapman & Hall/CRC Res. Notes Math. 419, G. Nakamura, S. Saitoh, and J. K. Seo, eds., Chapman & Hall/CRC, Boca Raton, FL, 2000, pp. 57–68.

[29] M. Ikehata, *Reconstruction of the support function for inclusion from boundary measurements*, J. Inverse Ill-Posed Probl., 8 (2000), pp. 367–378, https://doi.org/10.1515/jiip.2000.8.4.367.

[30] M. Ikehata, *Analytical methods for extracting discontinuity in inverse problems: The probe method after 10 years*, Sugaku Expositions, 26 (2013), pp. 1–28.

[31] A. Maffucci, A. Vento, S. Ventre, and A. Tamburrino, *A novel technique for evaluating the effective permittivity of inhomogeneous interconnects based on the monotonicity property*, IEEE Trans. Components Packaging Manuf. Tech., 6 (2016), pp. 1417–1427.

[32] S. Nagayasu, G. Uhlmann, and J.-N. Wang, *Reconstruction of penetrable obstacles in acoustic scattering*, SIAM J. Math. Anal., 43 (2011), pp. 189–211, https://doi.org/10.1137/09076218X.

[33] G. Nakamura and K. Yoshida, *Identification of a non-convex obstacle for acoustical scattering*, J. Inverse Ill-Posed Probl., 15 (2007), pp. 611–624, https://doi.org/10.1515/jiip.2007.034.

[34] R. T. Rockafellar, *Convex Analysis*, Princeton Math. Ser. 28, Princeton University Press, Princeton, NJ, 1970.

[35] M. Salo and X. Zhong, *An inverse problem for the p-Laplacian: Boundary determination*, SIAM J. Math. Anal., 44 (2012), pp. 2474–2495.

[36] M. SINI AND K. YOSHIDA, *On the reconstruction of interfaces using complex geometrical optics solutions for the acoustic case*, Inverse Problems, 28 (2012), 055013, https://doi.org/10.1088/0266-5611/28/5/055013.

[37] Z. SU, A. TAMBURRINO, S. VENTRE, L. UDPA, AND S. UDPA, *Time domain monotonicity based inversion method for eddy current tomography*, in Proc. 31st International Review of Progress in Applied Computational Electromagnetics (ACES), IEEE Press, Piscataway, NJ, 2015, pp. 1–2.

[38] A. TAMBURRINO, *Monotonicity based imaging methods for elliptic and parabolic inverse problems*, J. Inverse Ill-Posed Probl., 14 (2006), pp. 633–642, https://doi.org/10.1163/156939406778474578.

[39] A. TAMBURRINO AND G. RUBINACCI, *A new non-iterative inversion method for electrical resistance tomography*, Inverse Problems, 18 (2002), pp. 1809–1829. Special section on electromagnetic and ultrasonic nondestructive evaluation, https://doi.org/10.1088/0266-5611/18/6/323.

[40] A. TAMBURRINO, Z. SU, N. LEI, L. UDPA, AND S. UDPA, *The monotonicity imaging method for PECT*, in Electromagnetic Nondestructive Evaluation (XVIII), Stud. Appl. Electromagn. Mech. 40, Z. Chen, S. Xie, and Y. Li, eds., IOS Press, Amsterdam, 2015, pp. 159–166.

[41] A. TAMBURRINO, Z. SUA, S. VENTRE, L. UDPA, AND S. S. UDPA, *Monotonicity based imaging method method in time domain eddy current testing*, in Electromagnetic Nondestructive Evaluation (XIX), Stud. Appl. Electromagn. Mech. 41, N. Yusa, T. Uchimoto, and H. Kikuchi, eds., 2016, IOS Press, Amsterdam, pp. 1–8.

[42] G. UHLMANN, *Inverse problems: Seeing the unseen*, Bull. Math. Sci., 4 (2014), pp. 209–279, https://doi.org/10.1007/s13373-014-0051-9.

[43] T. H. WOLFF, *Gap series constructions for the p-Laplacian*, J. Anal. Math., 102 (2007), pp. 371–394. Preprint written in 1984; paper completed by J. Garnett and J.-M. Wu.